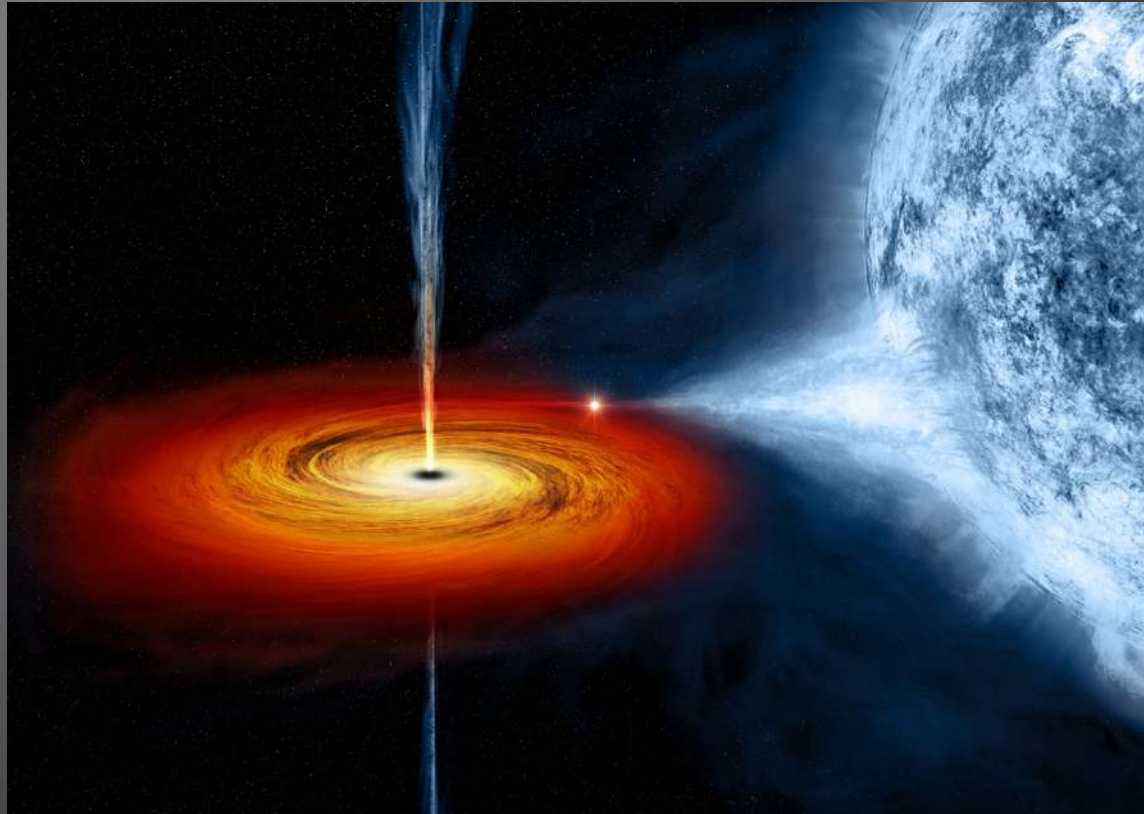


PHYSICS OF COMPACT OBJECTS AND THEIR BINARY INTERACTIONS

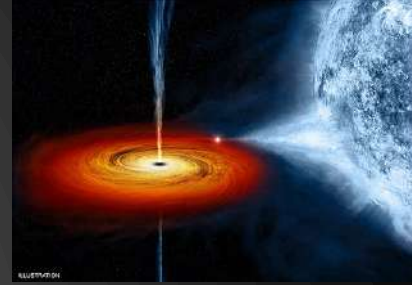


**AALBORG
UNIVERSITY**

Thomas Tauris – Physics, Aalborg University

Last Week

Summary



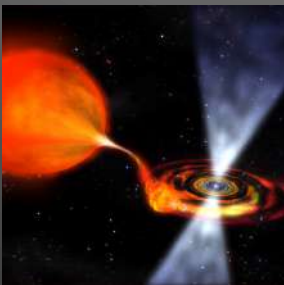
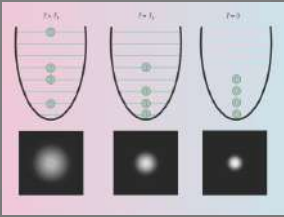
- Stellar mass BHs have spins: $0.1 < a_* < 0.99 \quad \wedge \quad a_* \equiv \frac{cJ}{GM^2}$
- BH spins can be determined via the continuum-fitting model

- TODAY!**
- BH spins can also be measured in merging BH binaries (LIGO)
 - The two classes of BHs (transient vs persistent) have different spins
 - The fast spins of the persistent BHs are natal
 - The BH spins seem to be correlated with the jet power

Literature

- Shapiro & Teukolsky (1983), Chapter 12 (14)
- McClintock, Narayan & Steiner (2013)
- Fabian & Lasenby (2015)
- Tauris & van den Heuvel (2023), Chapter 7.6

Programme

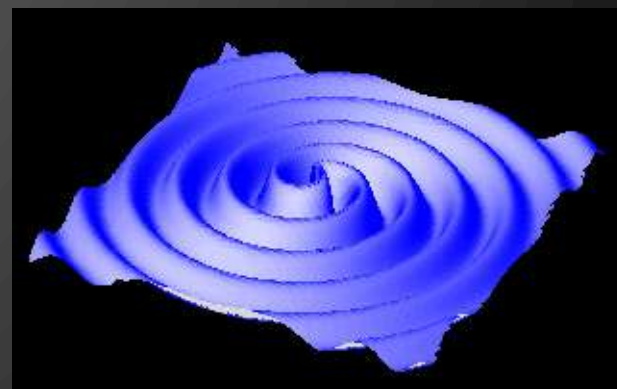


- * **Introduction**
- * **Degenerate Fermi Gases**
Non-relativistic and extreme relativistic electron / (n,p,e^-) gases
- * **White Dwarfs**
Structure, cooling models, observations
- * **Neutron Stars**
Structure and equation-of-state
- * **Radio Pulsars**
Characteristics, spin evolution, magnetars, observations
- * **Binary Evolution and Interactions**
X-ray binaries, accretion, formation of millisecond pulsars, recycling
- * **Black Holes**
Observations, characteristics and spins
- * **Gravitational Waves**
Sources and detection, kilonovae
- * **Exam**

Gravitational Waves

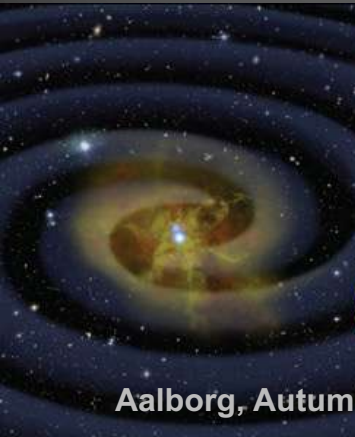
– sources and detection

- Concepts and emission of GWs
- Detection of GWs – LIGO, LISA, PTA
- Astrophysical sources
 - Burst emission sources (extra galactic)
 - Continuous emission sources (Galactic)
- Merger timescale
- GW150914 (first BHBH merger)
- GW170817 (first NSNS merger)
- Results from GWTC-3 (LIGO O1–O3)
- Kilonovae
- aLIGO detection rates
 - Population synthesis
 - Challenges



Introduction

The last 400 years of astronomy were about
"seeing" a silent movie.
LIGO is delivering the "sound track".



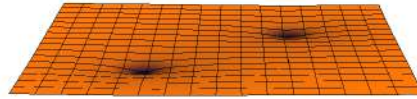
Aalborg, Autumn 2023



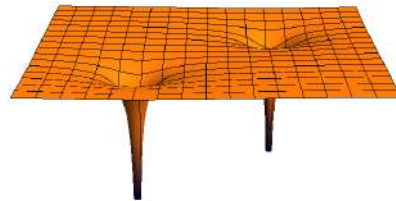
Thomas Tauris

Gravity regimes

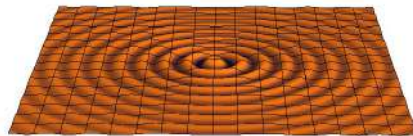
(1) Quasi-stationary weak-field regime



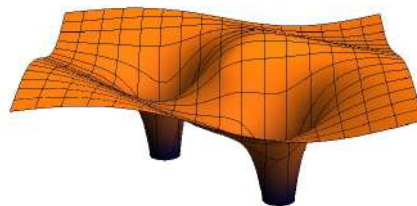
(2) Quasi-stationary strong-field regime



(3) Radiative regime



(4) Highly relativistic regime



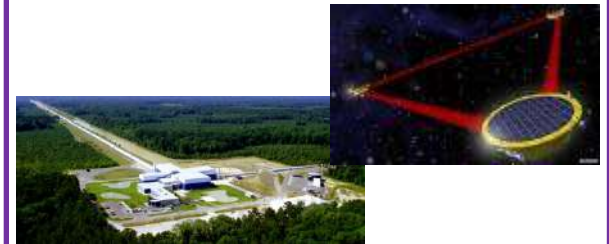
Solar system experiments



Binary pulsar experiments

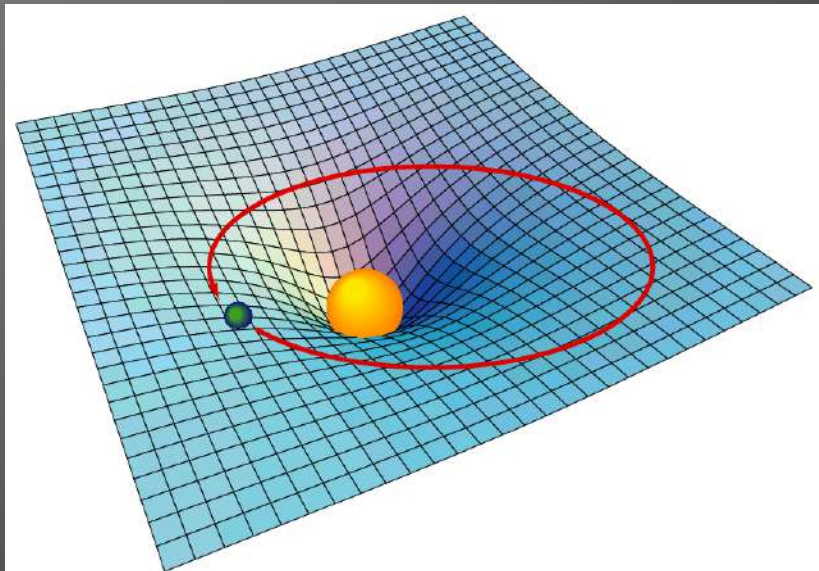


GW astronomy



General relativity in a nutshell

- Imagine space as a stretched rubber sheet
- A mass on the surface will cause a deformation
- Another mass dropped onto the sheet will roll towards that mass



*”The curvature of space determines how matter should move
- and matter determines the curvature of space”*

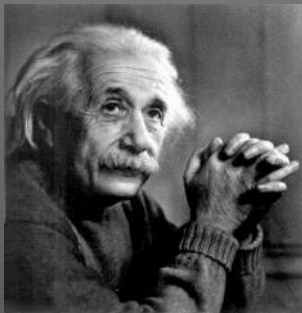
(Einstein’s field equations explained by John Wheeler)

Einstein's field equations

How does the distribution of mass-energy determine the geometry ?

$$G_{\mu\nu} = K T_{\mu\nu}$$

strain (curvature) = const. x stress (mass, energy)



space-time curvature tensor

stress-energy tensor (source term)

scalar constant "effectiveness of distorting space-time"

$$G_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu} - \Lambda g_{\mu\nu}$$

(cosmological constant)

Metric

Semi-Riemannian geometry
(curved space):

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu$$

metric tensor

$$g_{\mu\nu} = \begin{pmatrix} g_{00} & g_{01} & g_{02} & g_{03} \\ g_{10} & g_{11} & g_{12} & g_{13} \\ g_{20} & g_{21} & g_{22} & g_{23} \\ g_{30} & g_{31} & g_{32} & g_{33} \end{pmatrix}$$

Minkowski flat space:

$$ds^2 = -c^2 dt^2 + dx^2 + dy^2 + dz^2$$

(special relativity)

Weak field vacuum limit

Consider a small perturbation from a flat space-time:

Let the metric tensor be:

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$$

where

$$\eta_{\mu\nu} = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

and $|h_{\mu\nu}| \ll 1$ is a small perturbation

For a specific coordinate system it can be shown that

$$\nabla^2 h_{\mu\nu} - \frac{\partial^2 h_{\mu\nu}}{\partial t^2} = 0$$

is a solution to Einstein's field eq. (the equation for a plane wave)

Analogue to Hooke's law:

$$F = -kx$$

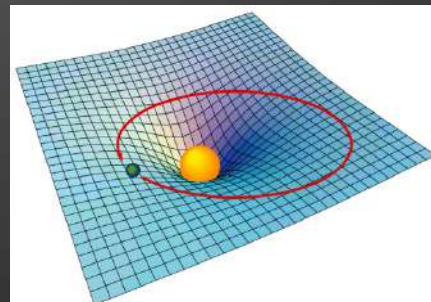
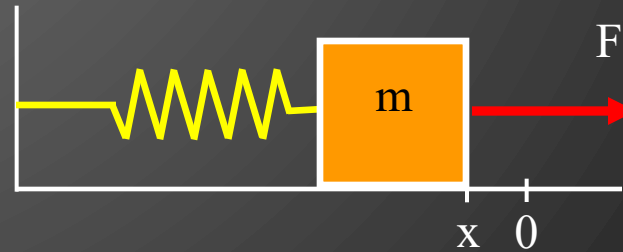
"force"

source

"displacement"

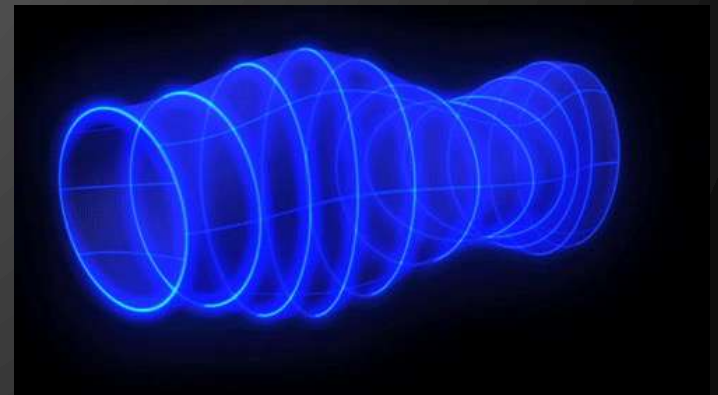
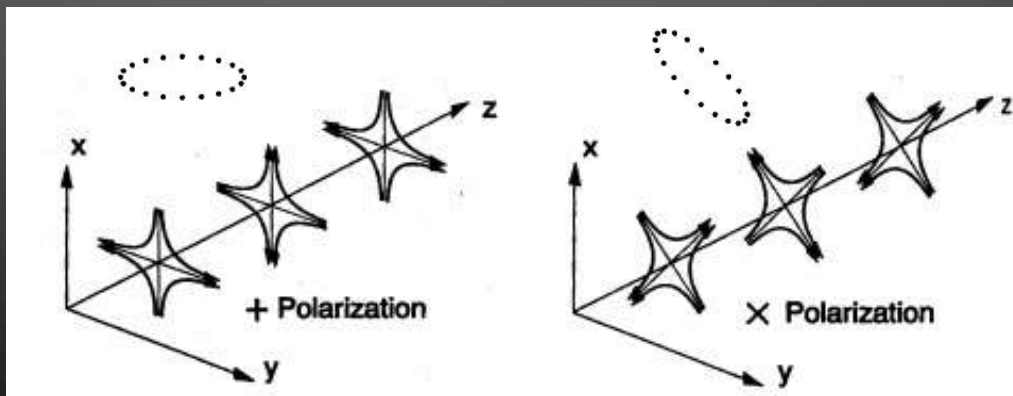
curvature

$$T_{\mu\nu} = \frac{c^4}{8\pi G} G_{\mu\nu}$$



Nature of the gravitational waves

- The emitted waves carry information of the changes in the gravitational field of the source as a result of a change in the distribution of mass, energy and momentum
- Gravitational waves propagate with the speed of light (the graviton has zero rest mass)
- They give rise to fluctuations in the metric where they pass through
- The waves' force field is transverse to its propagation direction and has quadrupolar symmetry (i.e. the graviton has $S=2$)



Gravitational wave emission

A time-varying quadrupole moment* gives rise to emission of gravitational waves with a strain amplitude:

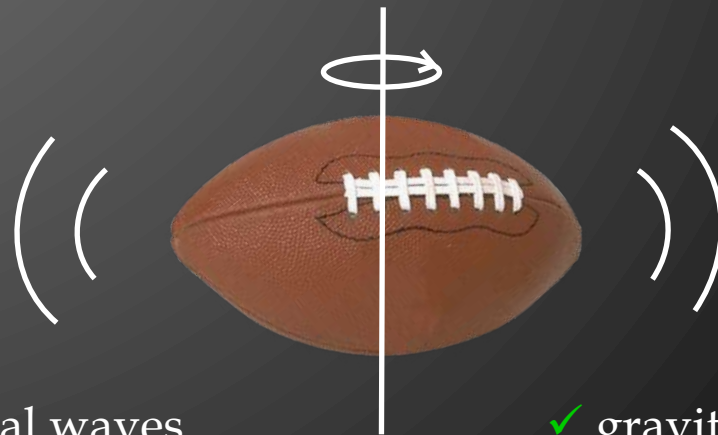
$$h_{\mu\nu} \approx \frac{2G}{c^4 d} \ddot{Q}_{\mu\nu}$$

quadrupole moment
distance to source
(Newtonian/quadrupole approximation)

* an asymmetric distribution of mass with respect to the rotation axis:



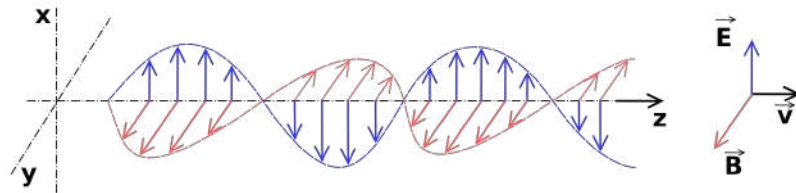
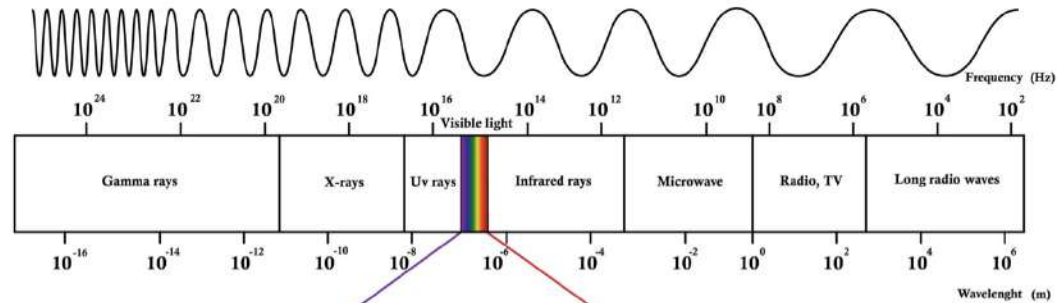
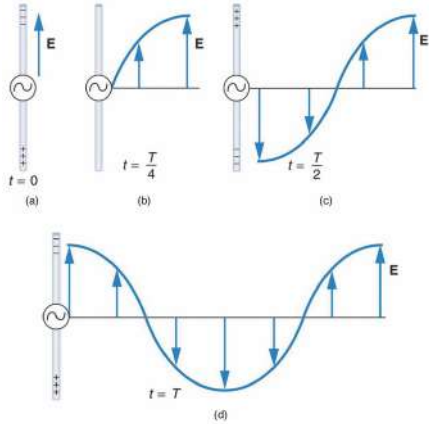
□ no gravitational waves



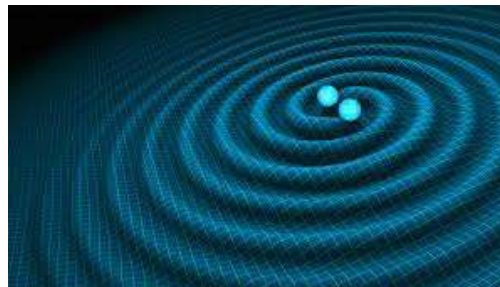
✓ gravitational waves

Gravitational waves – How are they created?

Acceleration of charged particles → electromagnetic waves:



Acceleration of masses → gravitational waves:



The physical meaning of "h"

Remember:

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu = (\eta_{\mu\nu} + h_{\mu\nu}) dx^\mu dx^\nu$$

Consider the following geometry:

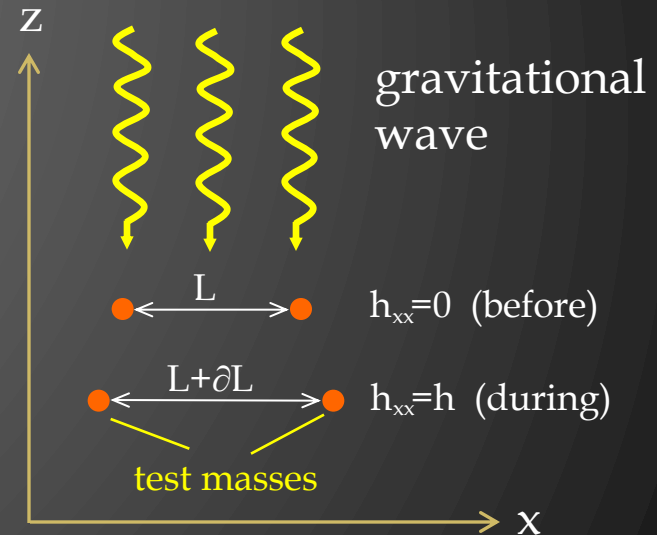
$$ds = L + \delta L \quad dx = L$$

$$dy = dz = dt = 0$$

$$h_{xx} = h$$



$$(L + \delta L)^2 = (1 + h)L^2 \quad \Leftrightarrow \quad \frac{\delta L}{L} \approx \frac{h}{2}$$



the wave (strain) amplitude is twice the relative length change

The effect on Leonardo da Vinci

Beware, only space is deformed - not matter!


wave



$$\frac{\delta L}{L} = \frac{1}{2}h \approx 10^{-21} \quad \text{for many astrophysical sources}$$

NS-NS collision at 200 Mpc

The value of h for astrophysical sources

$$h_{\mu\nu} = \frac{2G}{c^4 d} \ddot{Q}_{\mu\nu}$$

order-of-magnitude estimate

$$Q \sim MR^2 \Rightarrow \ddot{Q} \sim \frac{MR^2}{T^2} \sim Mv^2$$

where (M,R,T,v) are characteristic values of the source

$$h \sim \frac{2G}{c^4 d} Mv^2 = \frac{2G}{c^2} \frac{M}{d} \left(\frac{v}{c}\right)^2 \sim 1.0 \times 10^{-19} \frac{M / M_{\odot}}{d / \text{Mpc}} \left(\frac{v}{c}\right)^2$$

$$h = \begin{cases} 10^{-17} & \text{at outskirts of our Milky Way (10 kpc)} \\ 10^{-20} & \text{at the Virgo cluster of galaxies (15 Mpc)} \\ 10^{-21} & \text{at 200 Mpc} \\ 10^{-22} & \text{at the Hubble distance (3 Gpc)} \end{cases}$$

Gravitational wave luminosity

$$h_{\mu\nu} = \frac{2G}{c^4 d} \ddot{Q}_{\mu\nu}$$

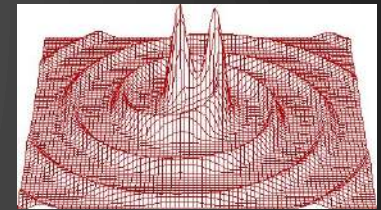
wave amplitude



$$F = \frac{c^3}{32\pi G} \langle \dot{h}_{\mu\nu} \dot{h}_{\mu\nu} \rangle$$

inner product (scalar)

energy flux detected



$$L = r^2 \int F d\Omega$$

luminosity



$$L_{\text{gwr}} \equiv \frac{dE}{dt} = \frac{G}{5c^5} \langle \ddot{Q}_{\mu\nu} \ddot{Q}_{\mu\nu} \rangle$$

Gravitational wave luminosity

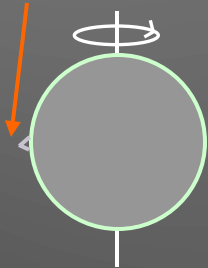
order-of-magnitude estimate

$$L_{\text{gwr}} \equiv \frac{dE}{dt} = \frac{G}{5c^5} \langle \ddot{Q}_{\mu\nu} \ddot{Q}_{\mu\nu} \rangle$$

$$Q \sim MR^2 \Rightarrow \ddot{Q} \sim \frac{MR^2}{T^3} \sim \frac{Mv^3}{R}$$

where (M,R,T,v) are characteristic values of the source

1 mm mountain on a neutron star:



M=1.4 M_⊙
R=10 km, Δr=1 mm
Ω=2πν=1000 rad/s

$$L_{\text{gwr}} = 10^{36} \text{ erg/s}$$

steel rod:



M=140.000 kg
R=20 meters
v=300 m/s

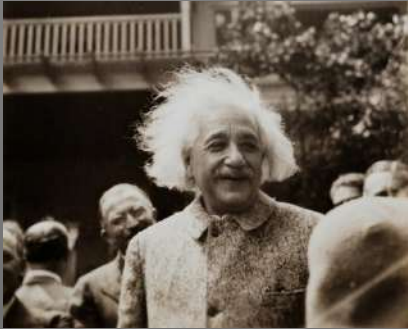
$$L_{\text{gwr}} = 10^{-24} \text{ erg/s}$$

$$L_{\text{gwr}} \equiv \frac{|dE|}{dt} = \frac{32G}{5c^5} I^2 \varepsilon^2 \Omega^6$$

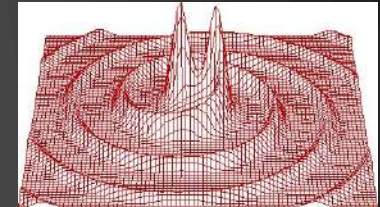
$$\varepsilon = \frac{a-b}{(a+b)/2} \approx \frac{\Delta r}{R} \quad I \sim MR^2$$

a factor of $\sim 10^{60}$ in difference!!!

Merging neutron star / black hole binaries



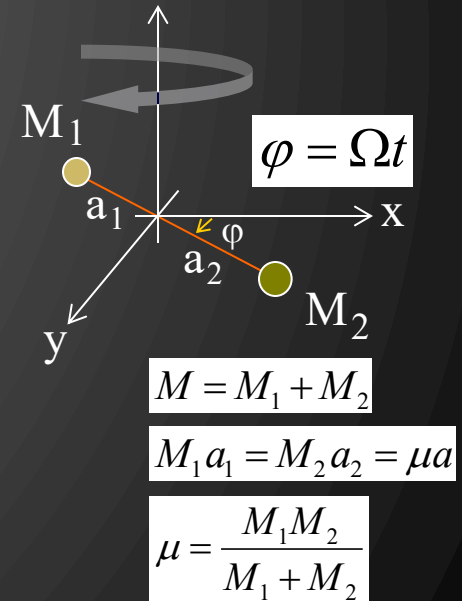
$$L_{\text{gwr}} \cong \frac{G}{5c^5} \langle \ddot{Q}_{\mu\nu} \ddot{Q}_{\mu\nu} \rangle$$



$$Q = \frac{1}{2} \mu a^2 \begin{pmatrix} \cos(2\varphi) + \text{const} & \sin(2\varphi) + \text{const} \\ \sin(2\varphi) + \text{const} & -\cos(2\varphi) + \text{const} \end{pmatrix}$$

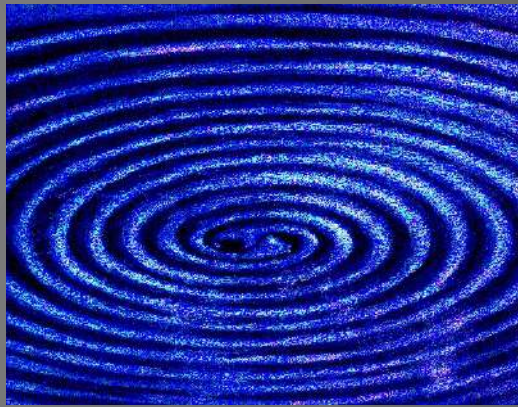
$$L_{\text{gwr}}(n, e) = \frac{32}{5} \frac{G^4}{c^5} \frac{M^3 \mu^2}{a^5} g(n, e)$$

Fourier decomposition factor
(harmonic number, eccentricity)



$$a = -\frac{GM\mu}{2E_{\text{orb}}} \Rightarrow \dot{a} = \frac{GM\mu}{2E_{\text{orb}}^2} \dot{E}_{\text{orb}} \quad \wedge \quad |\dot{E}_{\text{orb}}| = L_{\text{gwr}}$$

orbital decay!



$$\dot{a} = -\frac{GM\mu}{2E_{orb}^2} L_{gwr}$$

$$\frac{1}{a} \frac{da}{dt} = -\frac{1}{E} \frac{dE}{dt} \Big|_{e=0} f(e)$$

$$L_{gwr} \cong \frac{32}{5} \frac{G^4}{c^5} \frac{M^3 \mu^2}{a^5} \frac{1 + (73/24)e^2 + (37/96)e^4}{(1-e^2)^{7/2}}$$

$$\dot{a} \cong \frac{64}{5} \frac{G^3}{c^5} \frac{M^2 \mu}{a^3} \frac{1 + (73/24)e^2 + (37/96)e^4}{(1-e^2)^{7/2}}$$

⇓

$$\tau(a_0, e_0) \cong \frac{12}{19} \frac{C_0^4}{\beta} \int_0^{e_0} \frac{e^{29/19} [1 + (121/304)e^2]^{1181/2299}}{(1-e^2)^{3/2}} de$$

Merger timescale

determine C_0 from initial condition: $a=a_0, e=e_0$

$$\beta \equiv \frac{64}{5} \frac{G^3}{c^5} M^2 \mu$$

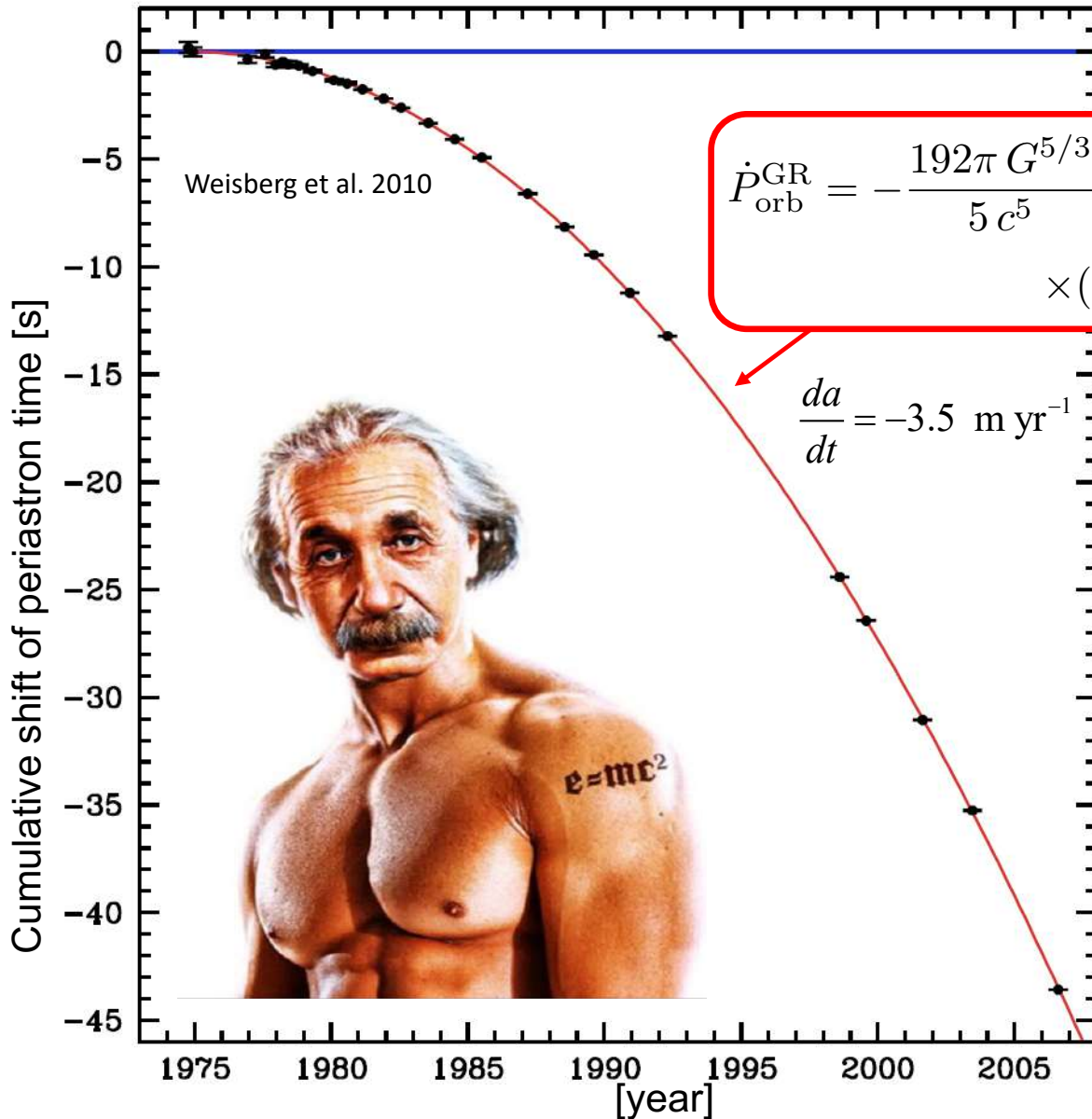
$$a(e) = \frac{C_0 e^{12/19}}{(1-e^2)} [1 + (121/304)e^2]^{870/2299}$$

For a circular binary:

$$\tau_{gwr}^{circ} = \frac{a_0^4}{4\beta}$$

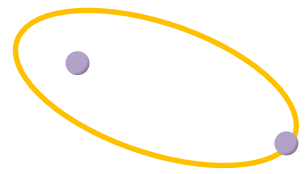
Gravitational waves do exist!

PSR B1913+16 (Hulse-Taylor pulsar, $P_{\text{orb}}=7.75$ hr, $\text{ecc}=0.61$)



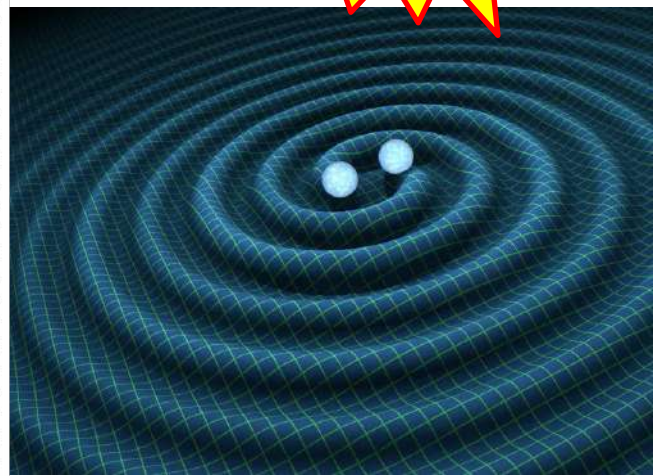
$$\dot{P}_{\text{orb}}^{\text{GR}} = -\frac{192\pi G^{5/3}}{5c^5} \left(\frac{P_{\text{orb}}}{2\pi}\right)^{-5/3} \left(1 + \frac{73}{24}e^2 + \frac{37}{96}e^4\right) \times (1 - e^2)^{-7/2} M_1 M_2 (M_1 + M_2)^{-1/3}$$

$$\frac{da}{dt} = -3.5 \text{ m yr}^{-1}$$

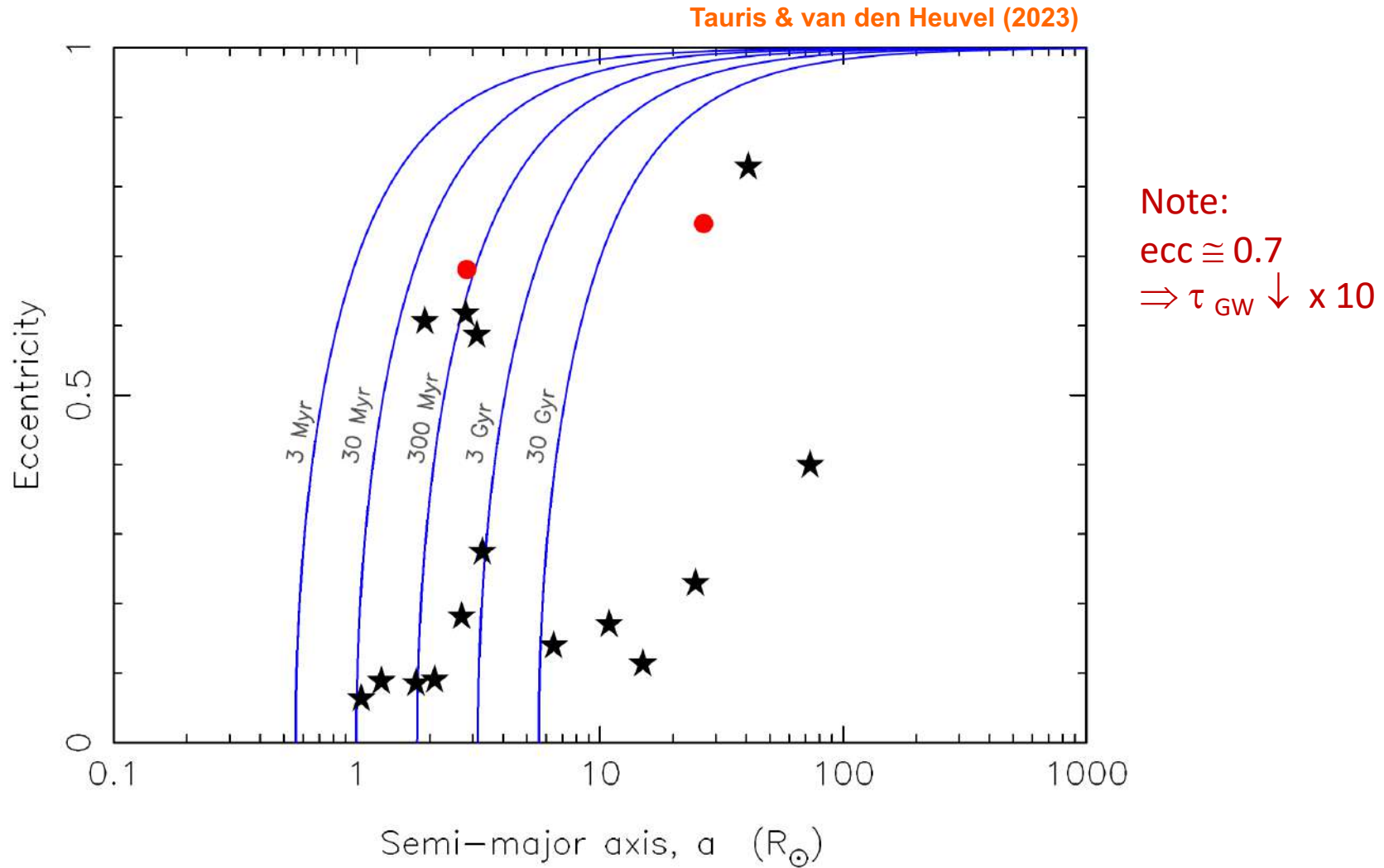


Peters (1964)

Collision in $\tau = 301$ mill. yr



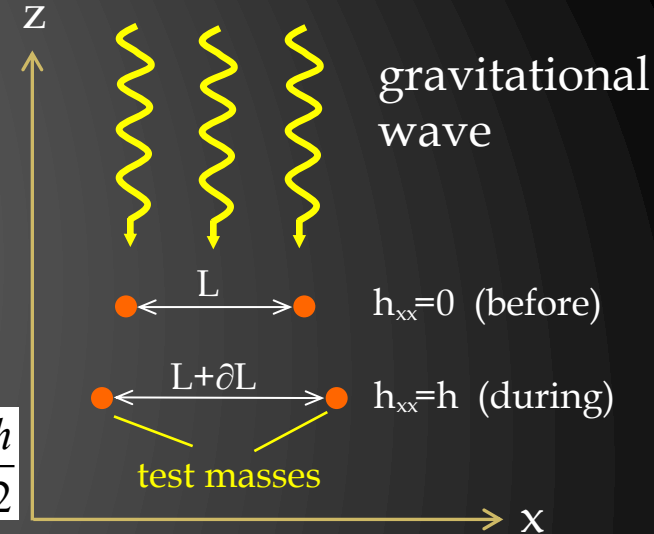
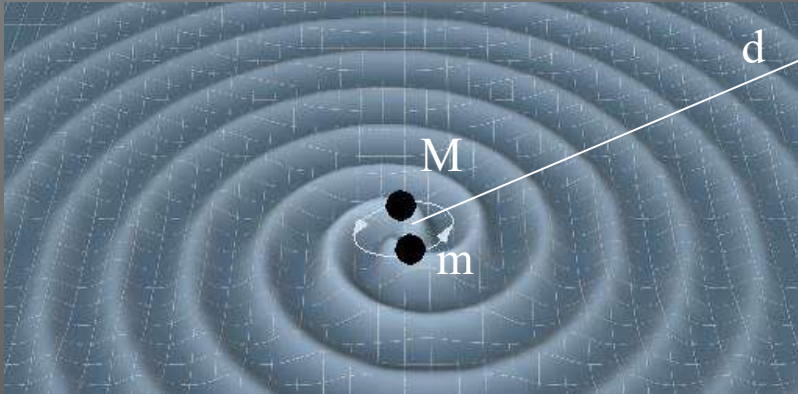
$$L_{\text{GW}} = 7 \times 10^{24} \text{ W} \quad (L_{\text{GW},\odot} = 5000 \text{ W})$$



Gravitational wave detection

$$f_{\text{gwr}} = 2 f_{\text{orb}} \quad (\text{ecc.}=0)$$

GWs in second harmonic ($n=2$)



the wave amplitude is twice the relative length change

strain (amplitude):

$$h(n, e) = \left(\frac{1}{2} [h_{+, \text{max}}^2 + h_{\times, \text{max}}^2] \right)^{1/2} = \left[\frac{16\pi G}{c^3 \omega_{\text{gwr}}^2} \frac{L_{\text{gwr}}(n, e)}{4\pi d^2} \right]^{1/2}$$

$$= 1.0 \times 10^{-21} \frac{\sqrt{g(n, e)}}{n} \left(\frac{Mm (M+m)^{-1/3}}{M_{\odot}^{5/3}} \right) \left(\frac{P_{\text{orb}}}{1 \text{ hr}} \right)^{-2/3} \left(\frac{d}{1 \text{ kpc}} \right)^{-1}$$

scale factor (Fourier decomposition factor)

massive
tight
nearby

LIGO

Laser Interferometer Gravitational wave Observatory

aLIGO observations began in 2015



aVIRGO observations began in 2016



wave amplitude: $h \sim 10^{-21}$

measurement accuracy to achieve: $\Delta L \sim 10^{-16}$ cm
(1/1000 the diameter of a proton!)

\Rightarrow interferometer arm length: $L = \Delta L / h \sim 4$ km



KAGRA joined in 2020



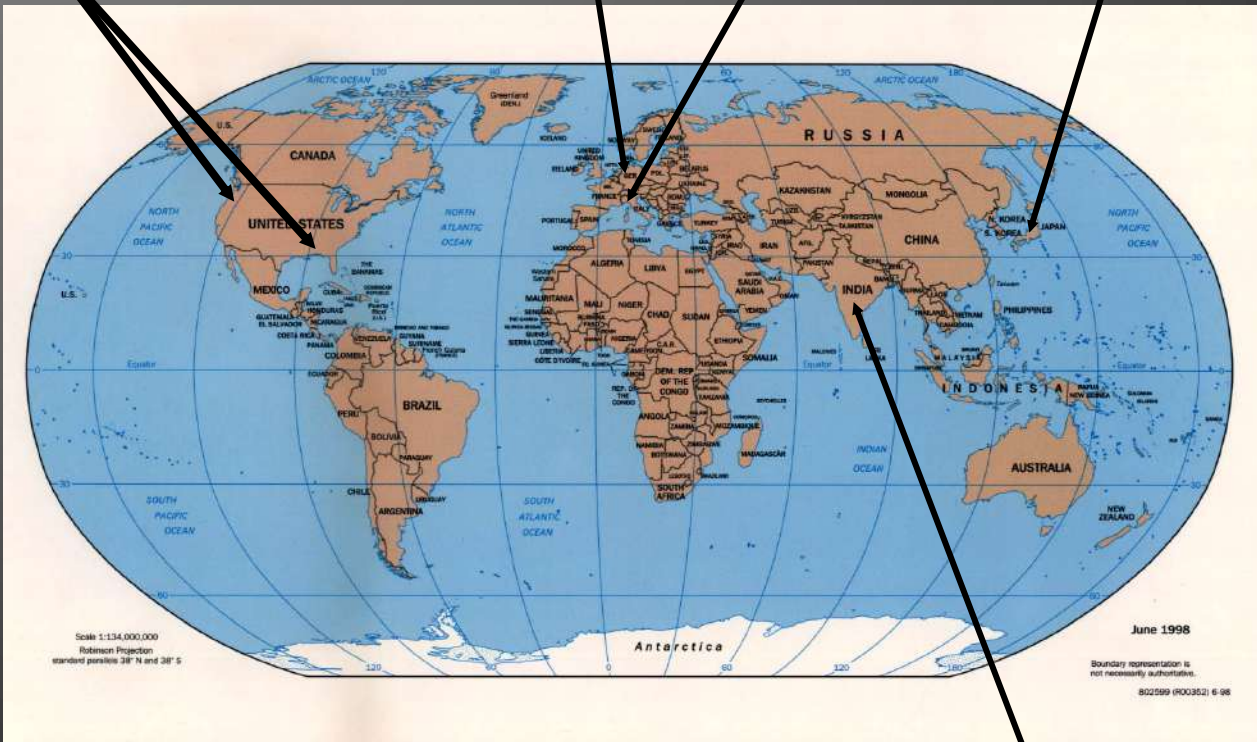
Gravitational wave observatories

LIGO

GEO

Virgo

KAGRA



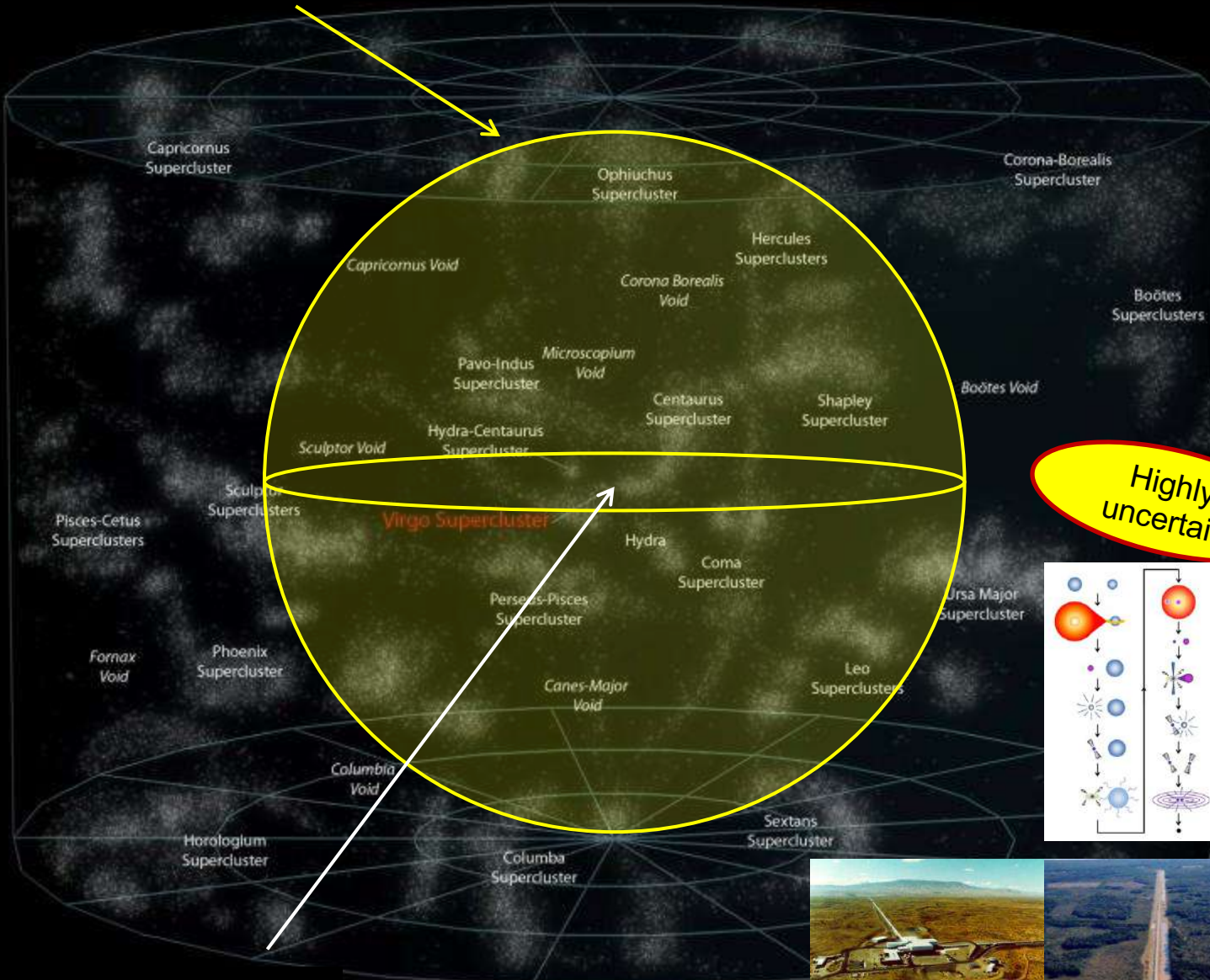
Angular resolution below 5-10 deg².

INDIGO

Merger rate of double neutron star binaries in a Milky Way-like galaxy: $3\text{--}10 \text{ Myr}^{-1}$



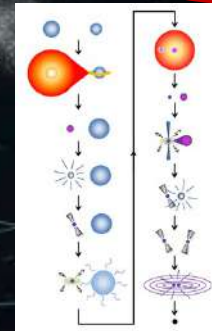
advanced LIGO



Range:
NSNS merger 200 Mpc
NSBH merger 600 Mpc
BHBH merger 5.0 Gpc
($Z=0.8$)

Detection rate:
3-10 per year
(Milky Way: 3-10 Myr⁻¹)

Highly uncertain

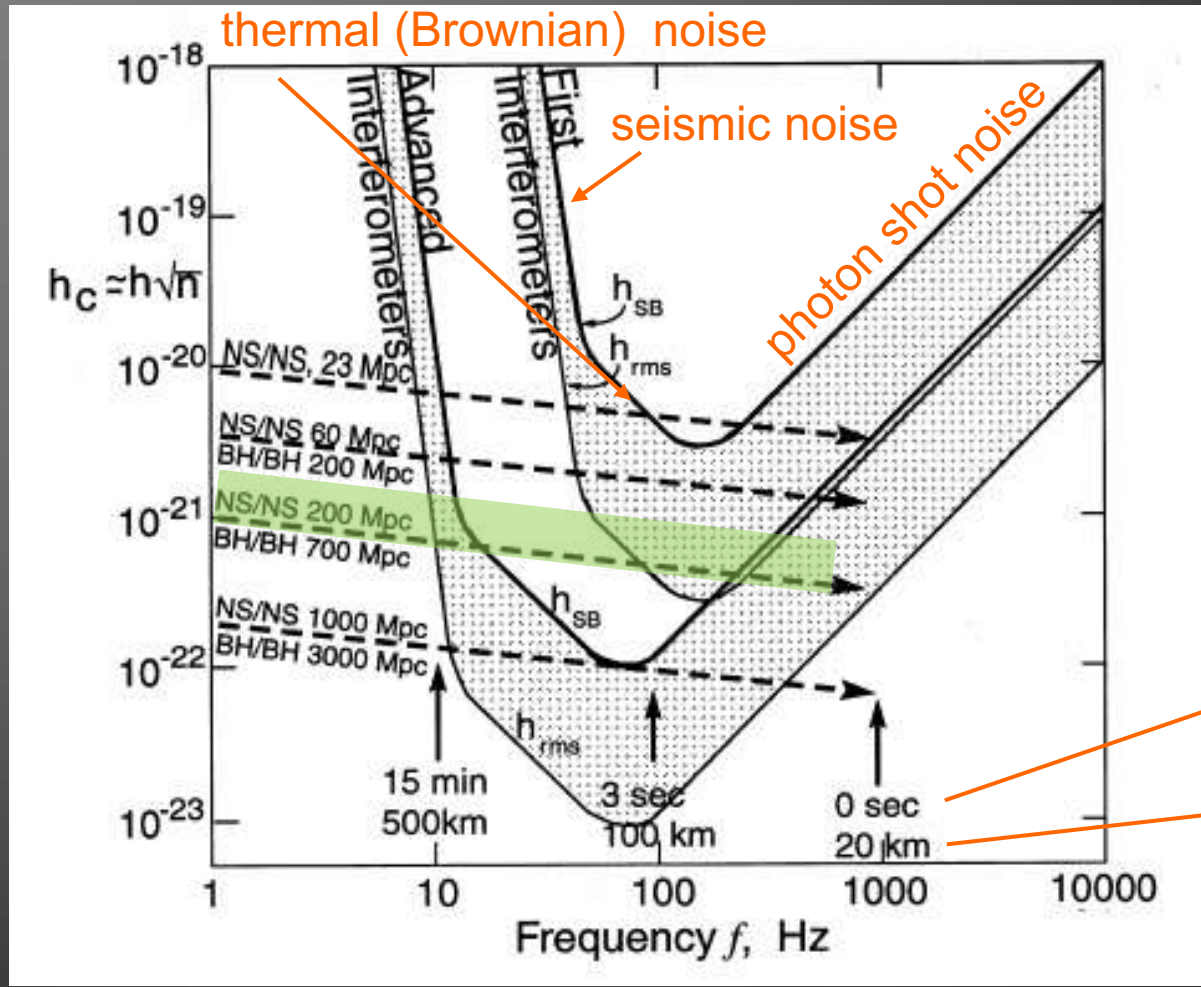


Main uncertainties:
CE evolution, kicks

You are here!



Sensitivity of LIGO/VIRGO

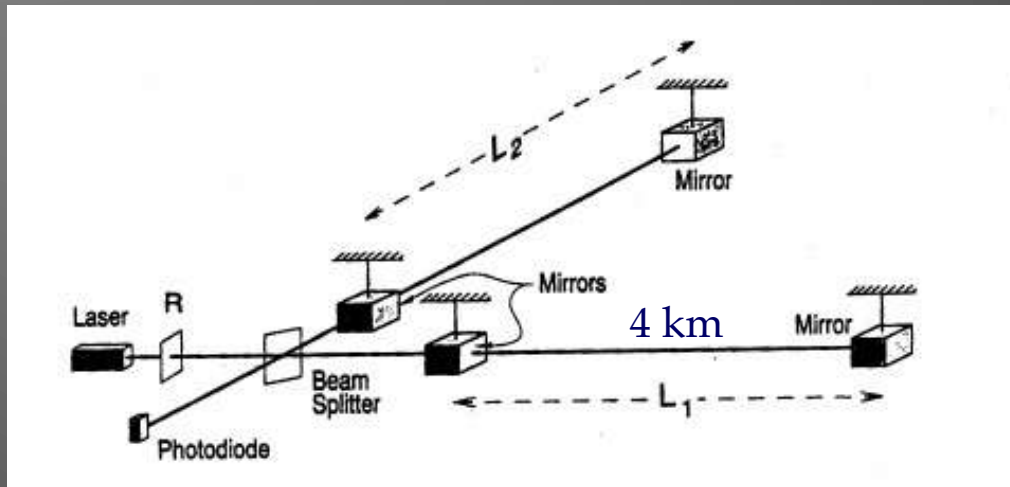


merger time

separation

Laser interferometer

- how to achieve $\Delta L \sim 10^{-16}$ cm ?



Michelson-Morley
interferometer

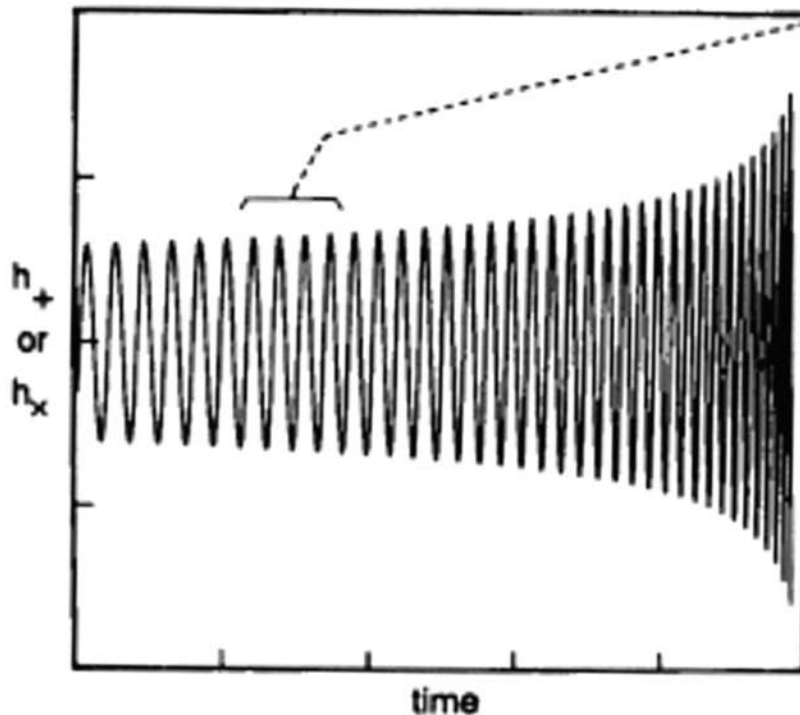
Sensitivity noise: photon, thermal and seismic.
5 cm wide laser beam shining on 10^{17} atoms.
Light is reflected 100 times \longleftrightarrow

$$\Delta I_{pd} \propto \Delta \Phi \propto \Delta L \propto h(t)$$

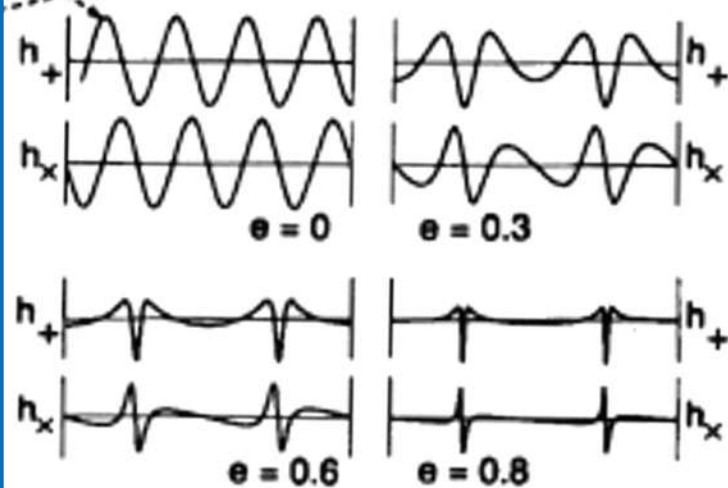


Waveform and dependence on inclination and eccentricity

WAVEFORM



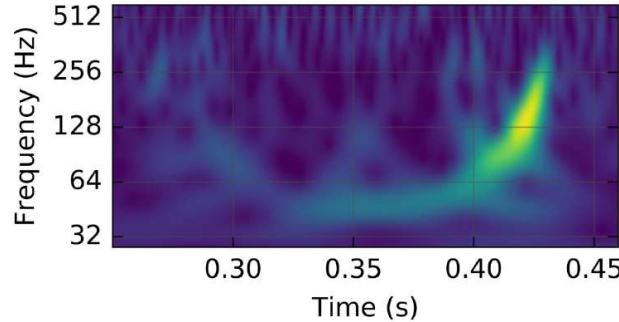
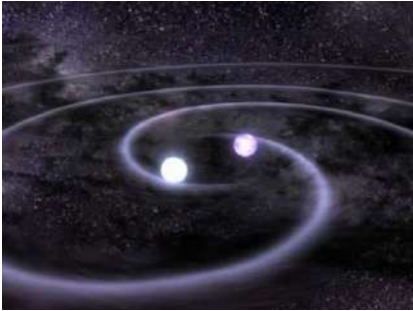
DEPENDENCE ON e , FOR $\iota = 90^\circ$:



DEPENDENCE ON ι , FOR $e = 0$:

$$\frac{\text{Amp}(h_x)}{\text{Amp}(h_+)} = \frac{2 \cos \iota}{1 + \cos^2 \iota}$$

Chirp mass



$$\mathcal{M} = \frac{(m_1 m_2)^{3/5}}{M^{1/5}} \simeq \frac{c^3}{G} \left[\frac{5}{96} \pi^{-8/3} f^{-11/3} \dot{f} \right]^{3/5} \quad (\mathcal{M}^{5/3} = \mu M^{2/3})$$

The chirp mass determines how fast the signal sweeps ("chirps") through the frequency band. It can be determined to within 1% error.

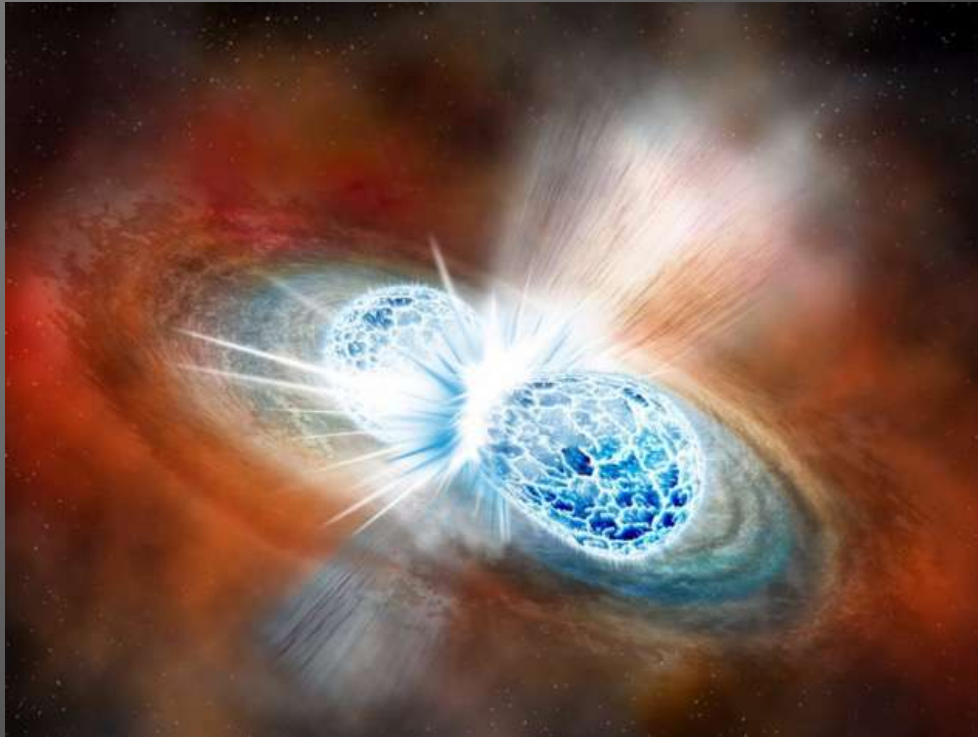
$$f_{GW}(t) = \frac{1}{8\pi} 5^{3/8} \left[\frac{c^3}{GM_{chirp}} \right]^{5/8} \frac{1}{(t_{merge} - t)^{3/8}}$$

$$h(t) = \frac{1}{r} \left[\frac{5G^5 M_{chirp}^5}{c^{11}} \right]^{1/4} \frac{1}{(t_{merge} - t)^{1/4}}$$

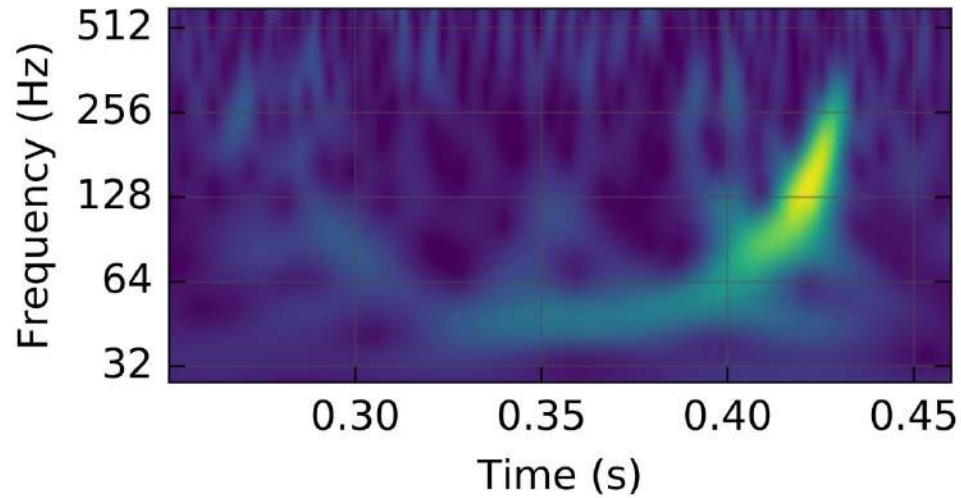
Riles (2013)

Note, this expression breaks down as
 $r \rightarrow 0$ ($v/c \rightarrow 1$)

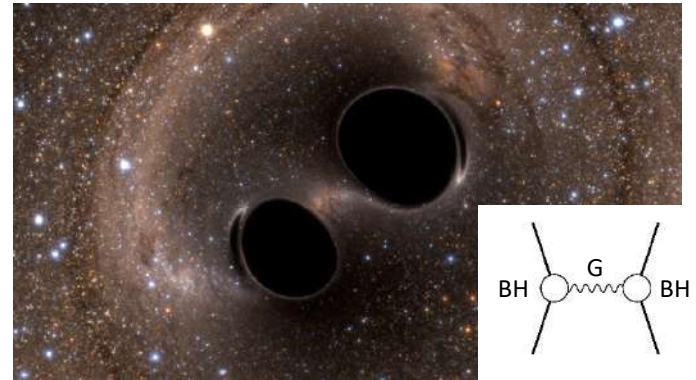
Detected Black Hole and Neutron Star Mergers



A brief chirp from a galaxy 1.4 billion light years away...

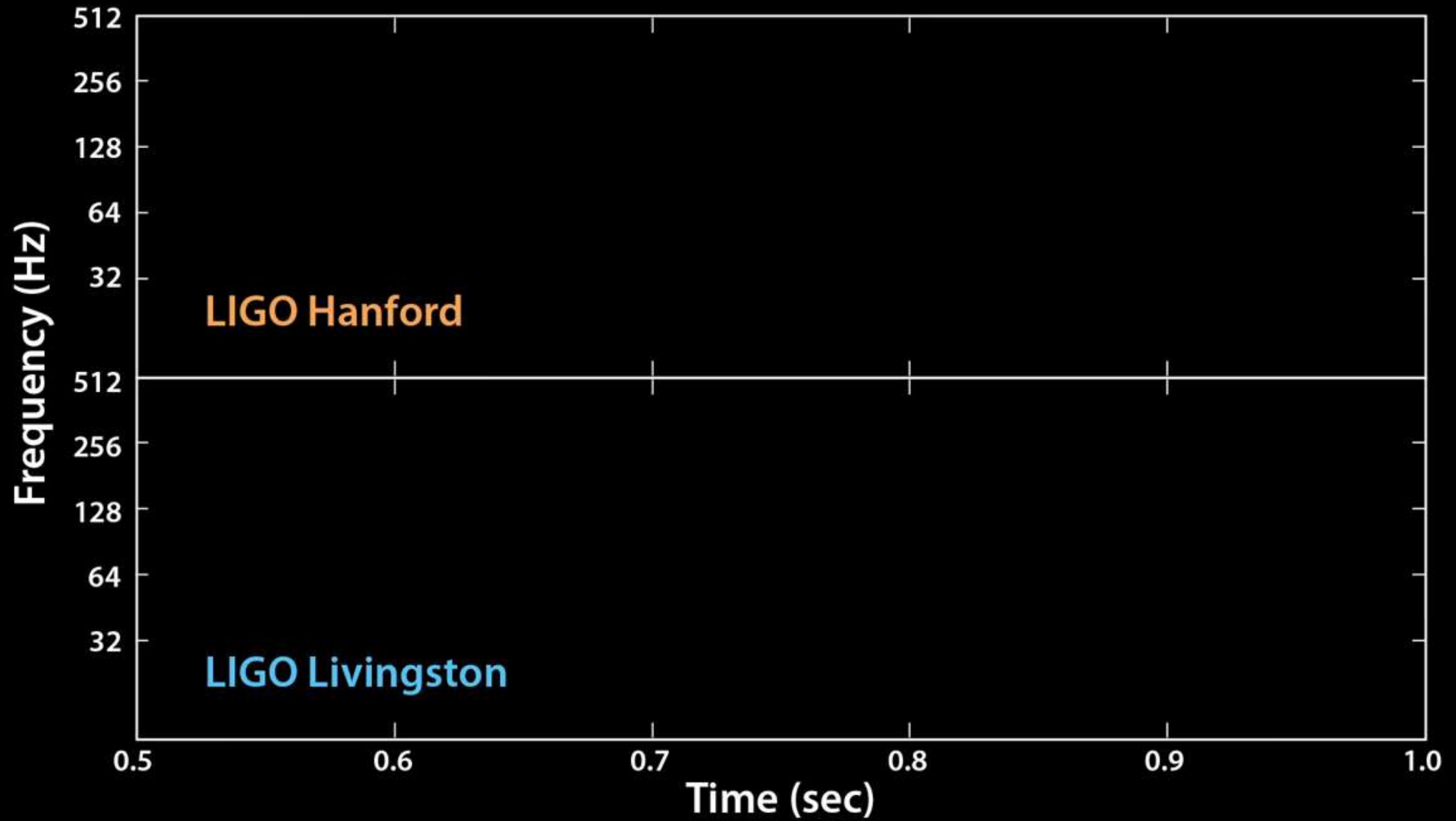


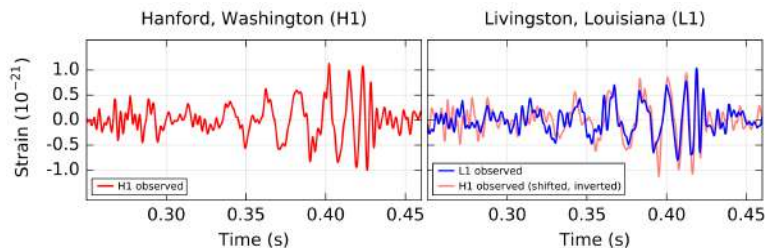
Detection of gravitational waves!



GW150914, GW151012, GW151226,
 GW170104, GW170608, GW170729,
 GW170809, GW170814, **GW170817**,
 GW170818, GW170823, (O3.....)



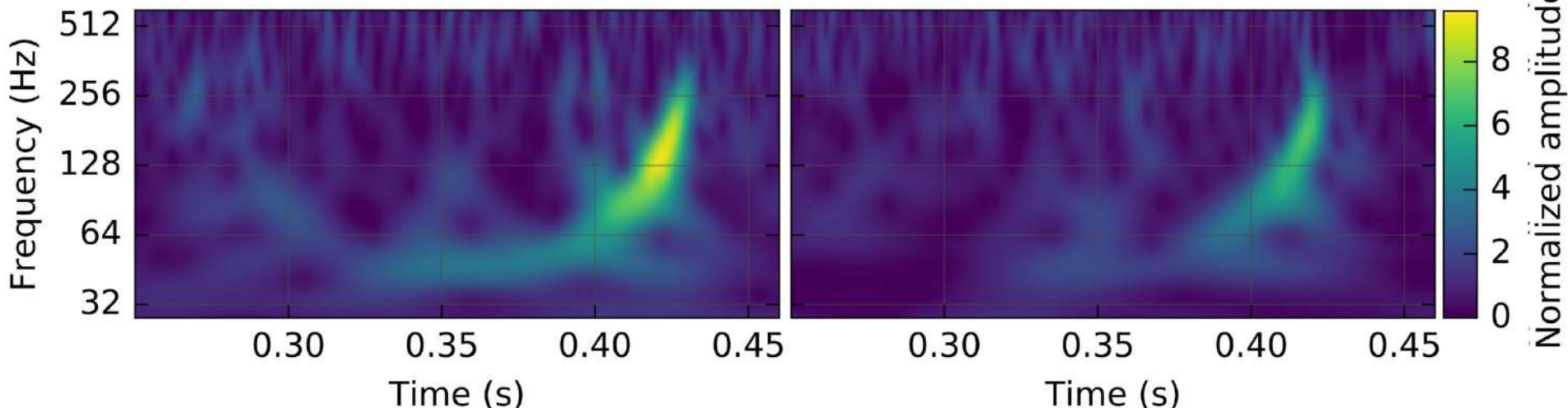




$$\mathcal{M} = \frac{(m_1 m_2)^{3/5}}{M^{1/5}} \simeq \frac{c^3}{G} \left[\frac{5}{96} \pi^{-8/3} f^{-11/3} \dot{f} \right]^{3/5}$$

Hanford, Washington (H1)

Livingston, Louisiana (L1)



Collision of two black holes:

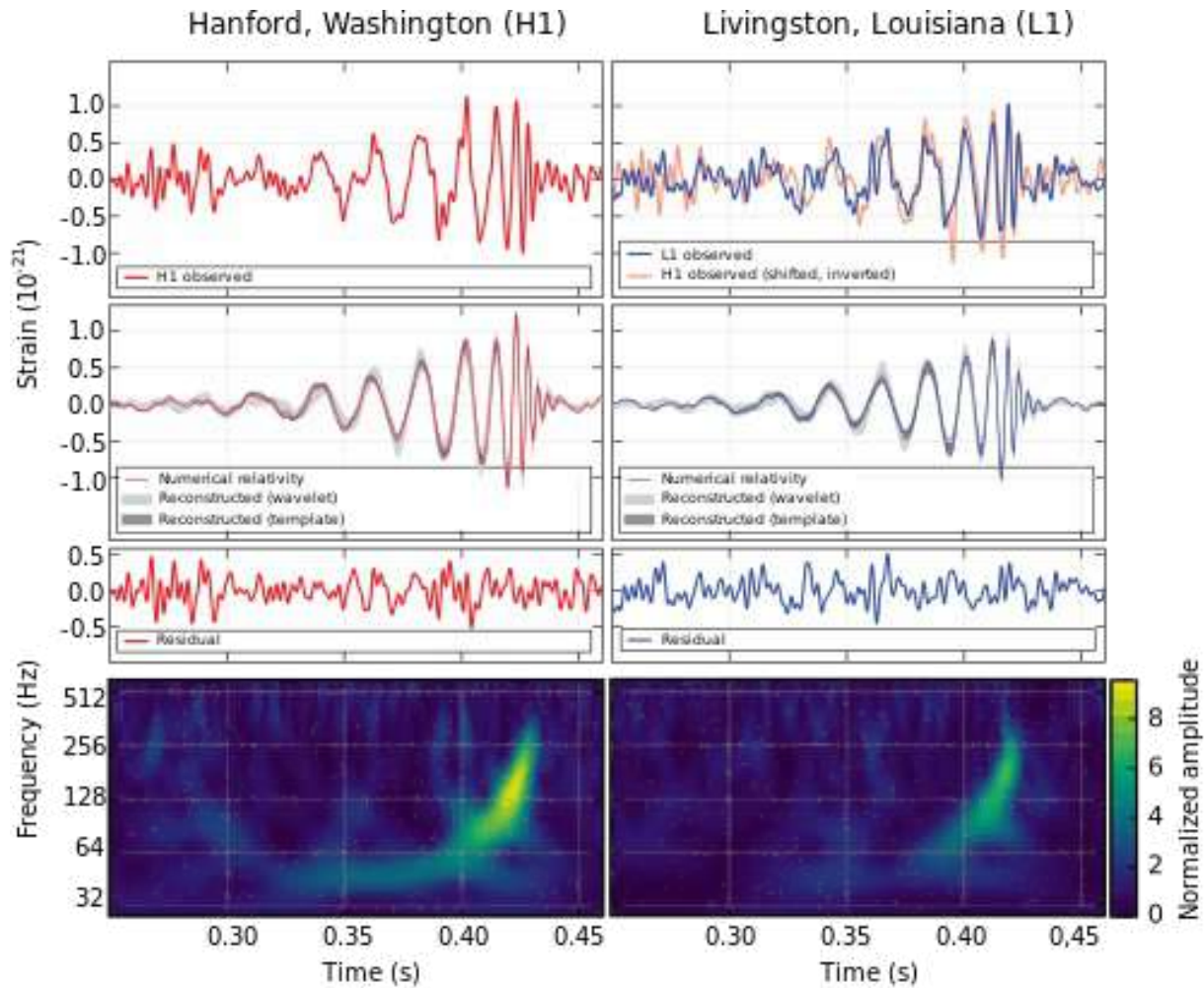
$36 M_{\text{sun}} + 29 M_{\text{sun}} \neq 65 M_{\text{sun}}$

$M_{\text{merger}} = 62 M_{\text{sun}}!$ (3 $M_{\text{sun}} \cdot c^2$ emitted as GWs)

35–250 Hz (8 cycles)

$h_{\text{top}} = 3.4 \times 10^{-22}$

Z=0.09 (400 Mpc = 1.4 bill. light years)



$$\mathcal{M} = \frac{(m_1 m_2)^{3/5}}{M^{1/5}} \simeq \frac{c^3}{G} \left[\frac{5}{96} \pi^{-8/3} f^{-11/3} \dot{f} \right]^{3/5}$$

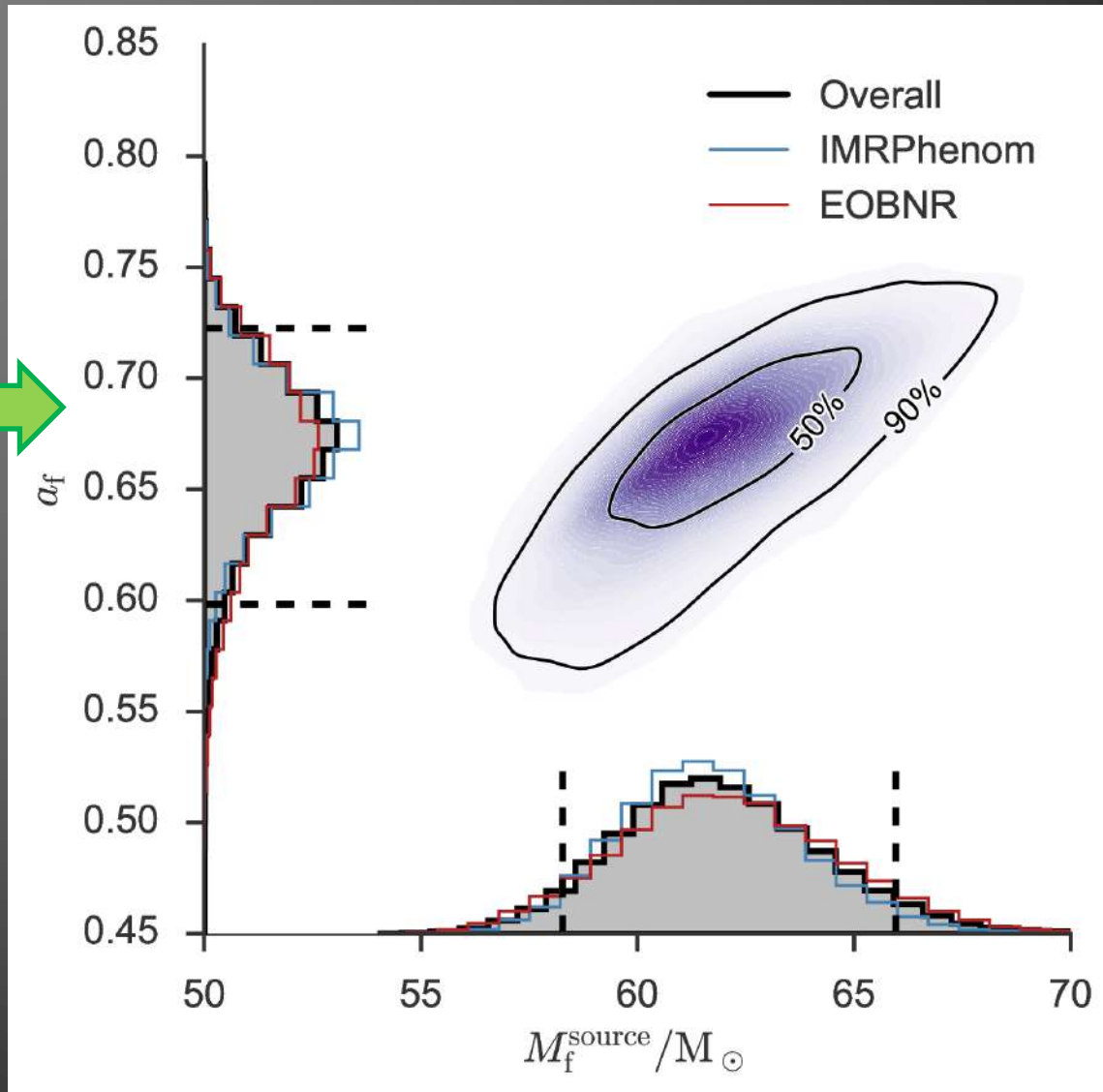
Chirp mass = $29 M_{\text{sun}}$

Must be BH-BH merger!
 (also the GW freq. at merger points to radii equal to Schwarzschild radii of massive BHs)

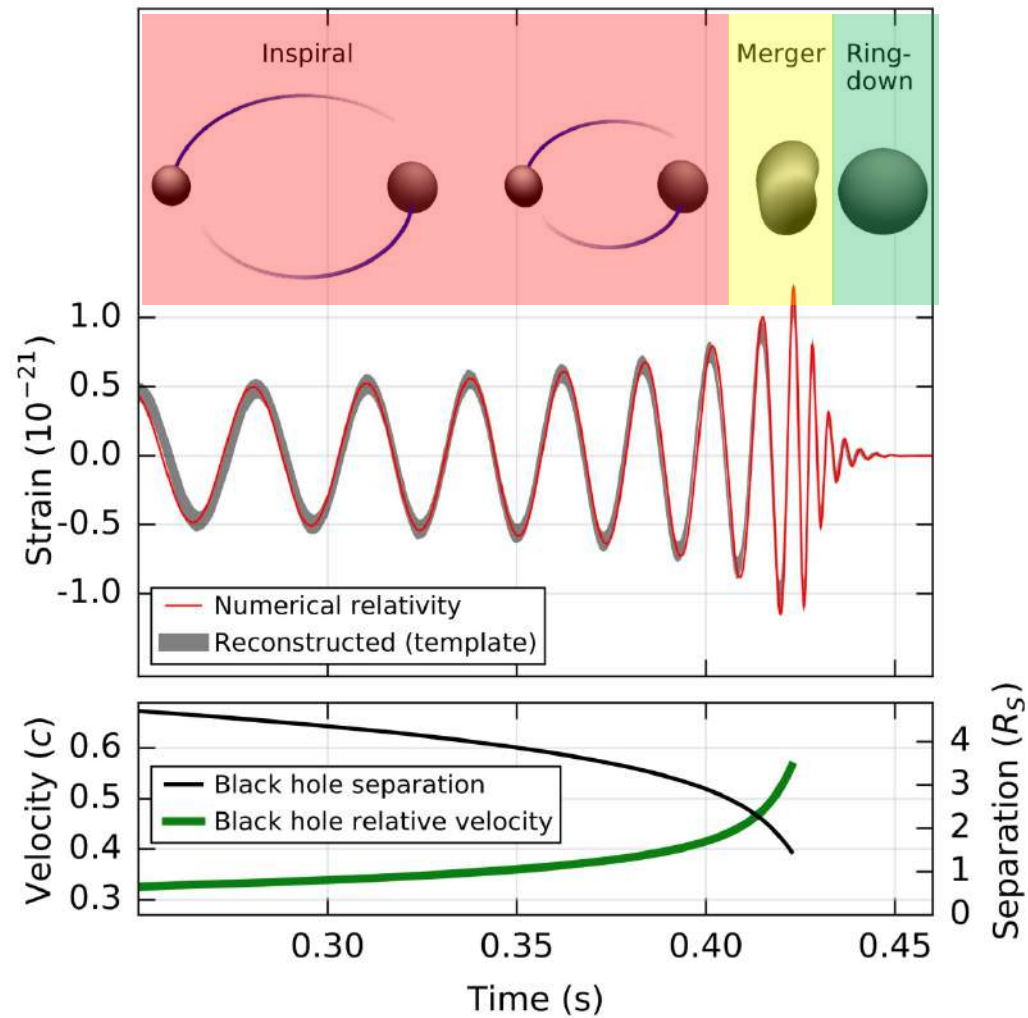
GW150914 and the revolution

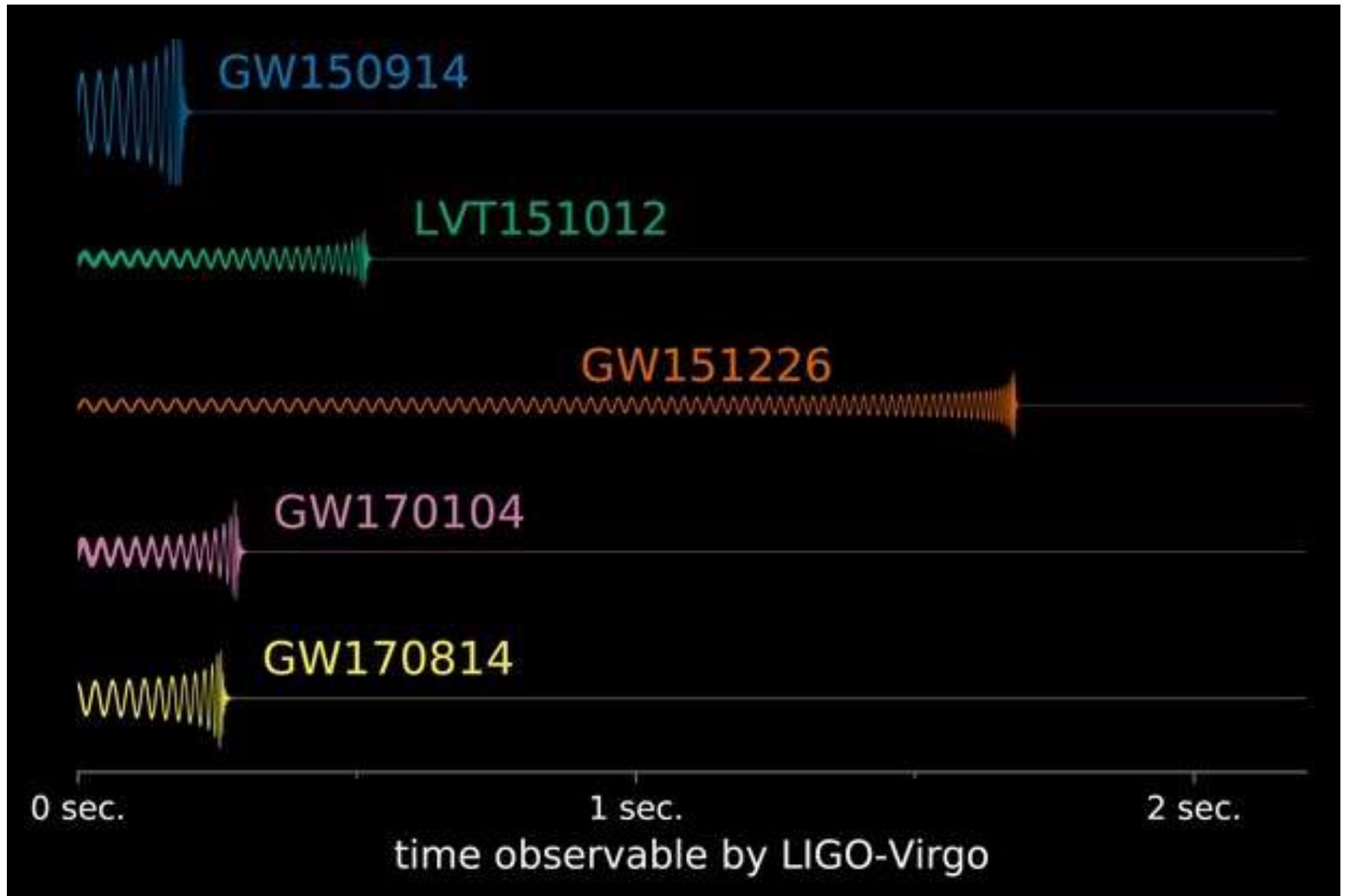
The ring-down

Typical value



3 phases:





11* events in O1 and O2

10 BH+BH mergers

1 NS+NS merger

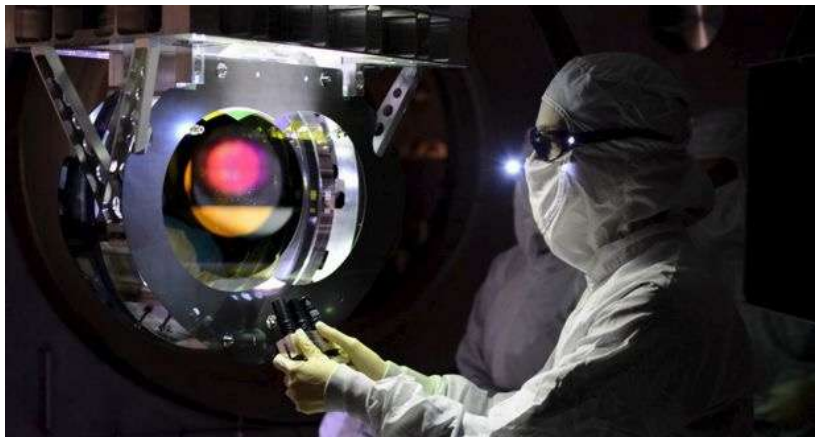


~ 80 events in O3a + O3b

?	x	BH+BH mergers	75
	y	NS+NS mergers	1(3)
	z	BH+NS mergers	2(3)
		mass gap	2(3)

- We predict ~10 times more detections of mixed BH/NS mergers compared to double NS mergers

Observational selection bias against?



Mainly BHBH mergers
(selection effect: detection vol. $\sim M_{\text{chirp}}^{5/2}$)

NSNS
BHNS

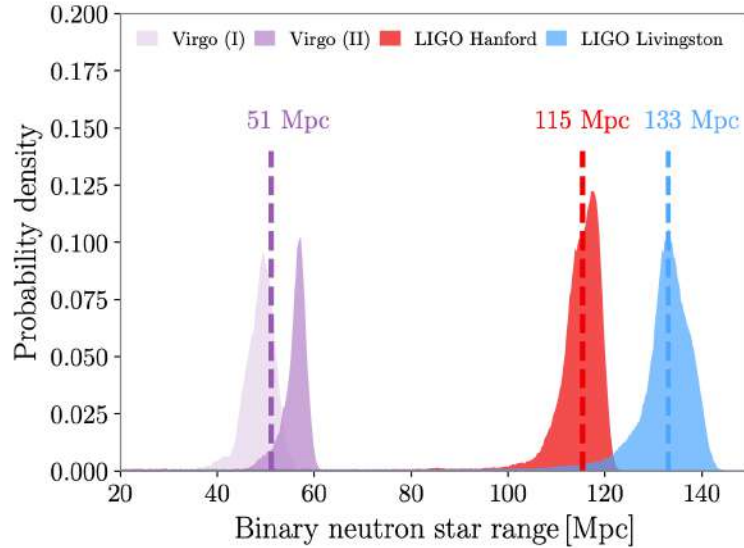
very massive

See GWTC-3
Population
analysis

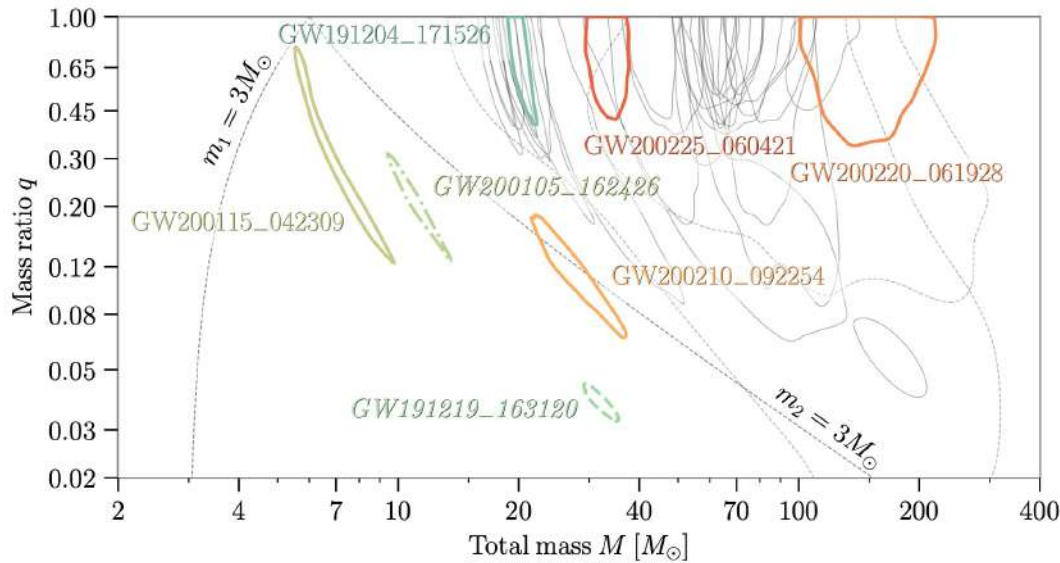
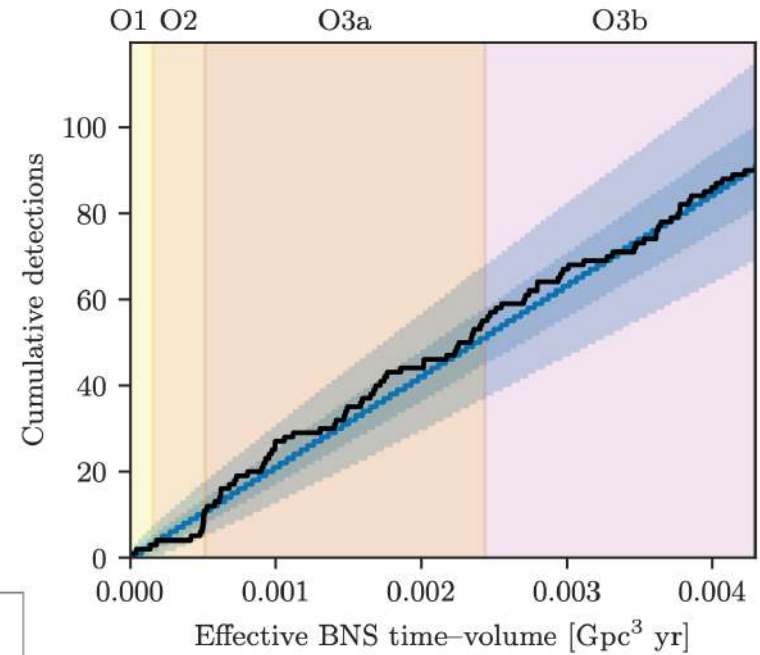
$q \sim 0.10$

Event	M (M_{\odot})	M (M_{\odot})	m_1 (M_{\odot})	m_2 (M_{\odot})	χ_{eff}	D_L (Gpc)	z	M_f (M_{\odot})	χ_f	$\Delta\Omega$ (deg)	SNR
GW190408.181802	42.9 ^{+4.1} _{-2.9}	18.3 ^{+1.8} _{-1.2}	24.5 ^{+5.1} _{-3.4}	18.3 ^{+3.2} _{-3.5}	-0.03 ^{+0.13} _{-0.19}	1.58 ^{+0.40} _{-0.29}	0.30 ^{+0.06} _{-0.10}	41.0 ^{+3.8} _{-2.7}	0.67 ^{+0.06} _{-0.07}	140	15.3 ^{+0.2} _{-0.3}
GW190412	38.4 ^{+3.8} _{-3.7}	13.3 ^{+0.4} _{-0.3}	30.0 ^{+5.1} _{-5.1}	8.3 ^{+1.0} _{-0.9}	0.25 ^{+0.05} _{-0.11}	0.74 ^{+0.14} _{-0.17}	0.15 ^{+0.03} _{-0.03}	37.3 ^{+3.9} _{-3.9}	0.67 ^{+0.05} _{-0.06}	21	18.9 ^{+0.2} _{-0.3}
GW190413.052954	56.9 ^{+13.1} _{-8.0}	24.0 ^{+5.4} _{-3.7}	33.4 ^{+12.4} _{-7.4}	23.4 ^{+5.7} _{-6.8}	0.01 ^{+0.29} _{-0.38}	4.10 ^{+2.41} _{-1.89}	0.66 ^{+0.30} _{-0.27}	54.3 ^{+12.4} _{-8.4}	0.69 ^{+0.12} _{-0.13}	1400	8.9 ^{+0.4} _{-0.8}
GW190413.134308	76.1 ^{+15.9} _{-10.6}	31.9 ^{+7.3} _{-4.6}	45.4 ^{+13.6} _{-9.6}	30.9 ^{+10.2} _{-9.0}	-0.01 ^{+0.24} _{-0.28}	5.15 ^{+2.44} _{-2.34}	0.80 ^{+0.30} _{-0.31}	72.8 ^{+15.2} _{-10.3}	0.69 ^{+0.10} _{-0.12}	520	10.0 ^{+0.4} _{-0.5}
GW190421.213856	71.8 ^{+12.3} _{-8.6}	30.7 ^{+5.5} _{-3.9}	40.6 ^{+10.4} _{-6.6}	31.4 ^{+7.5} _{-8.2}	-0.05 ^{+0.23} _{-0.20}	3.15 ^{+1.37} _{-1.42}	0.53 ^{+0.18} _{-0.21}	68.6 ^{+11.7} _{-8.1}	0.68 ^{+0.10} _{-0.11}	1000	10.7 ^{+0.2} _{-0.4}
GW190424.180648	70.7 ^{+13.4} _{-9.8}	30.3 ^{+5.7} _{-4.2}	39.3 ^{+10.0} _{-6.9}	31.8 ^{+7.4} _{-8.8}	0.15 ^{+0.22} _{-0.22}	2.55 ^{+1.56} _{-1.33}	0.45 ^{+0.22} _{-0.21}	67.1 ^{+12.5} _{-9.2}	0.75 ^{+0.08} _{-0.09}	26000	10.4 ^{+0.3} _{-0.4}
GW190425	3.4 ^{+0.3} _{-0.1}	1.44 ^{+0.05} _{-0.05}	2.0 ^{+0.0} _{-0.3}	1.4 ^{+0.1} _{-0.3}	0.06 ^{+0.11} _{-0.05}	0.16 ^{+0.07} _{-0.07}	0.03 ^{+0.02} _{-0.02}	-	-	9900	12.4 ^{+0.3} _{-0.4}
GW190426.152155	7.2 ^{+3.5} _{-1.5}	2.41 ^{+0.85} _{-0.6}	5.7 ^{+4.0} _{-2.3}	1.5 ^{+0.8} _{-0.5}	-0.03 ^{+0.33} _{-0.30}	0.38 ^{+0.19} _{-0.16}	0.08 ^{+0.04} _{-0.03}	-	-	1400	8.7 ^{+0.5} _{-0.4}
GW190503.185404	71.3 ^{+9.3} _{-8.0}	30.1 ^{+4.2} _{-4.0}	42.9 ^{+9.2} _{-7.8}	28.5 ^{+7.9} _{-7.9}	-0.02 ^{+0.20} _{-0.26}	1.52 ^{+0.71} _{-0.66}	0.29 ^{+0.11} _{-0.11}	68.2 ^{+8.7} _{-7.5}	0.67 ^{+0.09} _{-0.12}	94	12.4 ^{+0.2} _{-0.3}
GW190512.180714	35.6 ^{+3.9} _{-3.4}	14.5 ^{+1.3} _{-1.0}	23.0 ^{+5.4} _{-5.7}	12.5 ^{+3.5} _{-2.5}	0.03 ^{+0.13} _{-0.13}	1.49 ^{+0.53} _{-0.59}	0.28 ^{+0.08} _{-0.10}	34.2 ^{+3.9} _{-3.4}	0.65 ^{+0.07} _{-0.07}	230	12.2 ^{+0.2} _{-0.4}
GW190513.205428	53.6 ^{+8.8} _{-5.9}	21.5 ^{+3.6} _{-1.9}	35.3 ^{+9.6} _{-9.0}	18.1 ^{+7.3} _{-6.2}	0.12 ^{+0.29} _{-0.18}	2.16 ^{+0.94} _{-0.80}	0.39 ^{+0.14} _{-0.13}	51.3 ^{+8.1} _{-5.8}	0.69 ^{+0.14} _{-0.12}	490	12.9 ^{+0.3} _{-0.4}
GW190514.065416	64.2 ^{+16.9} _{-9.6}	27.4 ^{+6.9} _{-4.3}	36.9 ^{+13.4} _{-7.3}	27.5 ^{+8.2} _{-7.7}	-0.16 ^{+0.28} _{-0.32}	4.93 ^{+2.76} _{-2.41}	0.77 ^{+0.34} _{-0.33}	61.6 ^{+10.4} _{-9.2}	0.64 ^{+0.11} _{-0.14}	2400	8.2 ^{+0.3} _{-0.6}
GW190517.055101	61.9 ^{+10.0} _{-9.0}	26.0 ^{+4.2} _{-4.0}	36.4 ^{+11.9} _{-7.8}	24.8 ^{+6.9} _{-7.1}	0.53 ^{+0.20} _{-0.19}	2.11 ^{+1.79} _{-1.00}	0.38 ^{+0.26} _{-0.16}	57.8 ^{+9.4} _{-9.1}	0.87 ^{+0.05} _{-0.07}	460	10.7 ^{+0.4} _{-0.6}
GW190519.153544	104.2 ^{+24.5} _{-14.0}	43.5 ^{+6.8} _{-6.8}	64.5 ^{+11.3} _{-10.6}	39.9 ^{+11.0} _{-10.6}	0.33 ^{+0.19} _{-0.22}	2.85 ^{+2.02} _{-1.14}	0.49 ^{+0.27} _{-0.17}	98.7 ^{+13.5} _{-14.2}	0.80 ^{+0.07} _{-0.12}	770	15.6 ^{+0.2} _{-0.3}
GW190521	157.9 ^{+37.4} _{-20.4}	66.9 ^{+11.1} _{-9.6}	91.4 ^{+20.3} _{-17.5}	66.8 ^{+20.1} _{-20.7}	0.06 ^{+0.31} _{-0.37}	4.53 ^{+2.30} _{-2.13}	0.72 ^{+0.29} _{-0.29}	150.3 ^{+35.4} _{-20.4}	0.73 ^{+0.11} _{-0.14}	940	14.2 ^{+0.3} _{-0.3}
GW190521.074359	74.4 ^{+6.8} _{-4.6}	31.9 ^{+3.1} _{-2.4}	42.1 ^{+8.9} _{-7.8}	32.1 ^{+7.4} _{-6.2}	0.09 ^{+0.10} _{-0.13}	1.28 ^{+0.38} _{-0.57}	0.25 ^{+0.06} _{-0.10}	70.7 ^{+6.4} _{-4.2}	0.72 ^{+0.05} _{-0.07}	500	25.8 ^{+0.1} _{-0.2}
GW190527.092055	58.5 ^{+27.9} _{-10.9}	24.2 ^{+11.0} _{-4.4}	36.2 ^{+19.1} _{-9.5}	22.8 ^{+12.7} _{-8.1}	0.13 ^{+0.29} _{-0.28}	3.10 ^{+4.85} _{-1.84}	0.53 ^{+0.61} _{-0.25}	55.9 ^{+26.4} _{-10.1}	0.73 ^{+0.12} _{-0.16}	3800	8.1 ^{+0.4} _{-1.0}
GW190602.175927	114.1 ^{+18.5} _{-13.7}	48.3 ^{+8.6} _{-8.0}	67.2 ^{+16.0} _{-12.6}	47.4 ^{+13.4} _{-16.6}	0.10 ^{+0.25} _{-0.25}	2.99 ^{+2.02} _{-1.26}	0.51 ^{+0.27} _{-0.19}	108.8 ^{+17.2} _{-14.2}	0.71 ^{+0.10} _{-0.13}	720	12.8 ^{+0.2} _{-0.4}
GW190620.030421	90.1 ^{+17.3} _{-12.1}	37.5 ^{+5.7} _{-5.7}	55.4 ^{+15.8} _{-12.0}	35.0 ^{+11.6} _{-11.4}	0.34 ^{+0.21} _{-0.25}	3.16 ^{+1.67} _{-1.43}	0.54 ^{+0.22} _{-0.21}	85.4 ^{+15.9} _{-11.4}	0.80 ^{+0.08} _{-0.14}	6700	12.1 ^{+0.3} _{-0.4}
GW190630.185205	58.8 ^{+4.7} _{-4.8}	24.8 ^{+2.1} _{-2.0}	35.0 ^{+6.9} _{-5.7}	23.6 ^{+5.2} _{-5.1}	0.10 ^{+0.12} _{-0.13}	0.93 ^{+0.56} _{-0.40}	0.19 ^{+0.10} _{-0.07}	56.1 ^{+4.5} _{-4.0}	0.70 ^{+0.06} _{-0.07}	1300	15.6 ^{+0.2} _{-0.3}
GW190701.203306	94.1 ^{+11.6} _{-9.3}	40.2 ^{+5.2} _{-4.7}	53.6 ^{+11.7} _{-7.8}	40.8 ^{+8.3} _{-11.5}	-0.06 ^{+0.23} _{-0.28}	2.14 ^{+0.79} _{-0.73}	0.38 ^{+0.12} _{-0.12}	90.0 ^{+10.8} _{-8.6}	0.67 ^{+0.09} _{-0.12}	45	11.3 ^{+0.2} _{-0.2}
GW190706.222641	101.6 ^{+17.9} _{-13.5}	42.0 ^{+8.4} _{-6.2}	64.0 ^{+15.2} _{-11.2}	38.5 ^{+12.5} _{-12.4}	0.32 ^{+0.25} _{-0.30}	5.07 ^{+2.57} _{-2.11}	0.79 ^{+0.31} _{-0.28}	96.3 ^{+18.7} _{-13.2}	0.80 ^{+0.08} _{-0.17}	610	12.6 ^{+0.2} _{-0.4}
GW190707.093326	20.0 ^{+1.9} _{-1.3}	8.5 ^{+0.6} _{-0.4}	11.5 ^{+3.3} _{-1.7}	8.4 ^{+1.4} _{-1.9}	-0.05 ^{+0.10} _{-0.04}	0.80 ^{+0.37} _{-0.38}	0.16 ^{+0.07} _{-0.07}	19.2 ^{+1.9} _{-1.3}	0.66 ^{+0.03} _{-0.04}	1300	13.3 ^{+0.2} _{-0.2}
GW190708.232457	30.8 ^{+2.5} _{-1.8}	13.1 ^{+0.9} _{-0.6}	17.5 ^{+4.7} _{-2.3}	13.1 ^{+2.0} _{-2.7}	0.02 ^{+0.10} _{-0.08}	0.90 ^{+0.33} _{-0.40}	0.18 ^{+0.06} _{-0.07}	29.4 ^{+2.5} _{-1.7}	0.69 ^{+0.04} _{-0.04}	14000	13.1 ^{+0.2} _{-0.3}
GW190719.215514	55.8 ^{+16.3} _{-10.0}	22.7 ^{+5.9} _{-3.7}	35.2 ^{+16.9} _{-9.9}	20.2 ^{+8.1} _{-6.5}	0.35 ^{+0.28} _{-0.32}	4.61 ^{+2.84} _{-2.17}	0.73 ^{+0.35} _{-0.30}	52.9 ^{+15.5} _{-9.5}	0.80 ^{+0.10} _{-0.16}	2300	8.3 ^{+0.3} _{-1.0}
GW190720.000836	21.3 ^{+4.3} _{-2.3}	8.9 ^{+0.5} _{-0.8}	13.3 ^{+6.0} _{-3.0}	7.8 ^{+2.2} _{-2.2}	0.18 ^{+0.14} _{-0.12}	0.81 ^{+0.71} _{-0.33}	0.16 ^{+0.12} _{-0.06}	20.3 ^{+4.5} _{-2.3}	0.72 ^{+0.06} _{-0.05}	510	11.0 ^{+0.3} _{-0.8}
GW190727.060333	65.8 ^{+10.9} _{-7.4}	28.1 ^{+4.9} _{-3.4}	37.2 ^{+9.4} _{-5.9}	28.8 ^{+6.6} _{-7.9}	0.12 ^{+0.26} _{-0.25}	3.60 ^{+1.56} _{-1.51}	0.60 ^{+0.20} _{-0.22}	62.6 ^{+10.2} _{-7.0}	0.73 ^{+0.10} _{-0.10}	860	11.9 ^{+0.3} _{-0.5}
GW190728.064510	20.5 ^{+4.3} _{-1.3}	8.6 ^{+0.3} _{-0.3}	12.2 ^{+7.1} _{-2.2}	8.1 ^{+1.7} _{-2.6}	0.12 ^{+0.19} _{-0.07}	0.89 ^{+0.25} _{-0.37}	0.18 ^{+0.05} _{-0.07}	19.5 ^{+4.0} _{-1.3}	0.71 ^{+0.04} _{-0.04}	410	13.0 ^{+0.2} _{-0.4}
GW190731.140936	67.1 ^{+13.3} _{-10.2}	28.4 ^{+6.8} _{-4.5}	39.3 ^{+11.8} _{-8.2}	28.0 ^{+8.9} _{-8.4}	0.08 ^{+0.24} _{-0.24}	3.97 ^{+2.56} _{-2.07}	0.65 ^{+0.32} _{-0.30}	63.9 ^{+14.4} _{-9.8}	0.71 ^{+0.10} _{-0.12}	3000	8.6 ^{+0.2} _{-0.5}
GW190803.022701	62.7 ^{+11.8} _{-8.4}	26.7 ^{+5.2} _{-3.8}	36.1 ^{+10.2} _{-7.6}	26.7 ^{+7.1} _{-7.6}	-0.01 ^{+0.25} _{-0.26}	3.69 ^{+2.04} _{-1.69}	0.61 ^{+0.24} _{-0.24}	59.9 ^{+11.3} _{-7.9}	0.69 ^{+0.11} _{-0.11}	1500	8.6 ^{+0.3} _{-0.5}
GW190814	25.8 ^{+1.0} _{-0.9}	6.09 ^{+0.36} _{-0.36}	23.2 ^{+1.1} _{-1.0}	2.59 ^{+0.06} _{-0.09}	0.00 ^{+0.06} _{-0.06}	0.24 ^{+0.04} _{-0.05}	0.05 ^{+0.00} _{-0.010}	25.6 ^{+1.0} _{-0.9}	0.28 ^{+0.02} _{-0.02}	19	24.9 ^{+0.1} _{-0.2}
GW190828.063405	57.5 ^{+7.5} _{-4.4}	24.8 ^{+3.3} _{-2.0}	31.8 ^{+8.8} _{-5.9}	23.9 ^{+4.6} _{-4.6}	0.19 ^{+0.15} _{-0.18}	2.22 ^{+0.63} _{-0.95}	0.40 ^{+0.09} _{-0.15}	54.5 ^{+6.9} _{-4.0}	0.76 ^{+0.07} _{-0.07}	520	16.2 ^{+0.2} _{-0.3}
GW190828.065509	34.1 ^{+5.5} _{-4.5}	13.3 ^{+1.2} _{-0.9}	23.8 ^{+7.2} _{-7.0}	10.2 ^{+3.5} _{-2.1}	0.08 ^{+0.16} _{-0.18}	1.66 ^{+0.83} _{-0.61}	0.31 ^{+0.10} _{-0.10}	32.9 ^{+4.5} _{-4.5}	0.65 ^{+0.09} _{-0.08}	640	10.0 ^{+0.3} _{-0.5}
GW190909.114149	71.2 ^{+14.3} _{-15.0}	29.5 ^{+7.5} _{-6.3}	43.2 ^{+10.7} _{-12.2}	27.6 ^{+13.0} _{-10.9}	-0.03 ^{+0.44} _{-0.36}	4.77 ^{+3.70} _{-2.60}	0.75 ^{+0.45} _{-0.37}	68.3 ^{+14.5} _{-14.5}	0.68 ^{+0.16} _{-0.18}	4200	8.1 ^{+0.4} _{-0.7}
GW190910.112807	78.7 ^{+9.5} _{-9.0}	33.9 ^{+4.3} _{-3.9}	43.5 ^{+7.6} _{-6.2}	35.1 ^{+6.3} _{-7.0}	0.02 ^{+0.19} _{-0.18}	1.57 ^{+1.07} _{-0.84}	0.29 ^{+0.17} _{-0.11}	75.0 ^{+8.7} _{-8.5}	0.70 ^{+0.08} _{-0.07}	10000	14.1 ^{+0.2} _{-0.3}
GW190915.235702	59.5 ^{+7.5} _{-6.2}	25.1 ^{+3.1} _{-2.6}	34.9 ^{+8.5} _{-6.2}	24.4 ^{+5.5} _{-6.0}	0.03 ^{+0.19} _{-0.24}	1.70 ^{+0.71} _{-0.64}	0.32 ^{+0.11} _{-0.11}	56.8 ^{+7.1} _{-5.8}	0.71 ^{+0.09} _{-0.11}	380	13.6 ^{+0.2} _{-0.3}
GW190924.021846	13.9 ^{+5.1} _{-0.9}	5.8 ^{+0.2} _{-0.2}	8.8 ^{+7.0} _{-2.0}	5.0 ^{+1.3} _{-1.8}	0.03 ^{+0.30} _{-0.09}	0.57 ^{+0.22} _{-0.22}	0.12 ^{+0.04} _{-0.04}	13.3 ^{+5.2} _{-1.0}	0.67 ^{+0.05} _{-0.05}	380	11.5 ^{+0.3} _{-0.4}
GW190929.012149	90.6 ^{+21.2} _{-14.1}	34.3 ^{+8.6} _{-6.5}	64.7 ^{+22.4} _{-18.9}	25.7 ^{+14.4} _{-9.7}	0.03 ^{+0.27} _{-0.27}	3.68 ^{+2.98} _{-1.68}	0.61 ^{+0.38} _{-0.24}	87.5 ^{+20.7} _{-14.1}	0.64 ^{+0.17} _{-0.23}	1800	9.8 ^{+0.8} _{-0.6}
GW190930.193841	20.3 ^{+3.0} _{-1.1}	8.9 ^{+0.3} _{-0.5}	12.3 ^{+2.5} _{-2.3}	7.8 ^{+1.7} _{-1.3}	0.14 ^{+0.31} _{-0.15}	0.78 ^{+0.37} _{-0.33}	0.16 ^{+0.07} _{-0.06}	19.3 ^{+3.3} _{-1.5}	0.72 ^{+0.07} _{-0.06}	1800	9.5 ^{+0.5} _{-0.5}

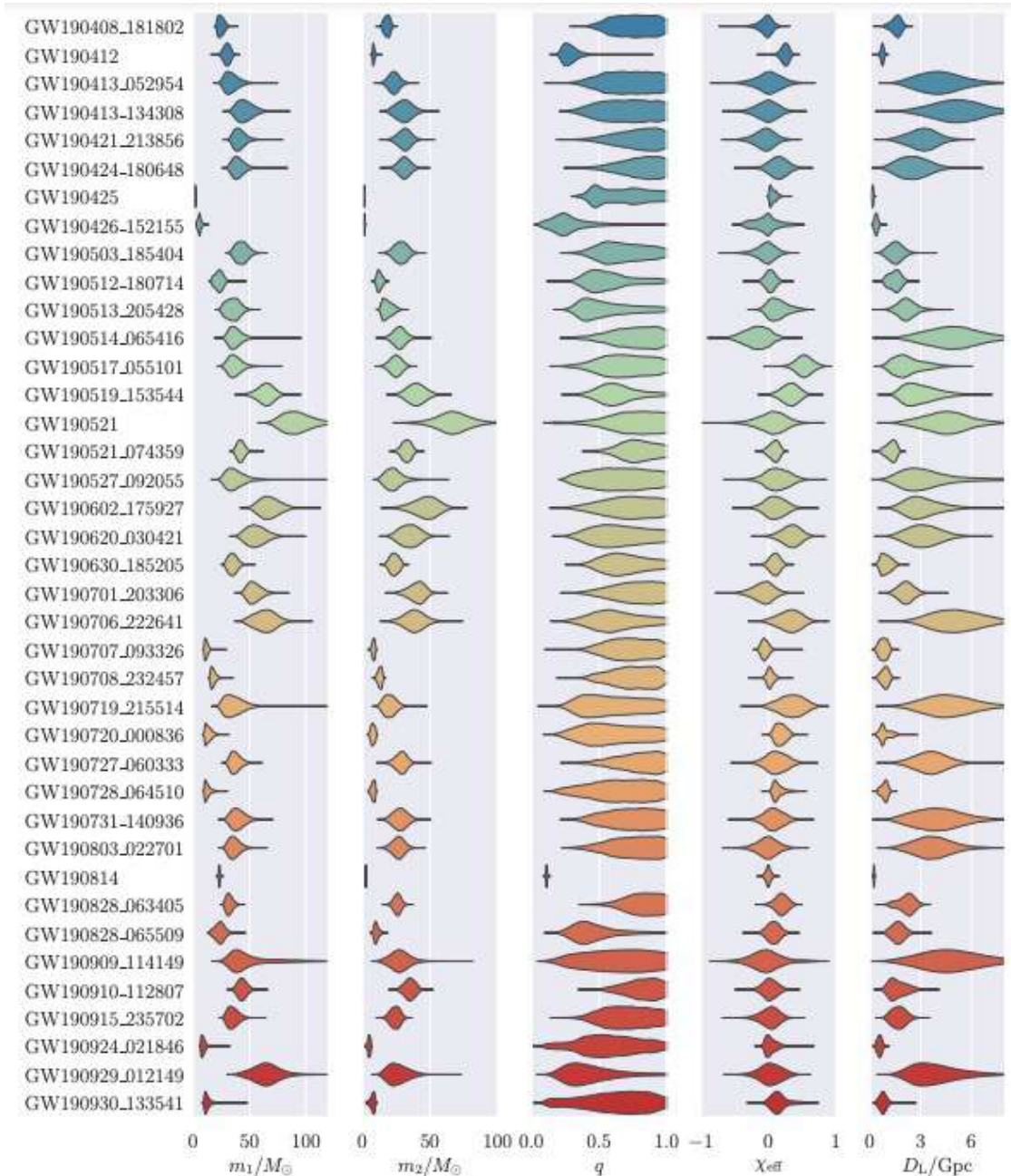
Detection range (NSNS)



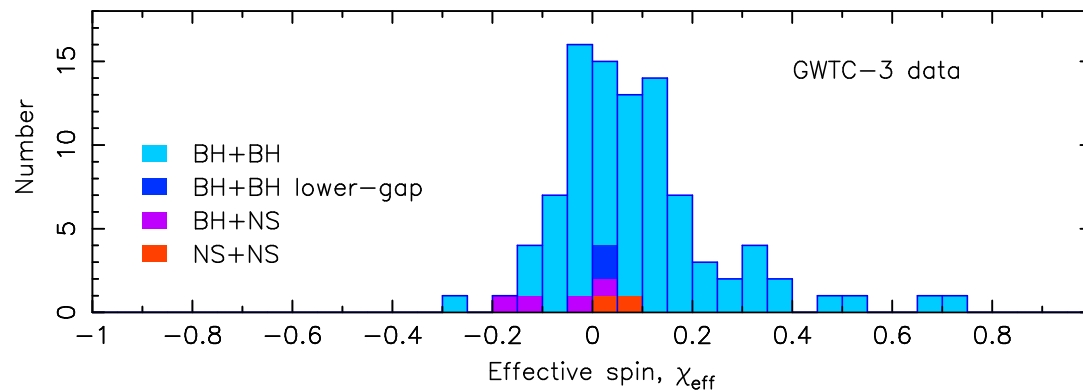
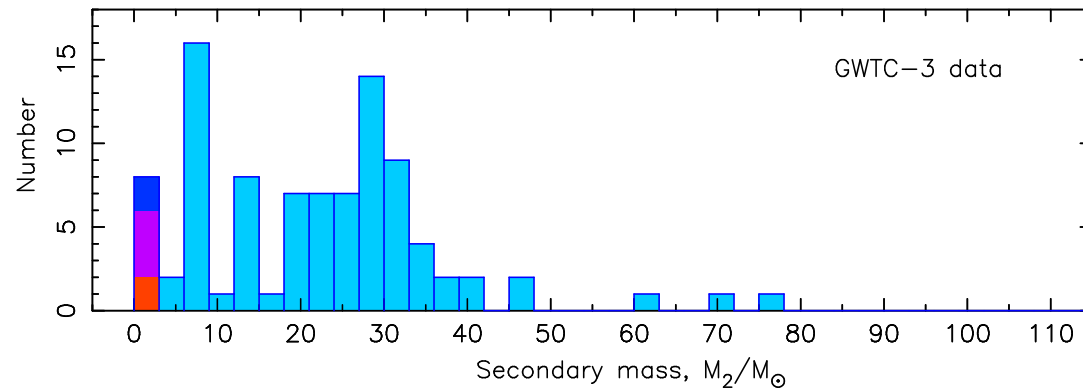
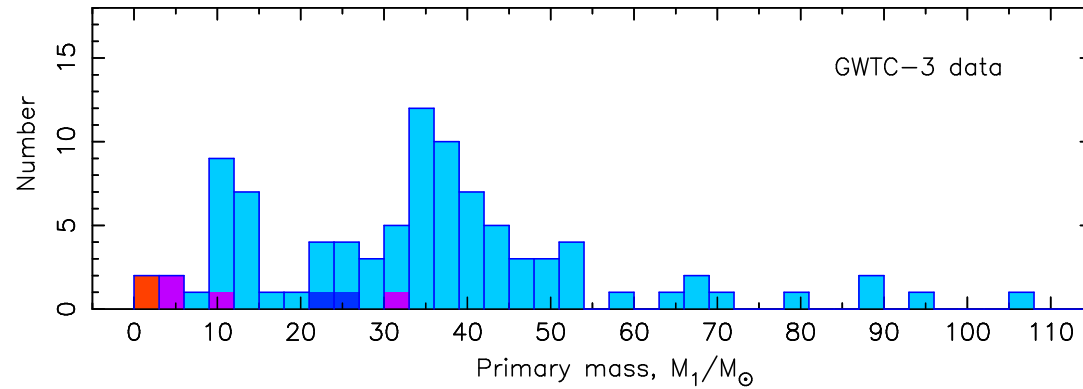
Detection versus time-volume



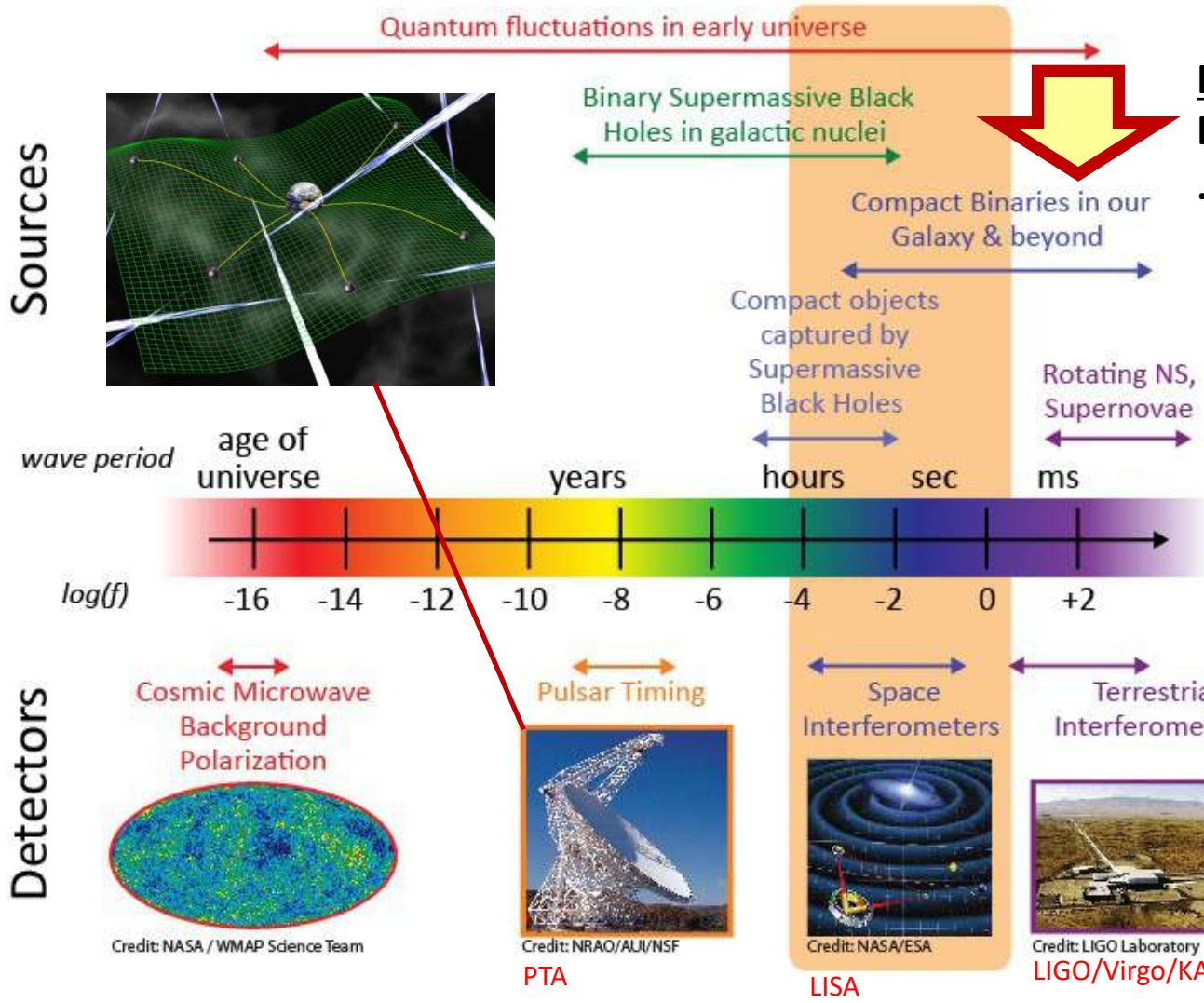
- Component masses
- Mass ratios
- Spins
- Distances



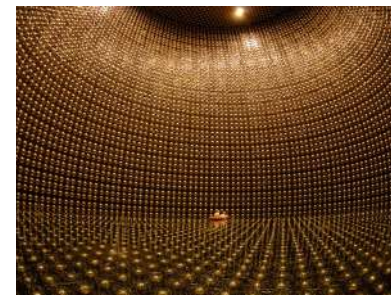
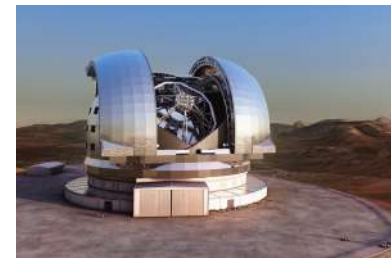
Tauris & van den Heuvel (2023)



The Gravitational Wave Spectrum © NASA



Follow-up:
EM transient counterparts
..... and neutrinos



high freq. GWs **LIGO: 10 Hz – 1 kHz**

low freq. GWs **LISA: 0.1 mHz – 10 mHz**

$$* \Delta E_{GW} < 10^{-8} M_{\odot} c^2$$

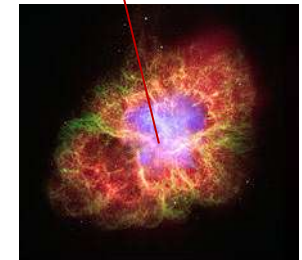
I. Transient (one-time) burst events: extragalactic

LIGO ❖ Colliding neutron star + black hole binaries

(**LISA** may detect these mergers too)

LIGO ❖ Supernova core collapse (**Galactic!**) *

LISA ❖ Supermassive black hole mergers



II. Persistent sources (continuous emission): Galactic

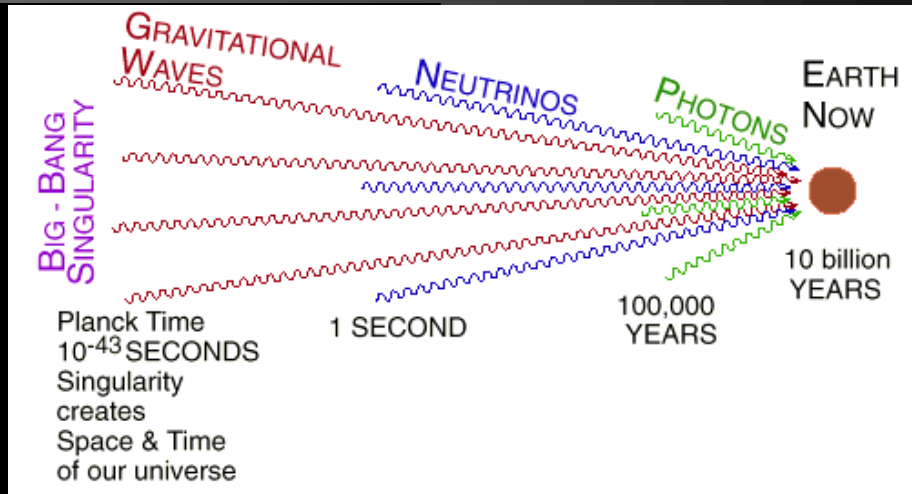
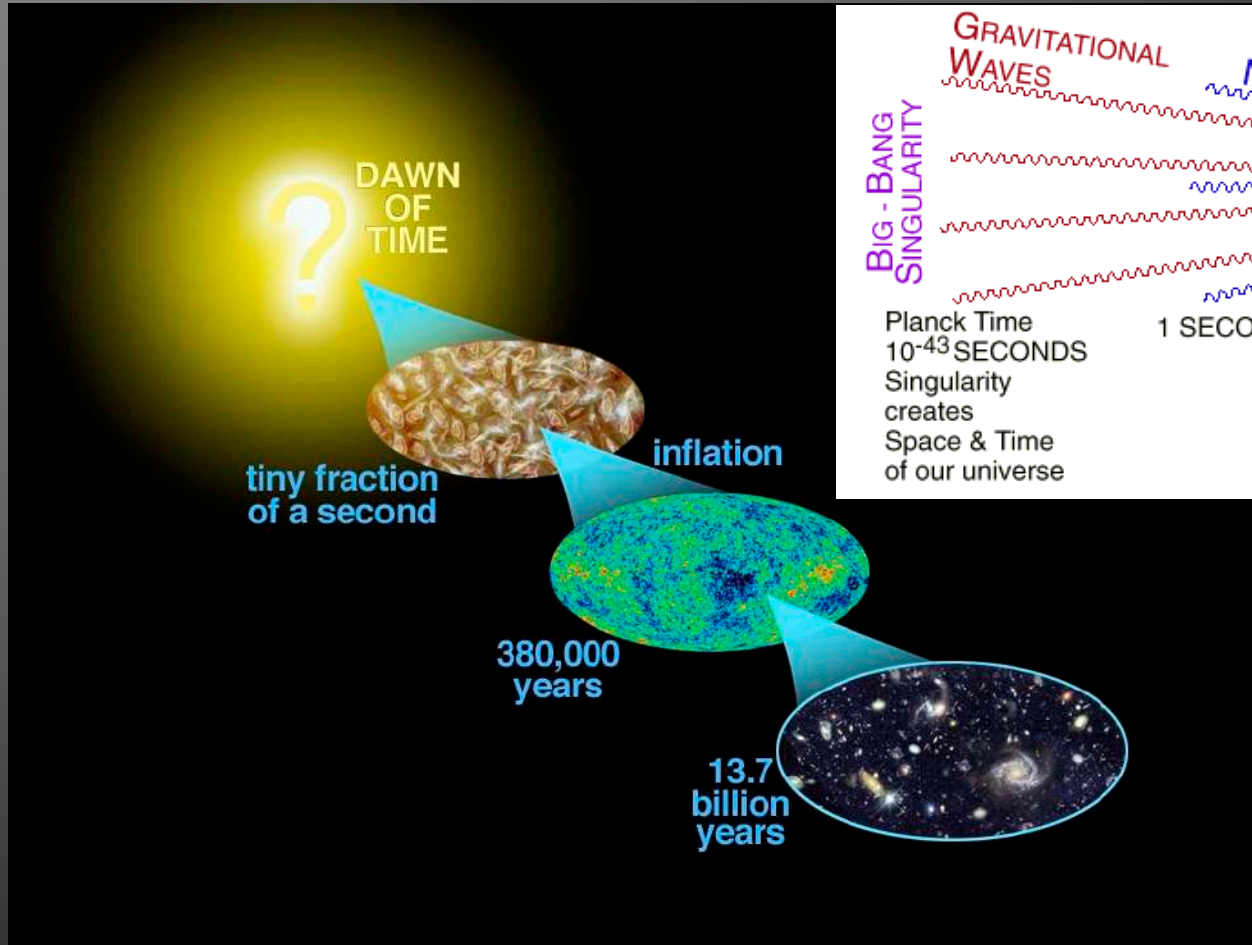
LIGO ❖ Pulsars or accreting NS

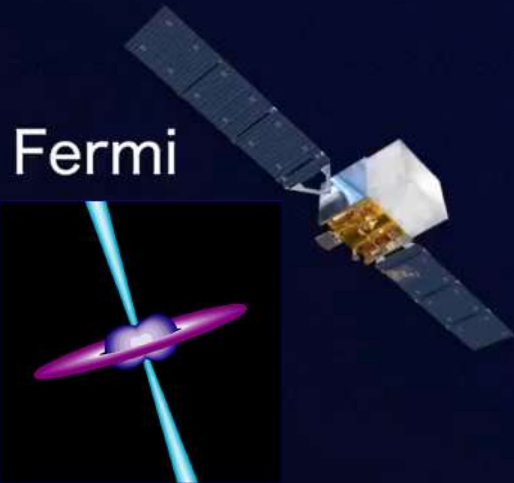
LISA ❖ Galactic resolved compact binaries (WD, NS, BH)



'Murmurs' from the Big Bang

signals from the early Universe

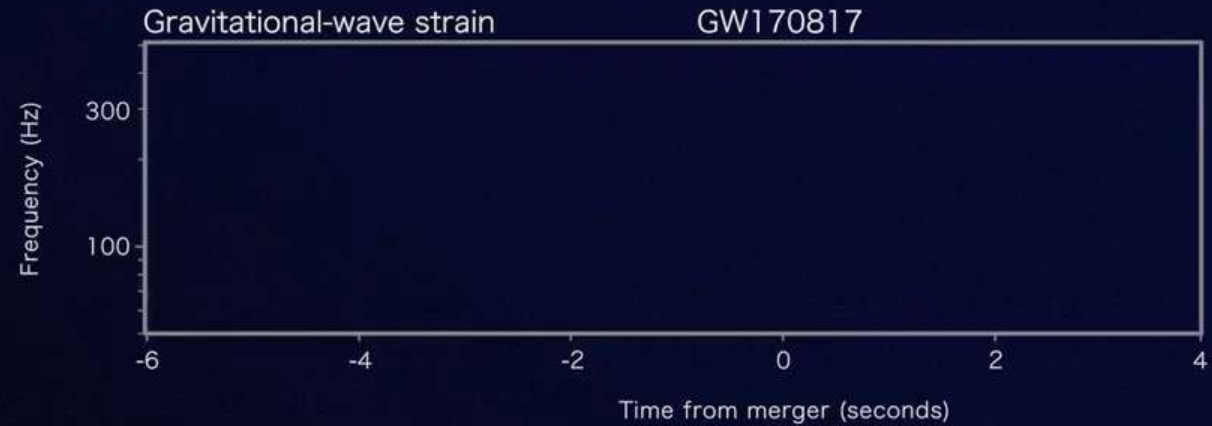




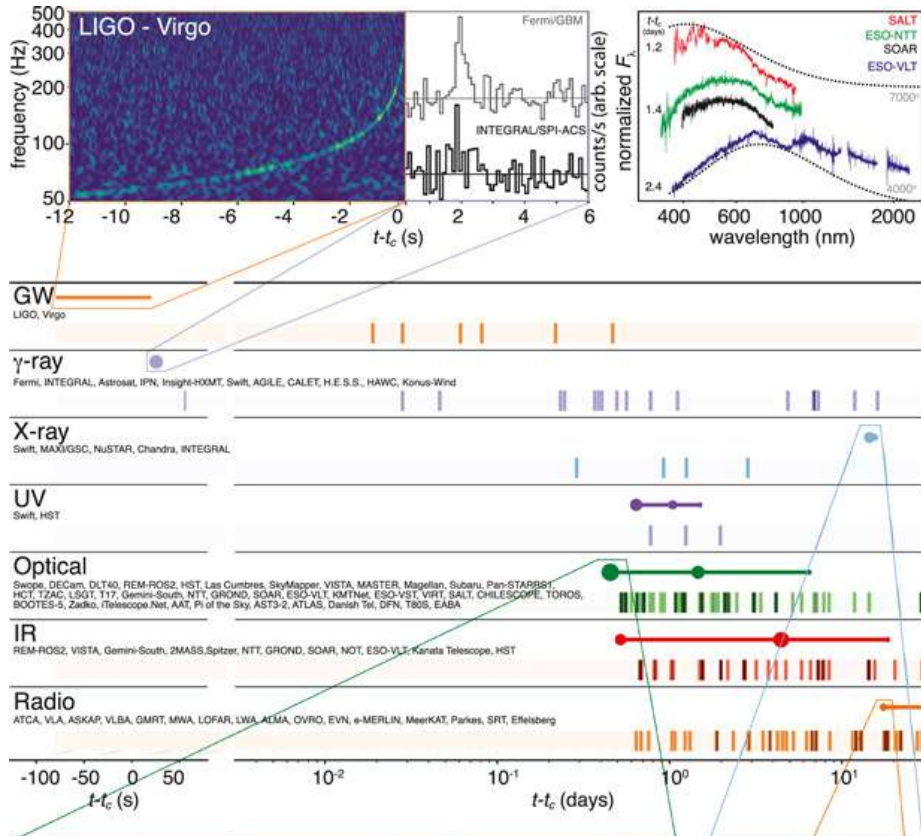
Fermi



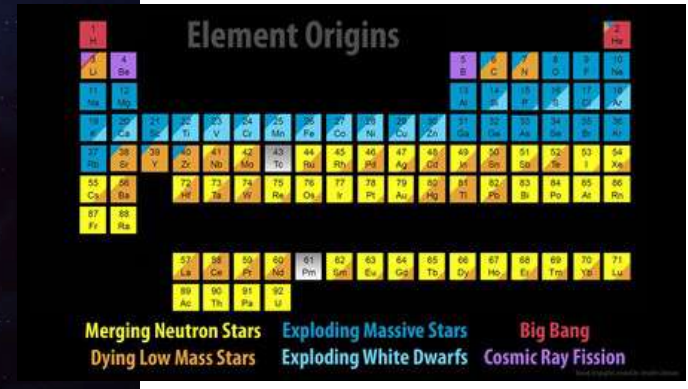
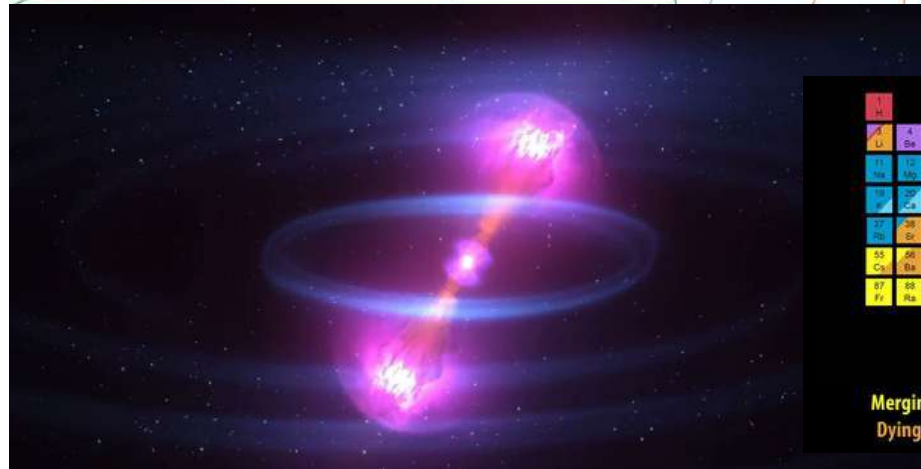
LIGO

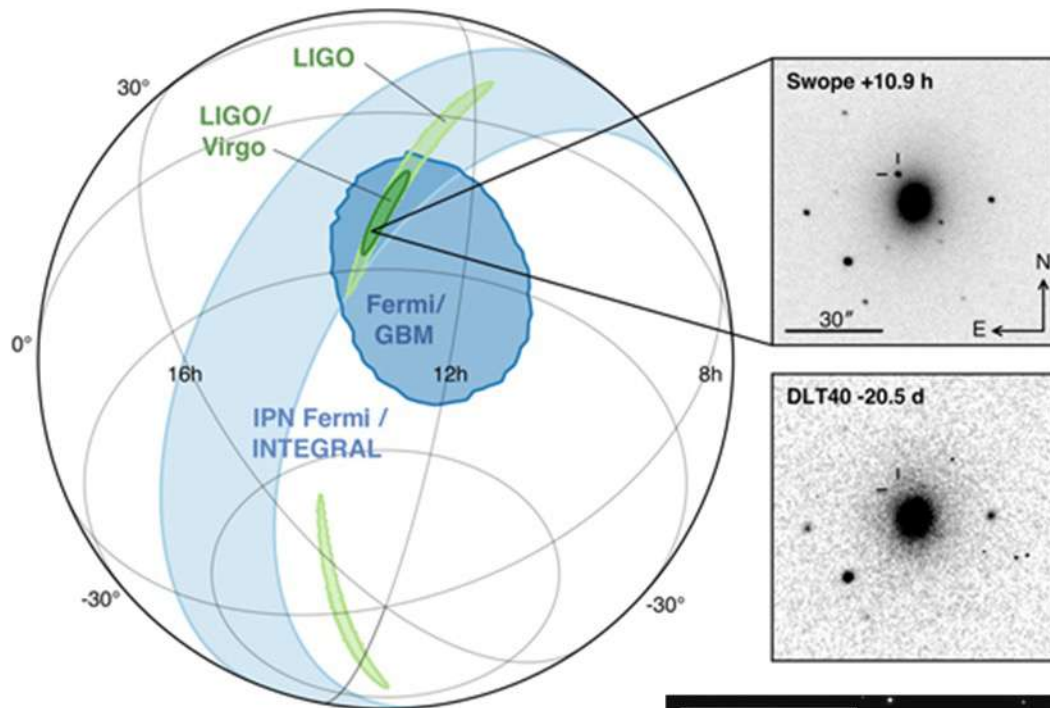


New era of **multi-messenger** astrophysics

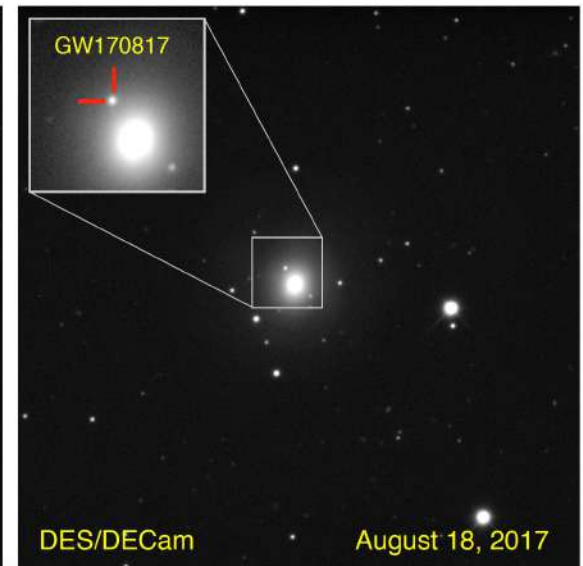
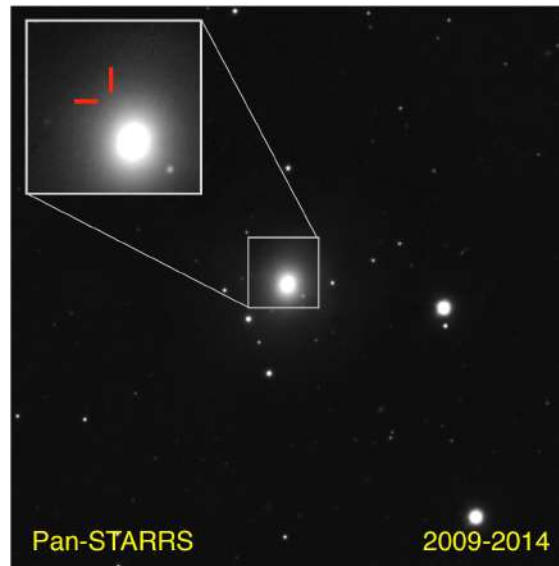


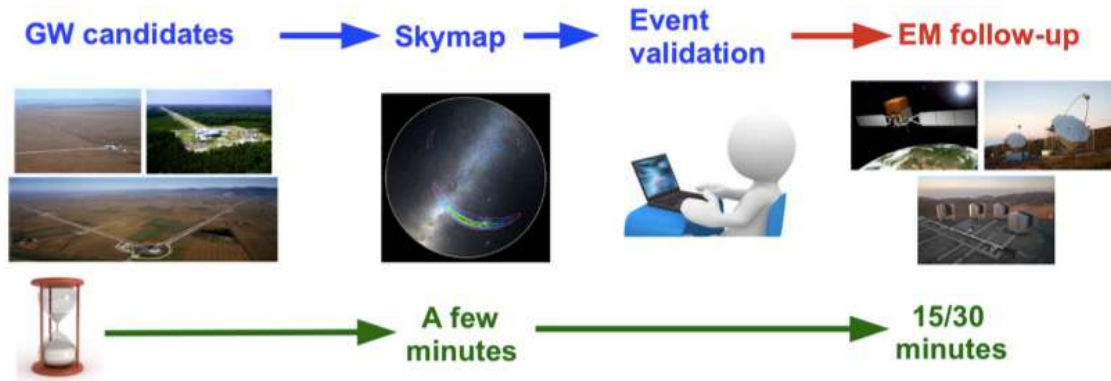
Detection of a **kilonova** (radioactive decay of heavy r-process elements)





NGC 4993
 d = 40 Mpc ($z=0.009$, 130 mill. ly)





© Patricelli

GW

- mass
- spin
- eccentricity
- luminosity distance
- system orientation
- BH/NS merger-rate density
- evolution over cosmic time (primordial BHs, SMBH seeds)

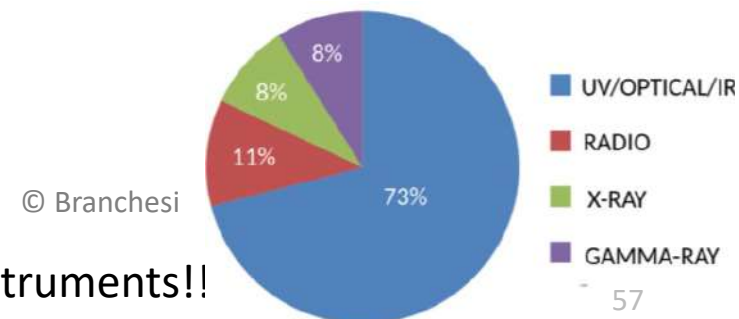
EM

- sky location
- host galaxy
- redshift
- local environment
- heavy r-process nucleosynthesis
- emission processes (kilonova)
- sGRB (central engine, beaming, jets, afterglow)

NEUTRINOS

- neutrino physics
- central engine
- SN explosions

- **Testing theories of gravity**
- **NS equation-of-state**
- **Cosmology**

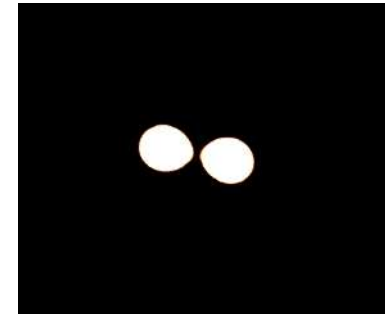
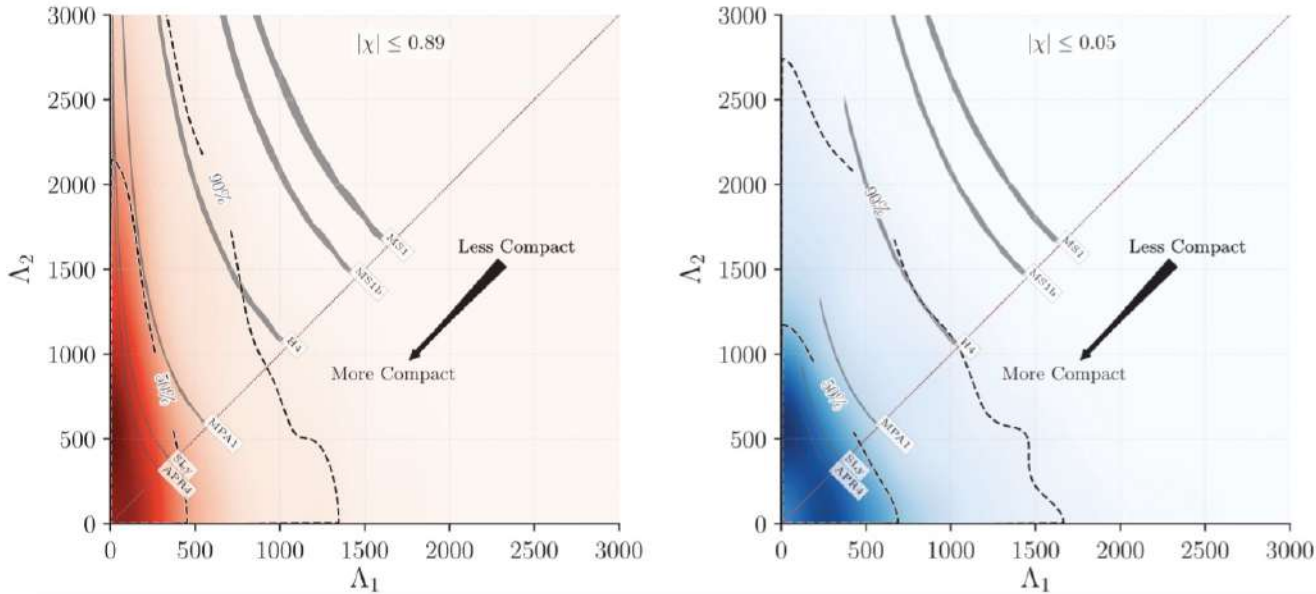


© Branchesi

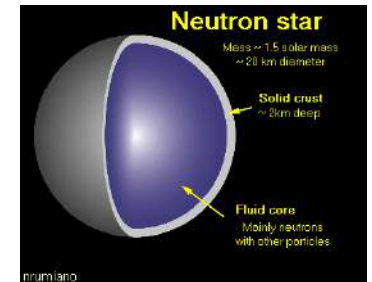
80 MoUs involving 170 instruments!!

Tidal deformation and NS EoS

L 119, 161101 (2017) PHYSICAL REVIEW LETTERS week ending 20 OCTOBER 2017



S. Roswogg



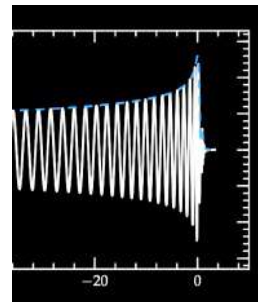
GW phase depends on tidal deformability parameter:

$$\lambda \equiv -\frac{Q_{ij}}{\mathcal{E}_{ij}} \quad \Lambda \equiv \frac{\lambda}{m^5} = \frac{2}{3} k_2 \left(\frac{Gm}{Rc^2} \right)^{-5}$$

$$\tilde{\Lambda} = \frac{16}{13} \frac{(12q + 1)\Lambda_1 + (12 + q)q^4\Lambda_2}{(1 + q)^5}$$

$$q = m_2/m_1 \leq 1$$

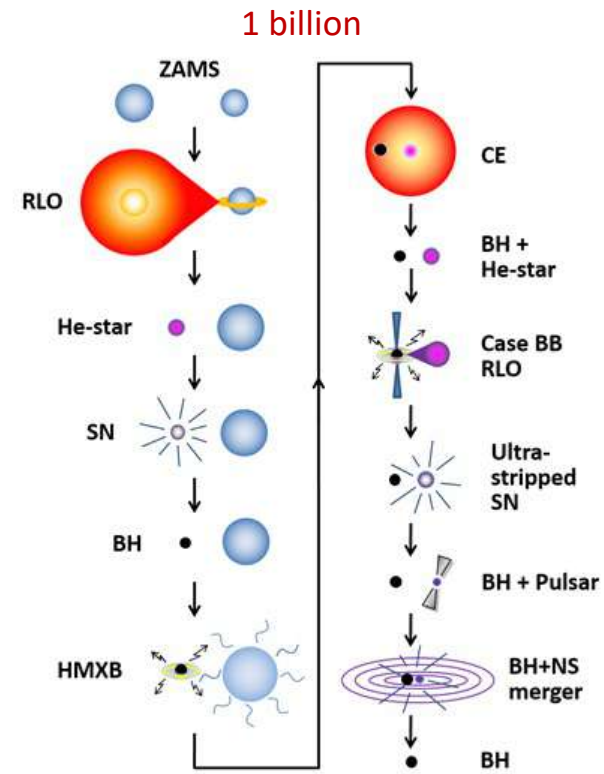
$$\Lambda \propto k(R/M)^5$$



Flanagan & Hinderer (2008)

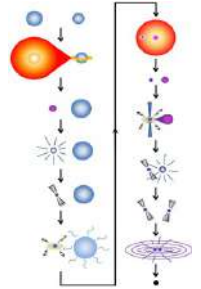
RECIPE

- Binary stellar evolution
- Population synthesis
(input distributions and stellar grids)
- Galactic star formation rate
(formation history of massive binaries)
- Galactic potentials
(to probe location of mergers in host galaxies)
- Extrapolation to local Universe
(scaling-law of galaxy number density)



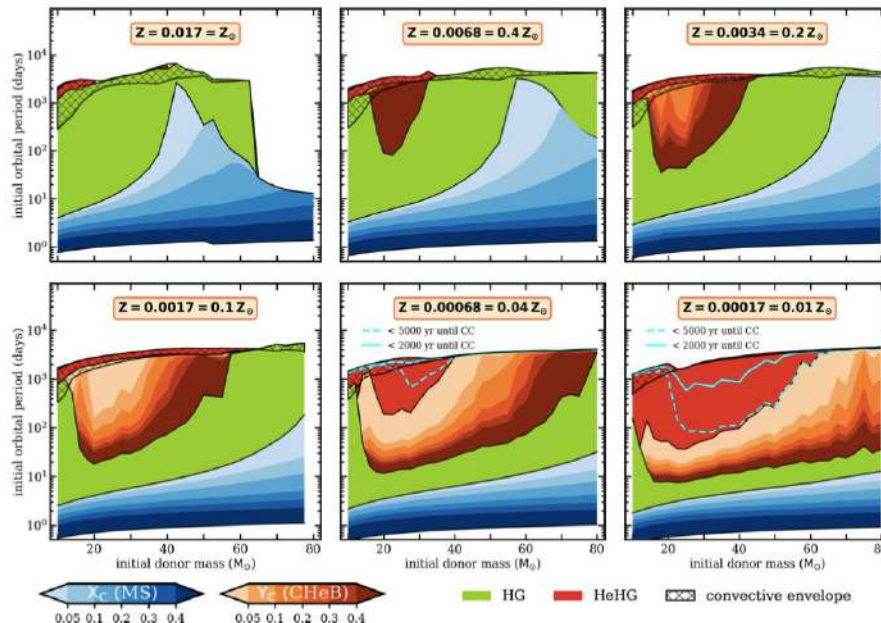
Notes on

POP. SYNTHESIS

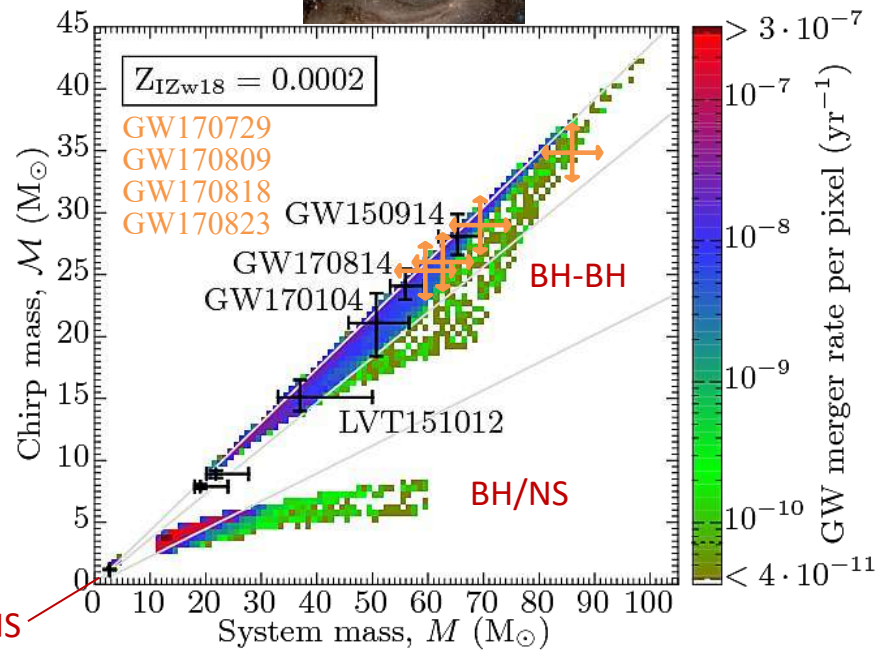
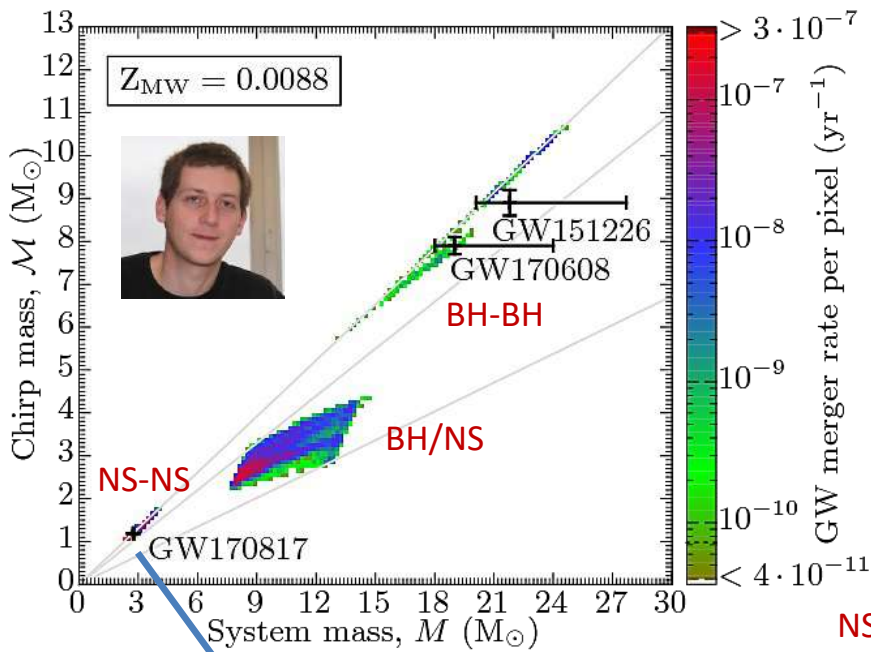


1. Reproduction of LIGO rates is *no* success criterion on its own
2. Can Galactic sources be reproduced? (properties of HMXBs, DNSs, etc.)
3. Is the input physics reasonable? Is the evolution self consistent? *
4. Watch out for papers that claim they can explain everything!

* E.g. Stellar evolution (Z)
Klencki et al. (2020)



Kruckow et al. (2018), MNRAS

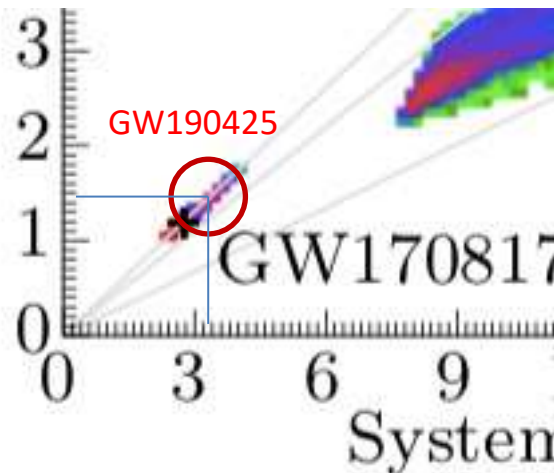


GW190425

$M_{\text{chirp}} = 1.44 M_{\text{sun}}$

$M_{\text{total}} = 3.3 M_{\text{sun}}$

$(\chi < 0.05)$



GW170817:

$R = 1540 \pm \begin{matrix} 3200 \\ 1220 \end{matrix} \text{ Gpc}^{-3} \text{ yr}^{-1}$

GW190425:

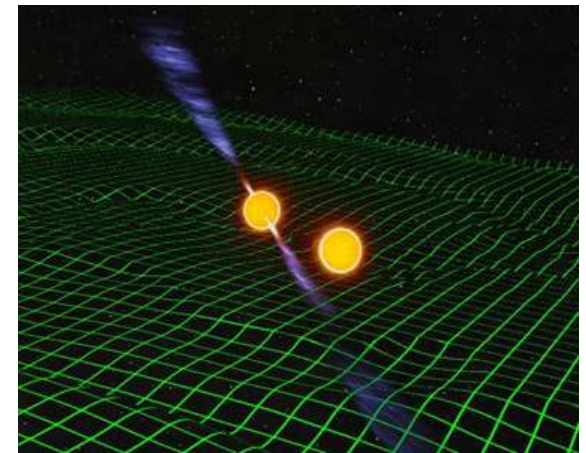
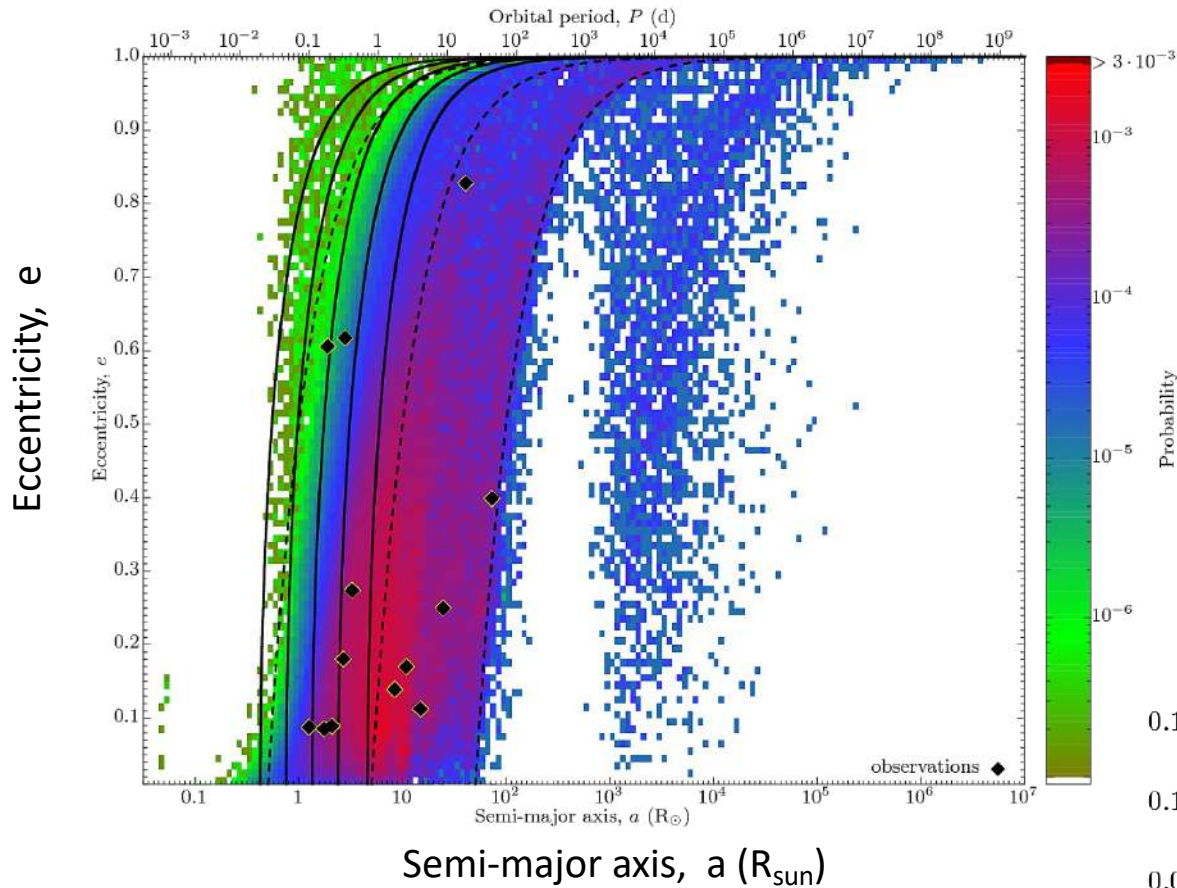
$R = 1090 \pm \begin{matrix} 1720 \\ 800 \end{matrix} \text{ Gpc}^{-3} \text{ yr}^{-1}$

GWTC-2 rate:

$R = 320 \pm \begin{matrix} 490 \\ 240 \end{matrix} \text{ Gpc}^{-3} \text{ yr}^{-1}$

GWTC-3 rate:

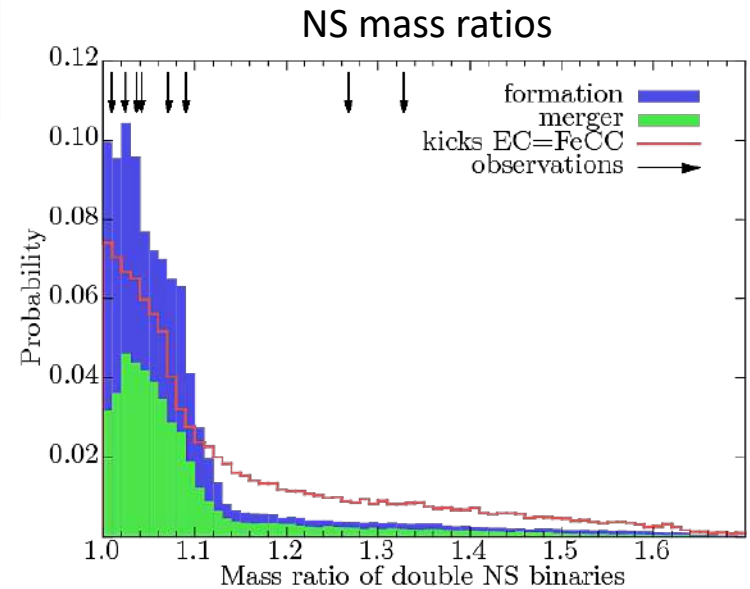
$R = 99 \pm \begin{matrix} 260 \\ 86 \end{matrix} \text{ Gpc}^{-3} \text{ yr}^{-1}$

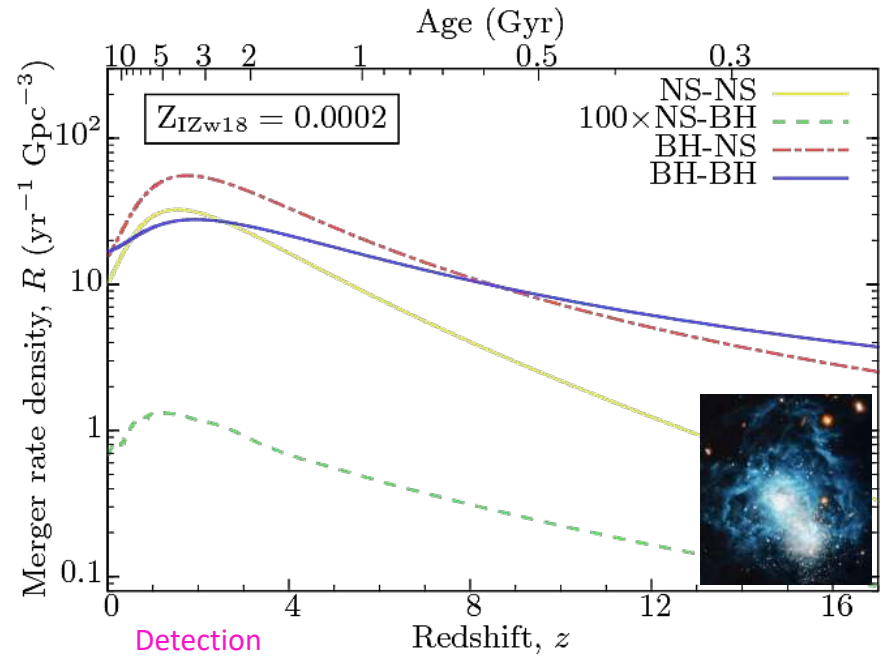
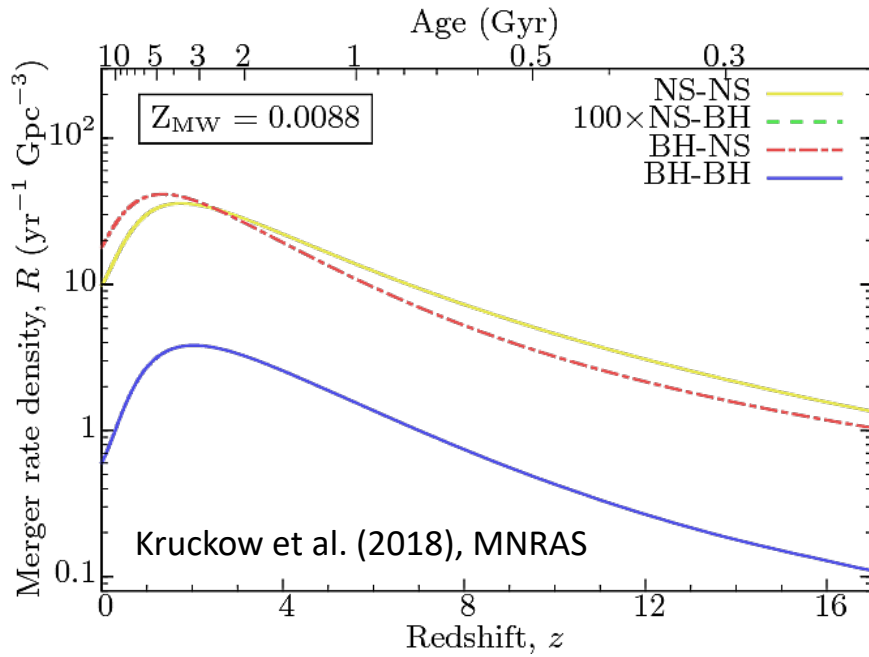


Do not trust population synthesis if it cannot reproduce observed Galactic DNS systems

Kruckow, Tauris, et al. (2018), MNRAS

Important calibration data!!!



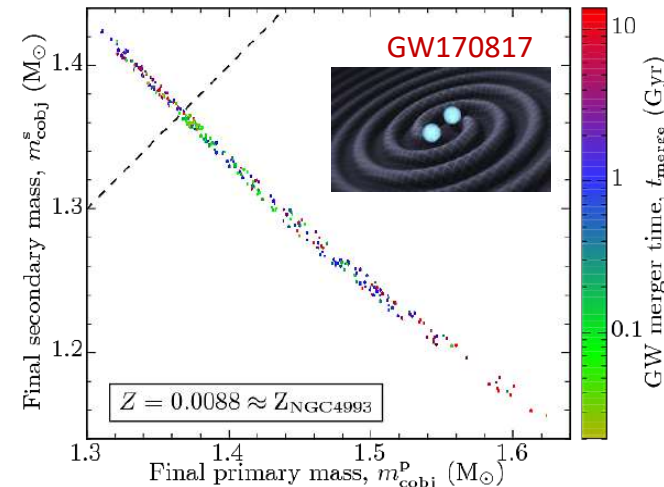


local Universe

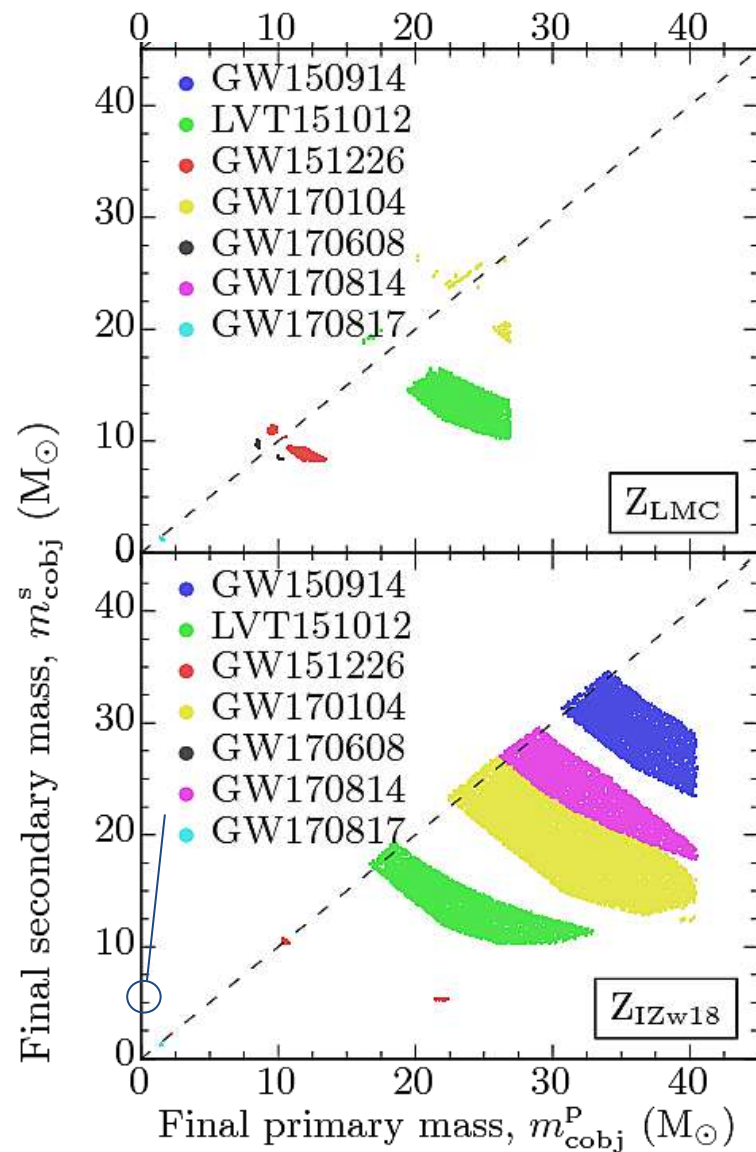
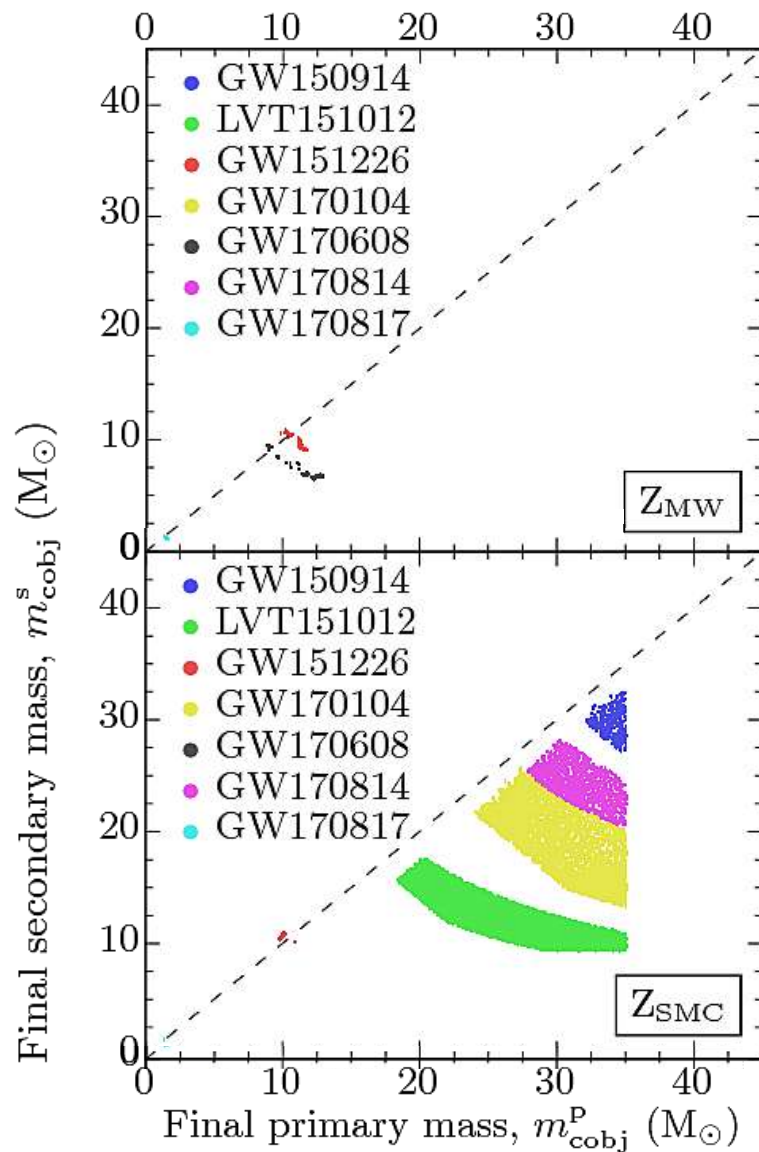
Detection rates

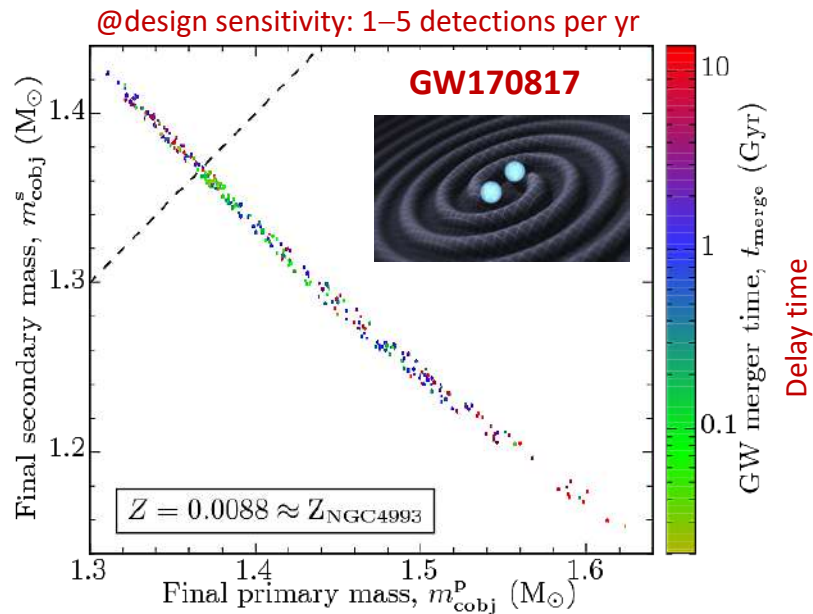
Z_{MW}	$\langle M^{2.5} \rangle$	$R_{z=0}$	R_D	R_{cSFR}	$R_{D,cSFR}$
NS-NS	$1.36 M_{\odot}^{2.5}$	$9.85 \times 10^0 \text{ yr}^{-1} \text{ Gpc}^{-3}$	0.28 yr^{-1}	$3.47 \times 10^1 \text{ yr}^{-1} \text{ Gpc}^{-3}$	0.98 yr^{-1}
NS-BH	$20.0 M_{\odot}^{2.5}$	$0.00 \times 10^0 \text{ yr}^{-1} \text{ Gpc}^{-3}$	0.00 yr^{-1}	$0.00 \times 10^0 \text{ yr}^{-1} \text{ Gpc}^{-3}$	0.00 yr^{-1}
BH-NS	$15.7 M_{\odot}^{2.5}$	$1.80 \times 10^1 \text{ yr}^{-1} \text{ Gpc}^{-3}$	5.88 yr^{-1}	$4.72 \times 10^1 \text{ yr}^{-1} \text{ Gpc}^{-3}$	15.43 yr^{-1}
BH-BH	$233 M_{\odot}^{2.5}$	$6.01 \times 10^{-1} \text{ yr}^{-1} \text{ Gpc}^{-3}$	2.92 yr^{-1}	$3.08 \times 10^0 \text{ yr}^{-1} \text{ Gpc}^{-3}$	14.95 yr^{-1}
Z_{IZw18}	$\langle M^{2.5} \rangle$	$R_{z=0}$	R_D	R_{cSFR}	$R_{D,cSFR}$
NS-NS	$1.27 M_{\odot}^{2.5}$	$1.00 \times 10^1 \text{ yr}^{-1} \text{ Gpc}^{-3}$	0.27 yr^{-1}	$3.28 \times 10^1 \text{ yr}^{-1} \text{ Gpc}^{-3}$	0.87 yr^{-1}
NS-BH	$32.3 M_{\odot}^{2.5}$	$6.61 \times 10^{-3} \text{ yr}^{-1} \text{ Gpc}^{-3}$	0.00 yr^{-1}	$1.55 \times 10^{-2} \text{ yr}^{-1} \text{ Gpc}^{-3}$	0.01 yr^{-1}
BH-NS	$35.5 M_{\odot}^{2.5}$	$1.54 \times 10^1 \text{ yr}^{-1} \text{ Gpc}^{-3}$	11.40 yr^{-1}	$5.32 \times 10^1 \text{ yr}^{-1} \text{ Gpc}^{-3}$	39.34 yr^{-1}
BH-BH	$1720 M_{\odot}^{2.5}$	$1.68 \times 10^1 \text{ yr}^{-1} \text{ Gpc}^{-3}$	603.02 yr^{-1}	$3.45 \times 10^1 \text{ yr}^{-1} \text{ Gpc}^{-3}$	1235.27 yr^{-1}
optimistic	$\langle M^{2.5} \rangle$	$R_{z=0}$	R_D	R_{cSFR}	$R_{D,cSFR}$
NS-NS	$1.31 M_{\odot}^{2.5}$	$7.09 \times 10^1 \text{ yr}^{-1} \text{ Gpc}^{-3}$	1.94 yr^{-1}	$1.59 \times 10^2 \text{ yr}^{-1} \text{ Gpc}^{-3}$	4.37 yr^{-1}
NS-BH	$19.4 M_{\odot}^{2.5}$	$0.00 \times 10^0 \text{ yr}^{-1} \text{ Gpc}^{-3}$	0.00 yr^{-1}	$0.00 \times 10^0 \text{ yr}^{-1} \text{ Gpc}^{-3}$	0.00 yr^{-1}
BH-NS	$21.9 M_{\odot}^{2.5}$	$1.34 \times 10^1 \text{ yr}^{-1} \text{ Gpc}^{-3}$	6.11 yr^{-1}	$2.44 \times 10^1 \text{ yr}^{-1} \text{ Gpc}^{-3}$	11.17 yr^{-1}
BH-BH	$275 M_{\odot}^{2.5}$	$4.34 \times 10^1 \text{ yr}^{-1} \text{ Gpc}^{-3}$	248.34 yr^{-1}	$1.09 \times 10^2 \text{ yr}^{-1} \text{ Gpc}^{-3}$	623.03 yr^{-1}

@design sensitivity: 1–5 detections per yr



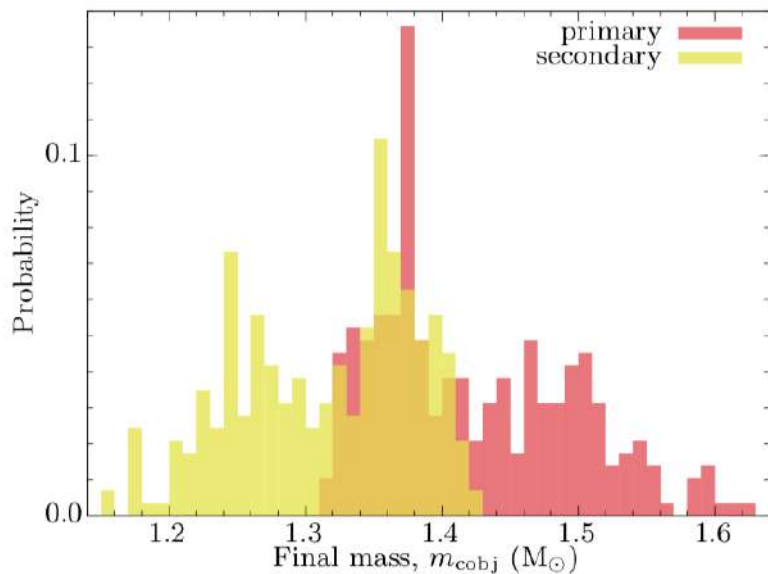
Kruckow et al. (2018), MNRAS



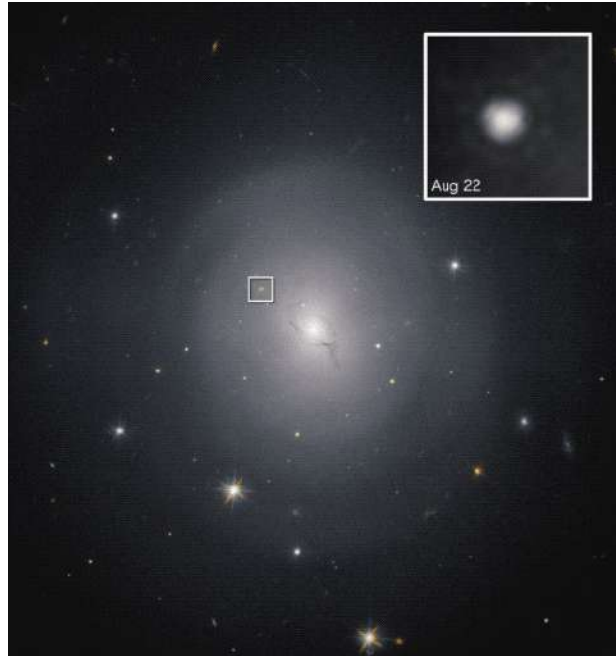


Kruckow, et al. (2018), MNRAS

We find age solutions from
<100 Myr to >10 Gyr

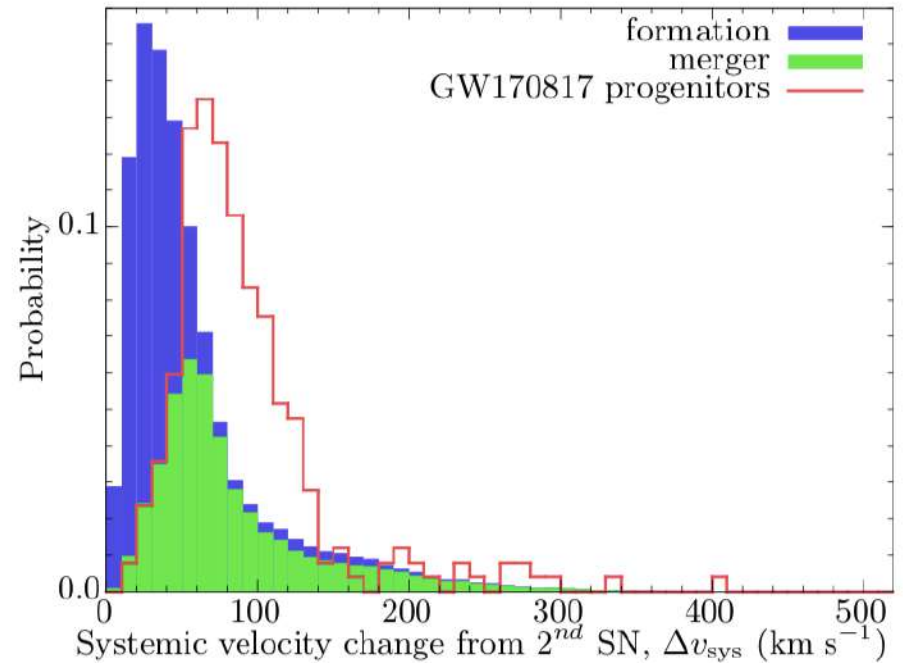


Our NS mass solutions for GW170817
are typical for Galactic DNS systems



NGC 4993

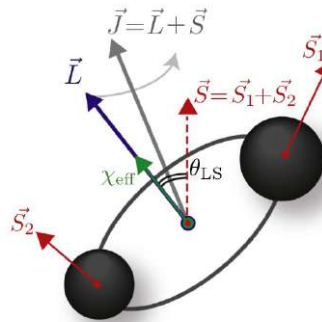
Kruckow et al. (2018), MNRAS



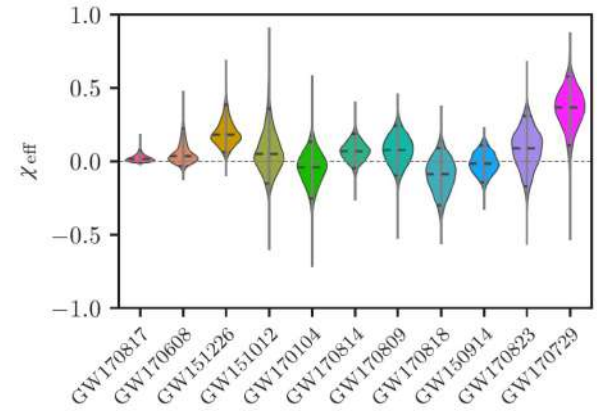
For NGC 4393, the escape velocity at the location of GW170817 is about 350 km s^{-1} (Pan et al. 2017), much larger than the typical systemic velocities we obtain in our simulations.

LIGO network measurements:

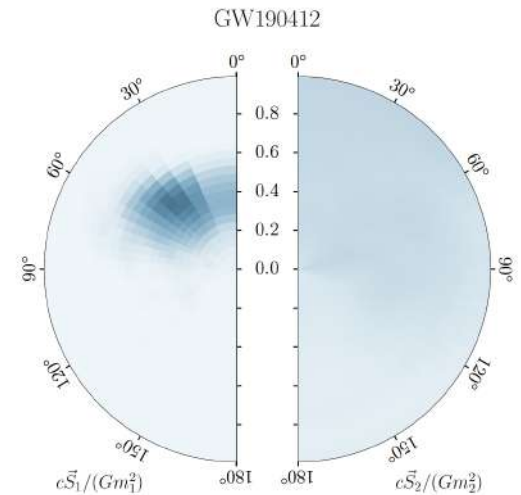
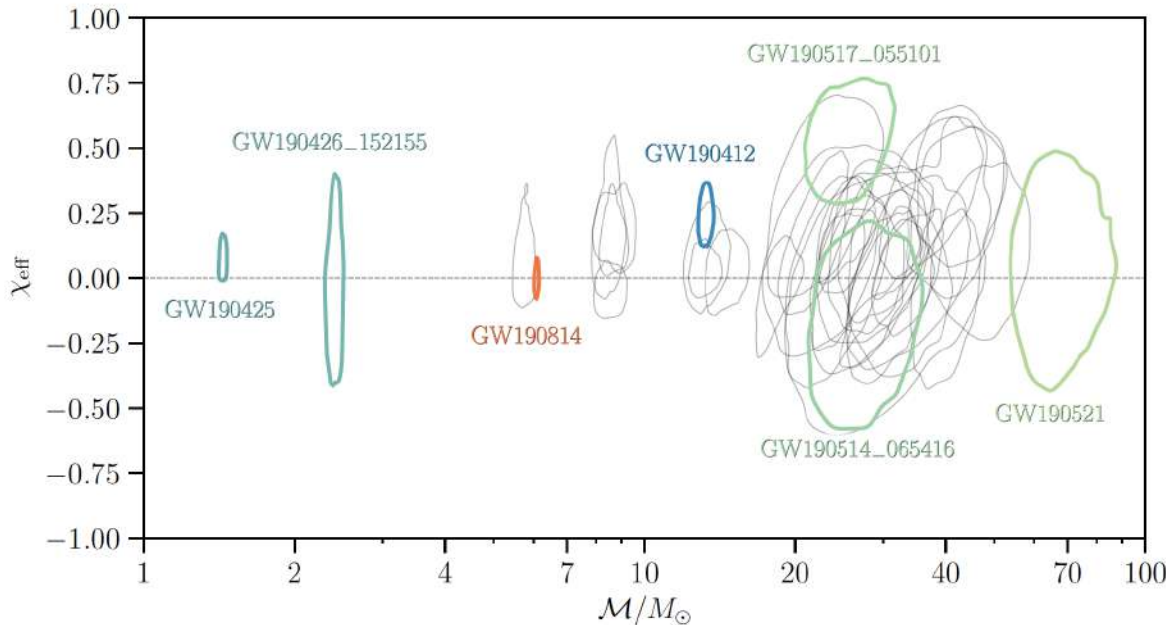
$$\chi_{\text{eff}} \equiv \frac{(M_1 \vec{\chi}_1 + M_2 \vec{\chi}_2)}{M} \cdot \frac{\vec{L}}{|\vec{L}|} = \frac{\chi_1 \cos \theta_1 + q \chi_2 \cos \theta_2}{1 + q}$$



LIGO run O1+O2:



LIGO run O3a:



Intrepretation of BHBH mergers spins

Given that the far majority of all BH-BH mergers reported so far have near-zero effective spins leads to only three potential explanations (e.g. Belczynski et al., 2020):

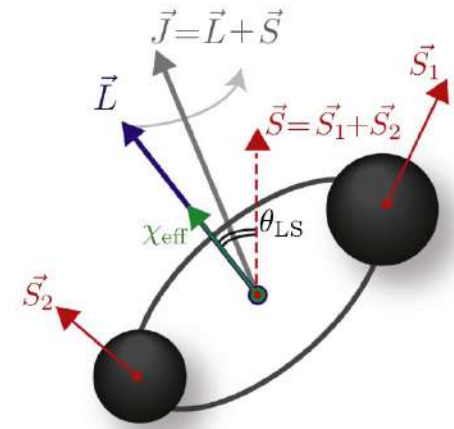
If the individual BH spin magnitudes are large, then:

- (i) Either both BH spin vectors must be nearly in the orbital plane, or
- (ii) the spin angular momenta of the BHs must be oppositely directed and similar in magnitude.

Finally, there is also the possibility that:

- (iii) the BH spin magnitudes are small.

Belczynski et al. (2020) demonstrate that they can reproduce the observed distribution of low χ_{eff} values within the classical isolated binary evolution scenario (the CE channel) of BH-BH formation assuming efficient angular momentum transport.



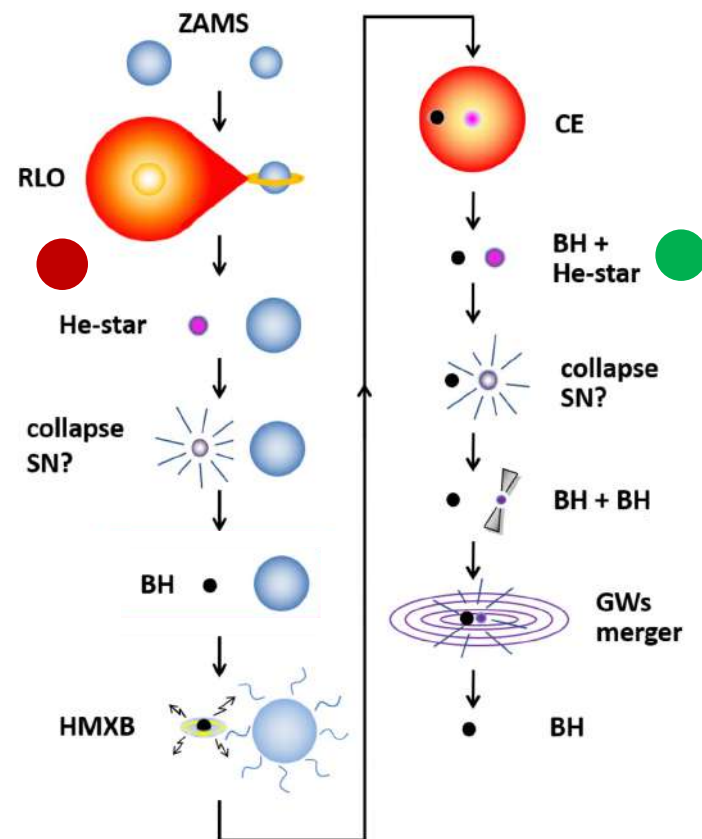
Expectations from stellar evolution:

See e.g.: Kushnir et al. (2016), Hotokezaka & Piran (2017), Zaldarriaga et al. (2018), Fuller & Ma (2019), Qin et al. (2019), Belczynski et al. (2020), Bavera et al. (2020)

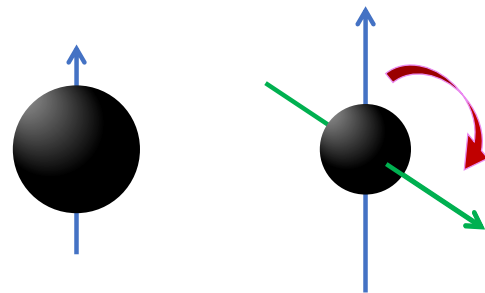
- First-born BH will be spinning rather slow
- Second-born BH will be spinning rather fast

1. Efficient angular momentum transport by viscosity will couple the stellar core to its envelope, thereby slowing the spin of the core as the envelope expands when it becomes a giant star. **Contradiction***

2. Tidal interactions between the first-born BH and the close-by naked-core WR-star (progenitor of the second-born BH) causes the latter to spin up efficiently.



* In clear tension with observations of BH spins in HMXBs (see Lecture 9)



sort hul 1

sort hul 2

THE ASTROPHYSICAL JOURNAL, 938:66 (14pp), 2022 October 10

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Tossing Black Hole Spin Axes

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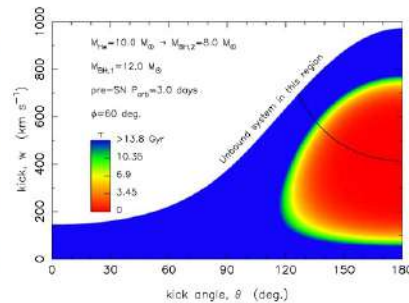
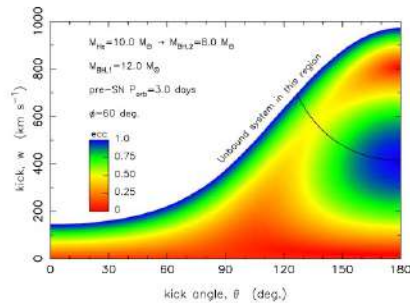
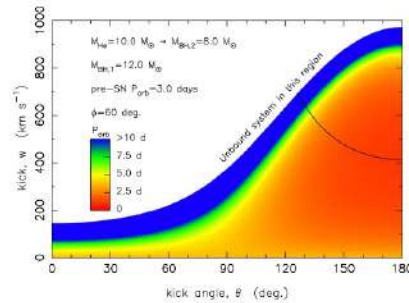
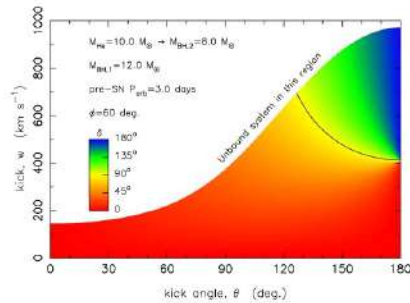
Received 2022 May 4; revised 2022 August 2; accepted 2022 August 2; published 2022 October 13

Abstract

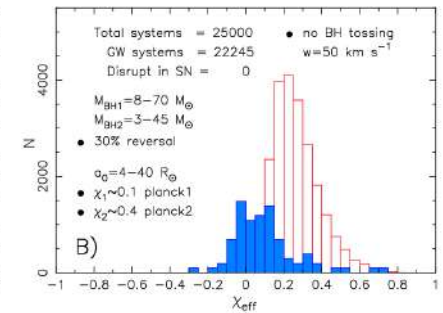
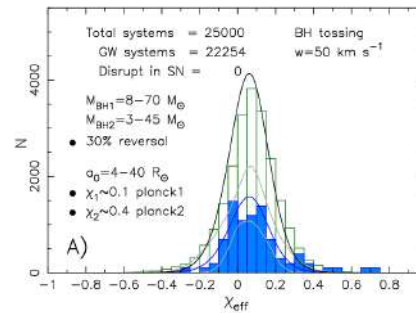
The detection of double black hole (BH+BH) mergers provides a unique possibility to understand their physical properties and origin. To date, the LIGO–Virgo–KAGRA network of high-frequency gravitational-wave

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Tauris



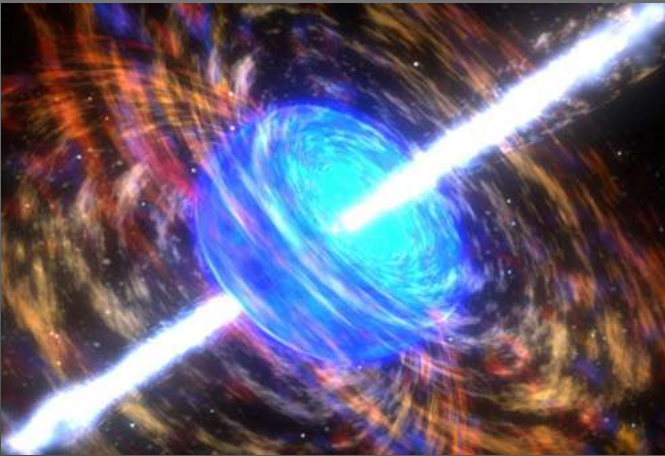
$$\chi_{\text{eff}} \equiv \frac{(M_1 \chi_1 + M_2 \chi_2)}{M_T} \cdot \frac{L}{|L|}$$



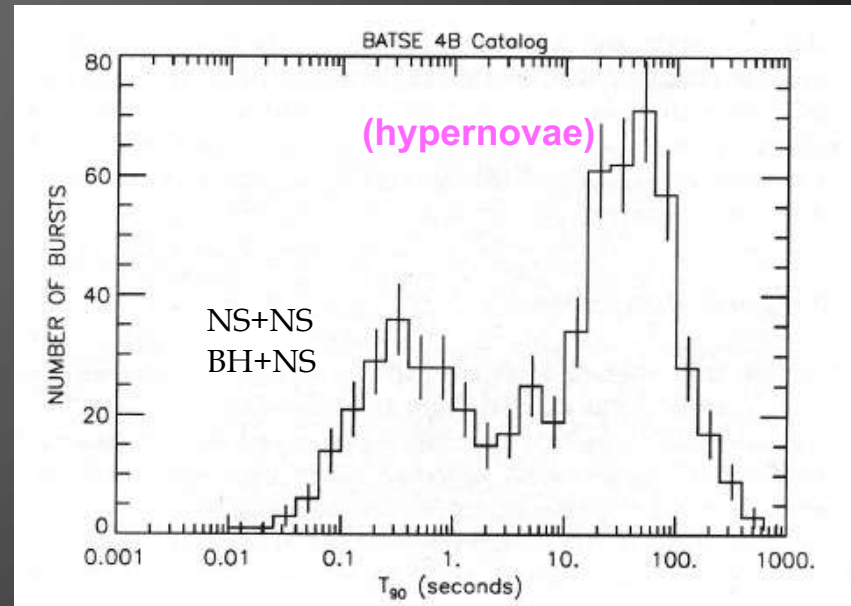
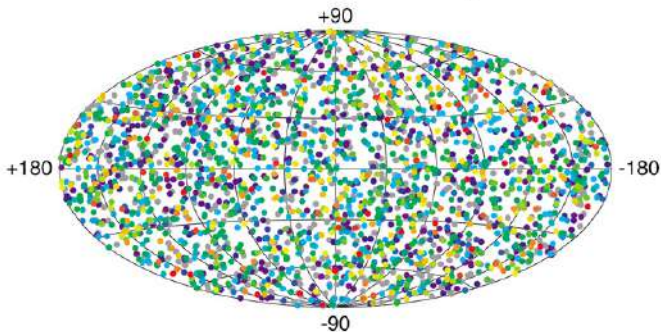
short Gamma-Ray Burst (sGRB)

sGRB may be launched (within 2 sec) via either:

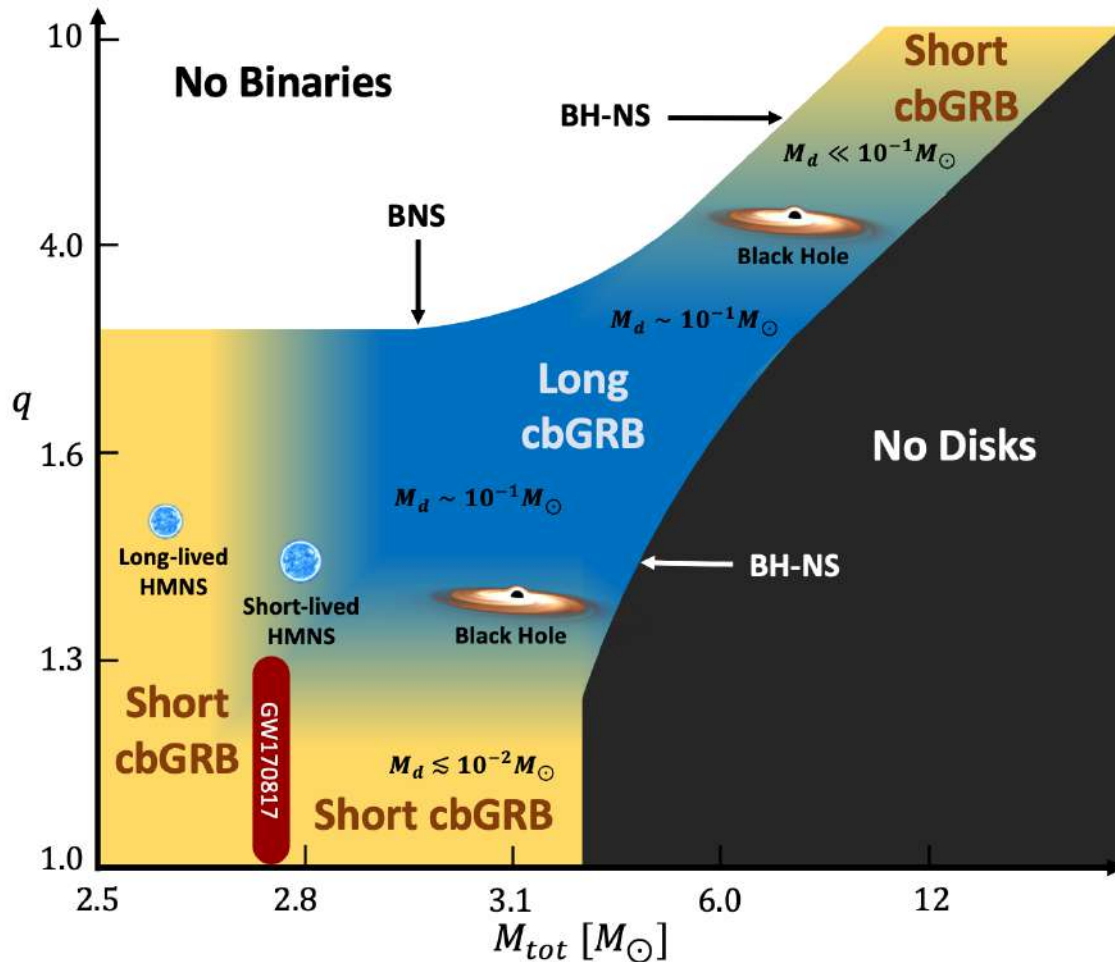
- a) Pair-annihilation of neutrinos
- b) Strong (and twisted) B-field (Blandford-Znajek mechanism, MRI)



2704 BATSE Gamma-Ray Bursts



Unified Picture of Gamma-Ray Bursts



Gottlieb et al. (2023)

<https://arxiv.org/abs/2309.00038>

Figure 2. The outcomes of compact object mergers and their ability to power various cbGRBs sub-classes as a function of the binary mass ratio and total mass. lbGRBs occur in high M_{tot} and high q BNS mergers that form a massive BH disk, or in high pre-merger BH spin and low mass ratio BH-NS mergers (blue region). sbGRBs may arise either from equal mass ratio BNS mergers (bottom yellow region) and low pre-merger BH spin/high mass ratio BH-NS mergers (top yellow region), or by HMNSs formed in BNS mergers with $M_{tot} \lesssim 2.8 M_{\odot}$ (left yellow region). The absence of evidence for distinct sub-classes of sbGRBs suggests that either BHs or HMNSs are likely to be the sole origin of these events, i.e. only one of the proposed sbGRB scenarios is correct. The Galactic BNS mass distribution, the bimodal GRB duration distribution, and GW170817 observations favor HMNSs as the engine of sbGRB jets.

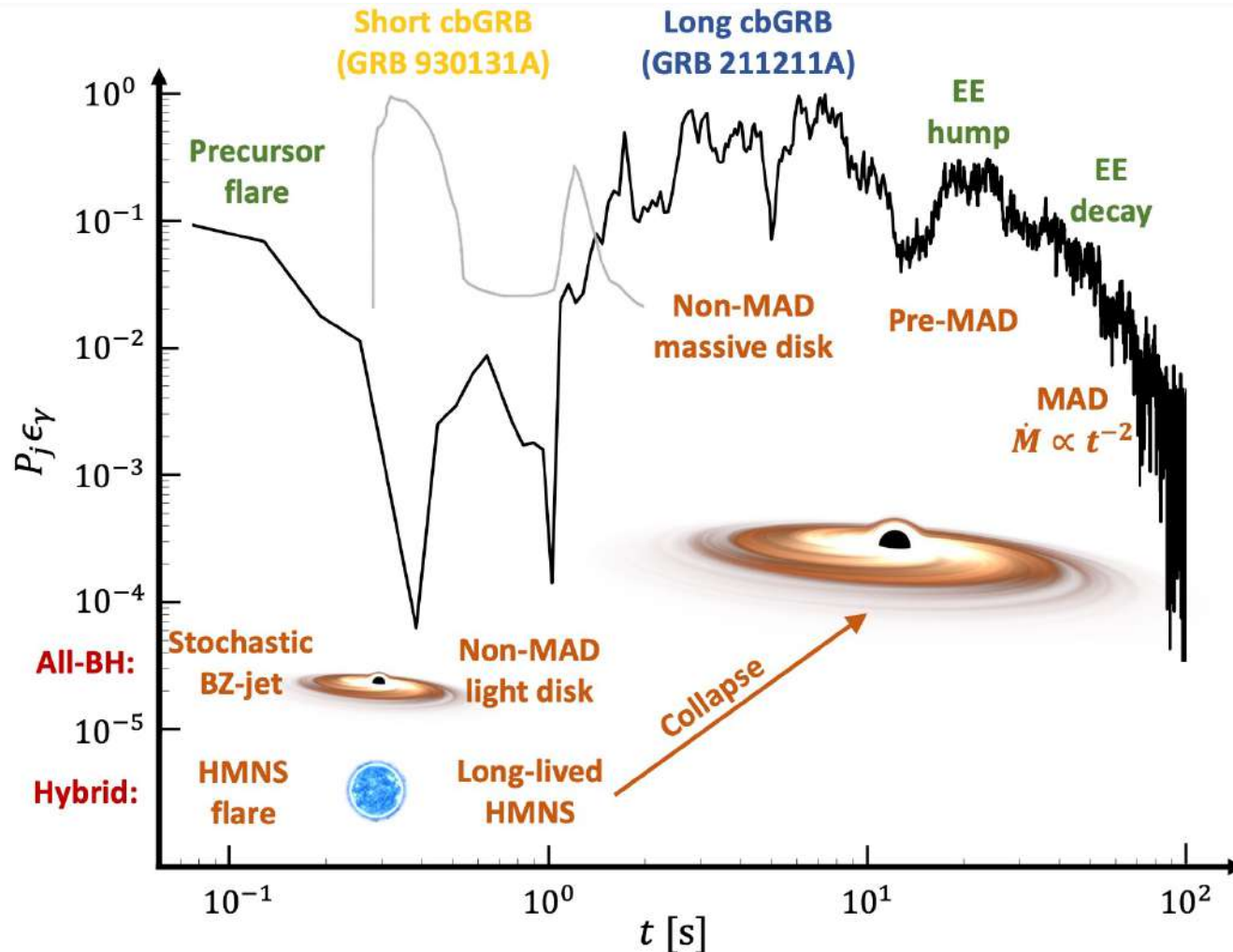


Figure 3. An illustration of how the underlying physics of the merger product (orange) in the hybrid and all-BH scenarios (red) translates into different phases in the cbGRB light curves: sbGRB (yellow), lbGRBs (blue) and preceding and succeeding phases (green). Representations of the light curves of the lbGRB 211211A (Rastinejad et al. 2022) and sbGRB 930131A (Kouveliotou et al. 1994) are shown in black and gray, respectively, in a log-log scale.

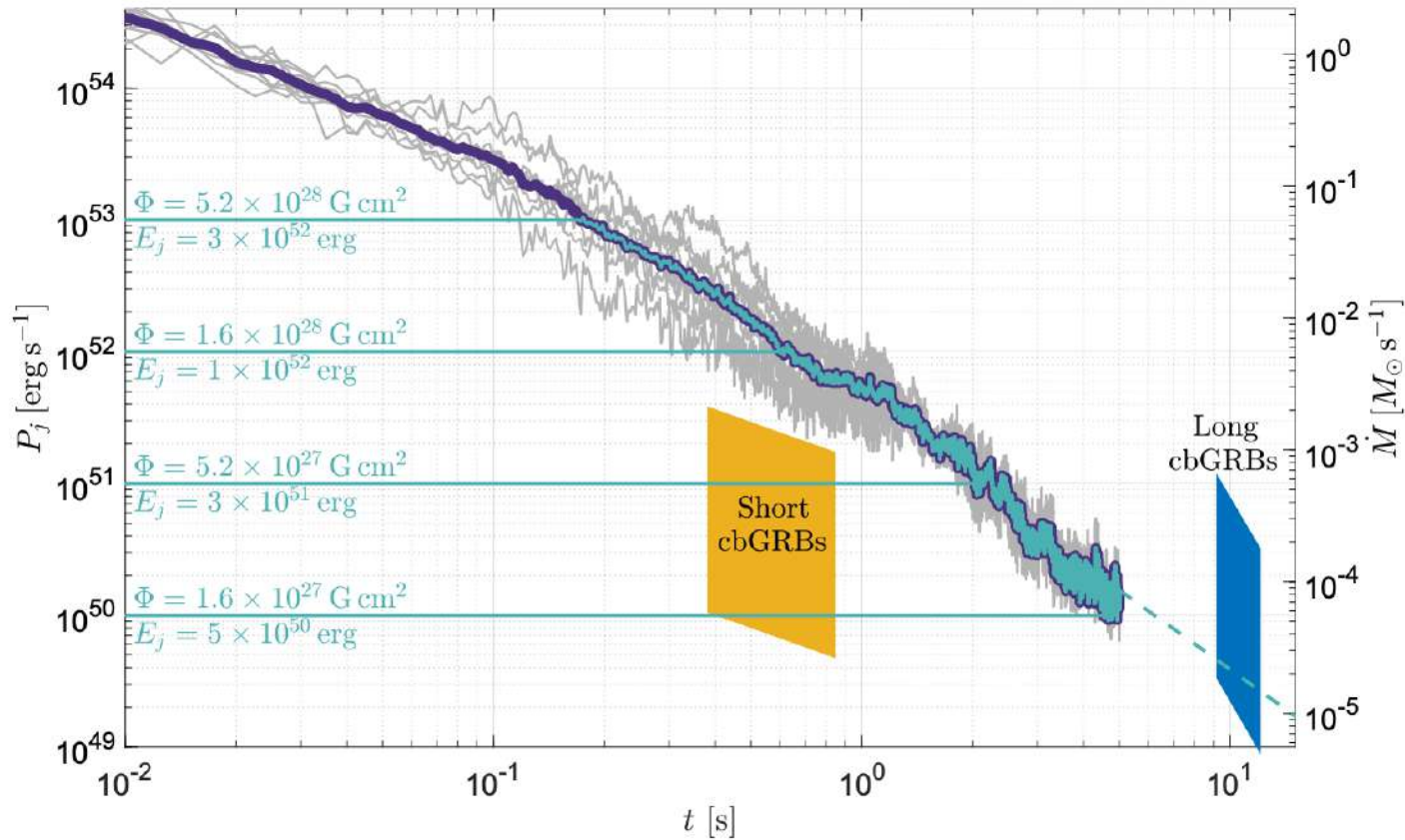
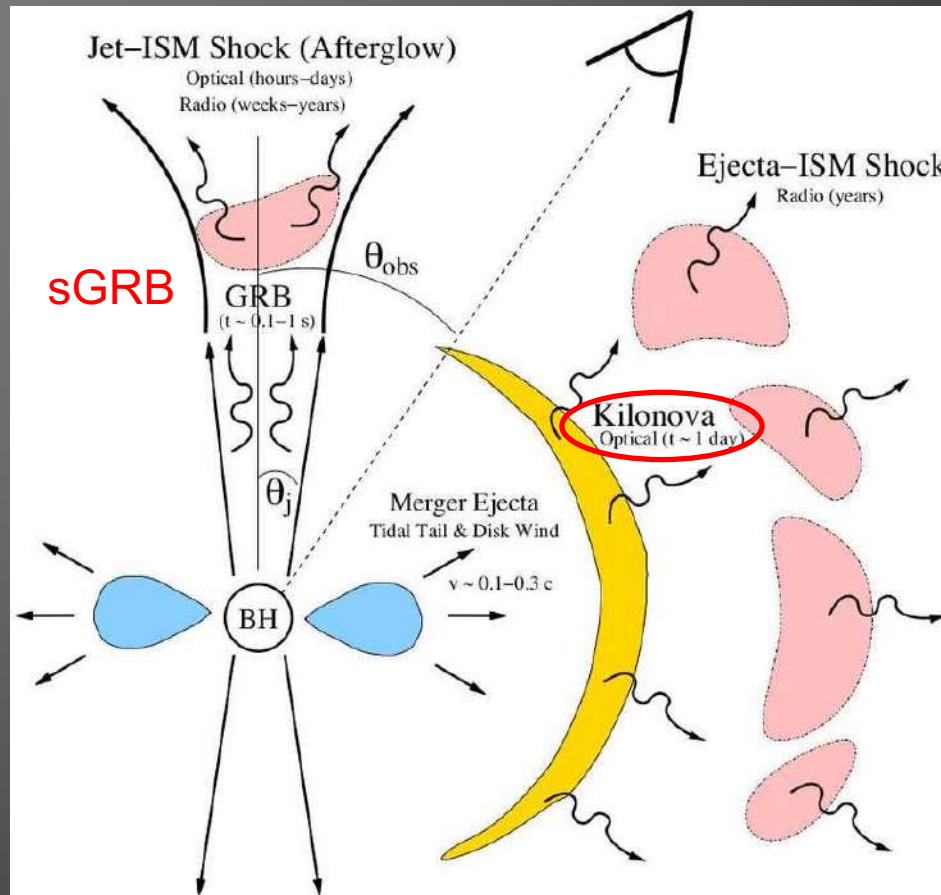


Figure 1. The jet power evolution of post-merger accretion disks for varying levels of magnetic flux ranging from non-MAD to MAD. Gray lines show the post-merger mass accretion rate evolution (right vertical axis) obtained for 4 BH–NS merger simulations (Gottlieb et al. 2023b) and the 5 BNS merger simulations presented here, all of which generate massive disks $M_d \approx 0.1 M_\odot$. The purple line delineates the logarithmic average of these mass accretion rates, which constitutes the maximum jet power assuming $\eta_a = 1$ corresponding to a BH spin $a \approx 0.87$ (left vertical axis). Turquoise lines illustrate schematically the jet power evolution for different assumptions about the dimensional magnetic flux threading the BH, Φ , and the corresponding total jet energy, E_j . Since the magnetic flux on the BH is likely accumulated early and hence remains nearly constant before the disk transitions to MAD, the jet power, P_j , is also predicted to be roughly constant at these times. However, once the dimensionless magnetic flux saturates in the MAD state, the jet power saturates at $P_j = \dot{M}c^2$ and thus follows the mass-accretion rate $\dot{M} \propto t^{-2}$ thereafter (we have extrapolated P_j by a dashed line to later times). The yellow (blue) region outlines the estimated average jet power and duration T_{90} (T_{50}) of the sbGRB (lbGRB) population based on prompt emission and afterglow observations (see text). While the jets from such massive disks are either too powerful, or operate for too long, compared to the sbGRB population, BH accretion from such massive disks nicely matches the observed properties of lbGRBs.

Kilonovae

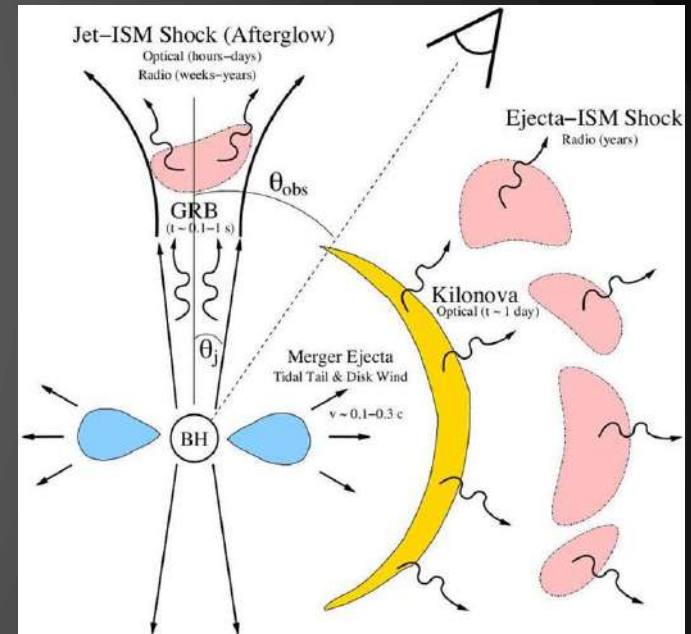
Metzger & Berger (2012)



Excellent reviews:
Giacomazzo, Eichler & Arcones (2019)
Shibata & Hotokezaka (2019)

Optical near-IR kilonova (makronova / sGRB afterglow)

- $0.01-0.1 M_{\text{sun}}$ ejected as dynamic + viscous ejecta (more disk mass for aligned spin axes, also dependence on NS+NS vs BH+NS)
- r-process nucleosynthesis (heavy n-rich nuclei).
- The decay to stability powers an EM transient. (Abbott et al. 2017, Metzger 2017, Rosswog et al. 2018)
- Peaks at $L=10^{42} \text{ erg s}^{-1}$ after about one day
Faint (21-24 mag) and fast decay (hrs-days)
The EM signal (kilonova) is affected by mass composition (Y_e), and temperature of ejecta.
- ISM powered X-ray and radio afterglow from synchrotron emission (a few $100 \mu\text{Jy}$, weeks-months)
- Neutrinos are not expected to be detected (even in next generation detectors) due to the large distances.



Metzger & Berger (2012)

Element Origins

1 H																	2 He	
3 Li	4 Be											5 B	6 C	7 N	8 O	9 F	10 Ne	
11 Na	12 Mg											13 Al	14 Si	15 P	16 S	17 Cl	18 Ar	
19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr	
37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe	
55 Cs	56 Ba			72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	84 Po	85 At	86 Rn
87 Fr	88 Ra																	
			57 La	58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb	71 Lu	
			89 Ac	90 Th	91 Pa	92 U												

Merging Neutron Stars
Dying Low Mass Stars

Exploding Massive Stars
Exploding White Dwarfs

Big Bang
Cosmic Ray Fission

Based on graphic created by Jennifer Johnson

Metzger & Berger (2012)

The r-process (see Giacomazzo et al. 2019)

Rapid neutron-capture process (r-process) was proposed by Burbidge et al. (1957) and Cameron (1957). It occurs due to fast neutron (n) captures in comparison to beta-decays, and runs very close to the n-drip line.

Accumulation of material along the path occurs whenever isotopes with a closed n-shell are reached (affecting both the n-capture cross sections and the beta-rates).

After the intense supply of free neutrons has ceased, the extremely n-rich isotopes undergo a series of beta-decays to stability.

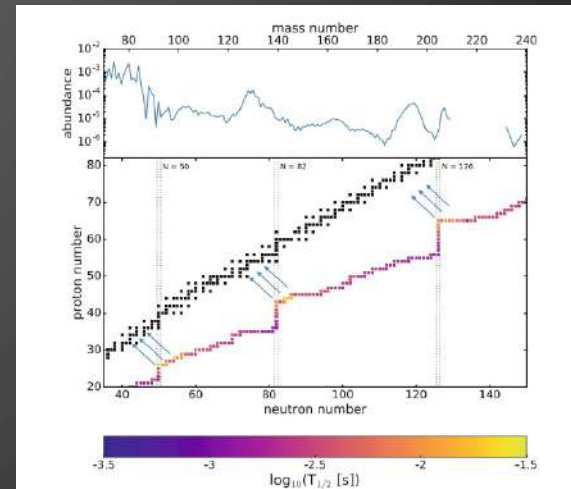
The final average mass can be estimated from:

final average mass number (e.g. Pu-244)
initial seed nuclei (typical iron group, e.g. 60)

$$\langle A \rangle_f = \langle A \rangle_i + \left(\frac{Y_n}{Y_{seed}} \right)_i$$

ratio (e.g. 184) = $\frac{\text{neutron abundance}}{\text{summed up abundances of seed nuclei}}$

Note, a high Y_n is possible where Y_e is low (b/c charge neutrality), e.g. in NS merger ejecta.



Sometimes fission cycles occur in which A_f is divided in two and builds up again.

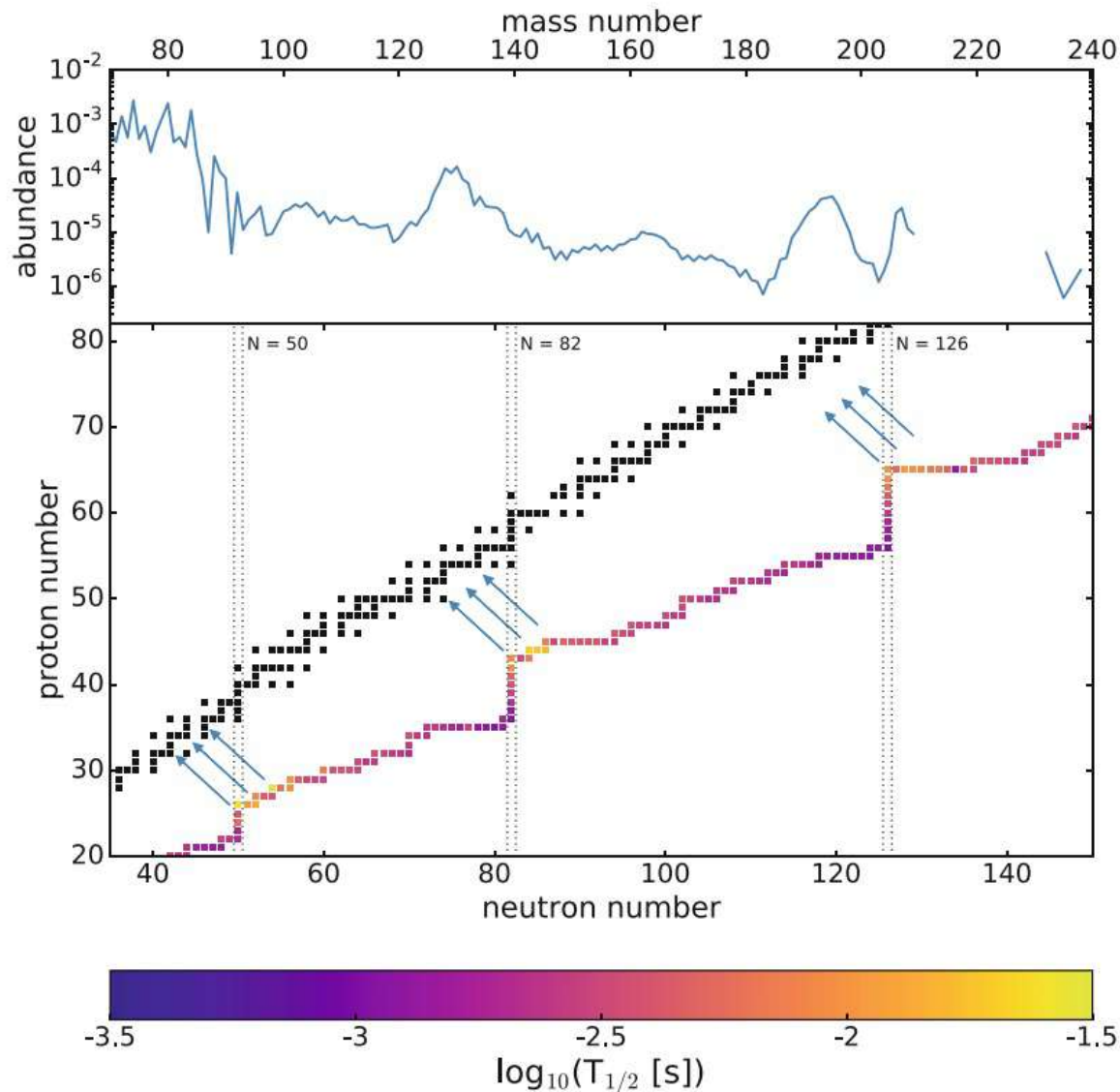


Fig. 11.2 Top panel: Solar r-process abundances as a function of nuclear mass number A . The values are taken from Sneden et al. (2008). Bottom panel: Typical r-process path in the nuclear chart and the corresponding β -decay half-lives according to Möller et al. (2003). Stable isotopes are marked in black and the **magic neutron numbers are indicated by vertical dotted lines**. The overlay of the two panels demonstrates how regions of large $T_{1/2}$ at the magic neutron numbers are responsible for the r-process abundance peaks after decay to stability

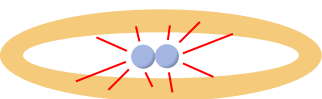
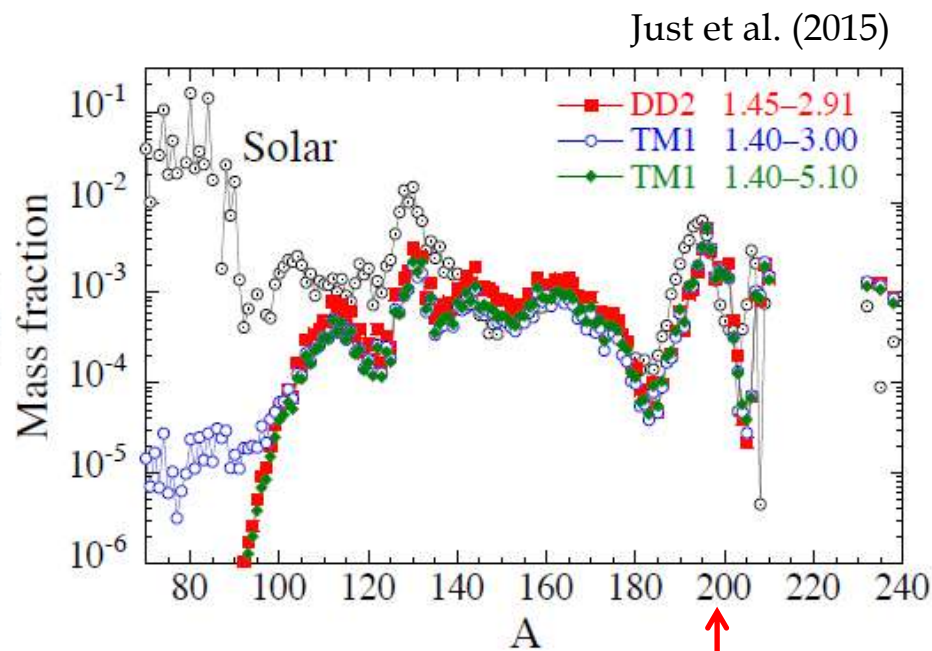
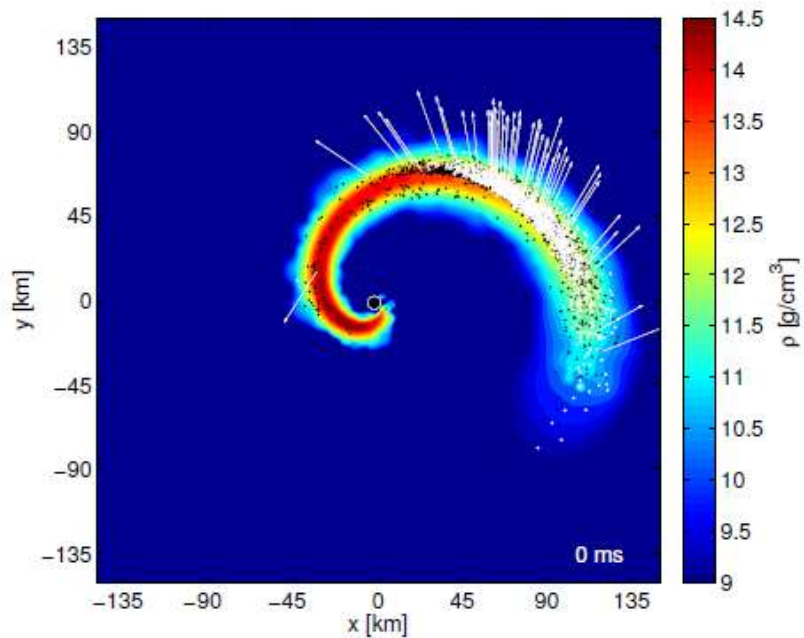
Modelling the r-process

Necessary information: characteristics of light to heavy nuclei between the valley of stability and the n-drip line.

Calibration: Solar spectroscopic data, meteoritic values, deep-sea sediments. In particular, the Galactic evolution of the Eu ($Z=63$) abundance is of interest.

The pattern of abundances of heavy n-capture elements (e.g. $Z=58-76$ or the **Lanthanides**) observed in r-process-rich metal-poor stars are remarkably similar to the Solar System measurements.

Actinides have a clear and unique r-process origin.



Double neutron star mergers:

→ ejection of $10^{-3} - 10^{-2} M_{\odot}$ heavy r-process elements

See also Rosswog (2013, [2013RSPTA.37120272R](https://arxiv.org/abs/2013RSPTA.37120272R))



1A																	2				
1																	2				
H																	He				
3	4															5	6	7	8	9	10
Li	Be															B	C	N	O	F	Ne
11	12											13	14	15	16	17	18				
Na	Mg											Al	Si	P	S	Cl	Ar				
19	20	21	22	23	24	25	26	27	28	29	30	31	32	33	34	35	36				
K	Ca	Sc	Ti	V	Cr	Mn	Fe	Co	Ni	Cu	Zn	Ga	Ge	As	Se	Br	Kr				
37	38	39	40	41	42	43	44	45	46	47	48	49	50	51	52	53	54				
Rb	Sr	Y	Zr	Nb	Mo	Tc	Ru	Rh	Pd	Ag	Cd	In	Sn	Sb	Te	I	Xe				
55	56	57	72	73	74	75	76	77	78	79	80	81	82	83	84	85	86				
Cs	Ba	La	Hf	Ta	W	Re	Os	Ir	Pt	Au	Hg	Tl	Pb	Bi	Po	At	Rn				
87	88	104	105	106	107	108	109	110	111	112			114			116	118				
Fr	Ra	(Rf)	Db	Sg	Bh	Hs	Mt														

The elements from actinium (element 89) to lawrencium (element 103) form a distinct group—the actinides—within the periodic table.

89														
Ac (227)	Actinides													
<i>Thorium Protactinium Uranium Neptunium Plutonium Americium Curium Berkelium Californium Einsteinium Fermium Mendeleevium Nobelium Lawrencium</i>														
90	91	92	93	94	95	96	97	98	99	100	101	102	103	
Th (232)	Pa (231)	U (238)	Np (237)	Pu (244)	Am (243)	Cm (247)	Bk (247)	Cf (251)	Es (252)	Fm (257)	Md (258)	No (259)	Lr (260)	

Mass ejecta and electron fraction

- 1) Dynamical ejecta (tidal disruption)
- 2) Disk ejecta (viscous heating and MHD)

Total amount of ejecta (few $0.001 M_{\text{sun}}$ to $0.1 M_{\text{sun}}$) depends on:

- NS+NS \rightarrow prompt BH formation **or** MNS (meta stable, $\Delta t = 10 \text{ ms} - 10 \text{ s}$)
- Mass ratio ($q < 0.8$ leads to larger yield)
- NS radius and BH spin

Important output parameters are: mass, velocity and electron fraction (Y_e).

Y_e is of key importance for determining the abundance of r-process elements, which again determine the opacity of the EM emission.

Electron fraction and opacity

NS material: $Y_e = 0.05 - 0.1$

However dynamical ejecta could be influenced by weak processes (>10 MeV)

which drive $Y_e = 0.5$ $n + e^+ \rightarrow p + \bar{\nu}_e$ $n + \nu_e \rightarrow p + e^-$

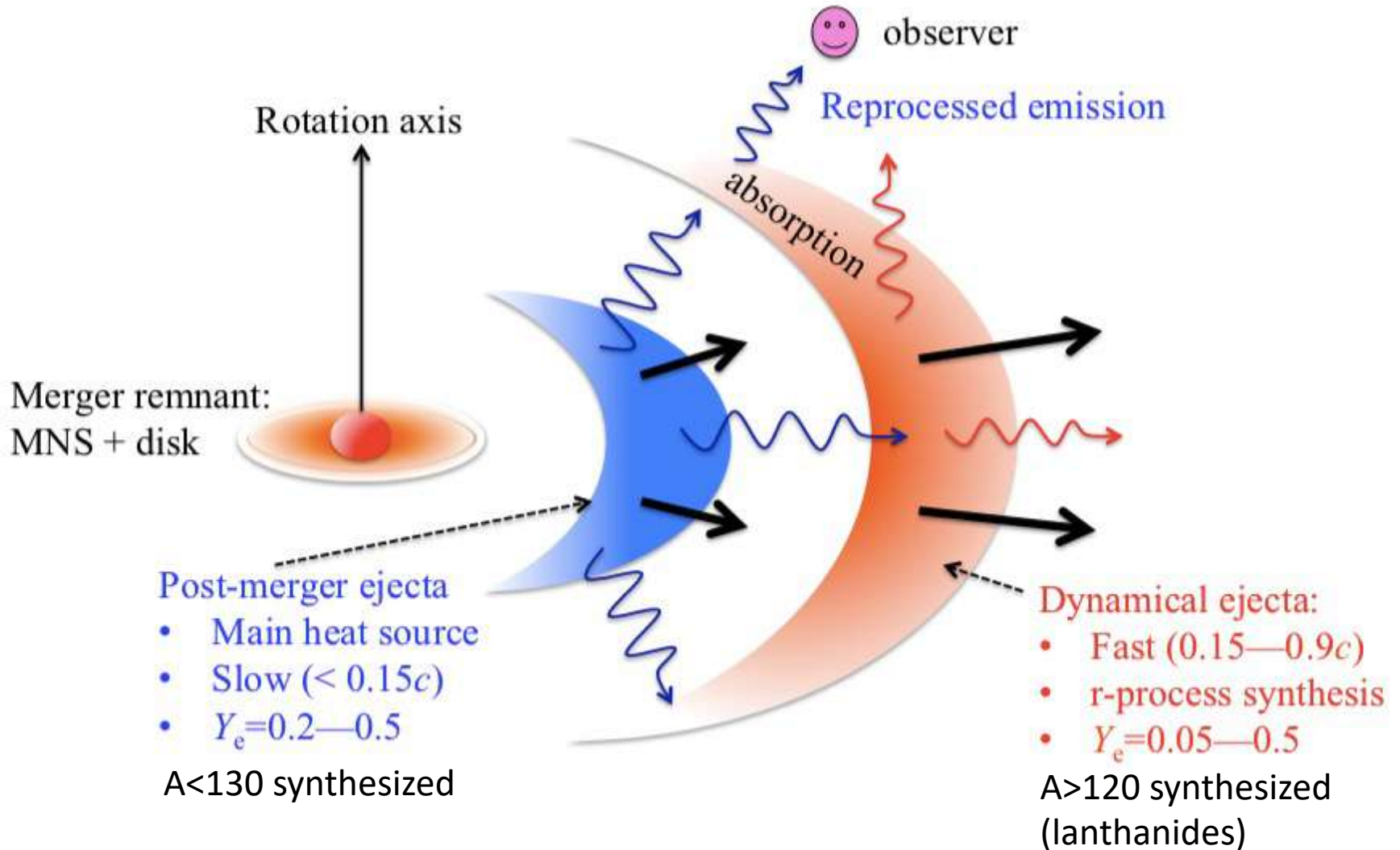
n-rich matter (Y_e small): r-process elements with $A > 120$ are robustly synthesized

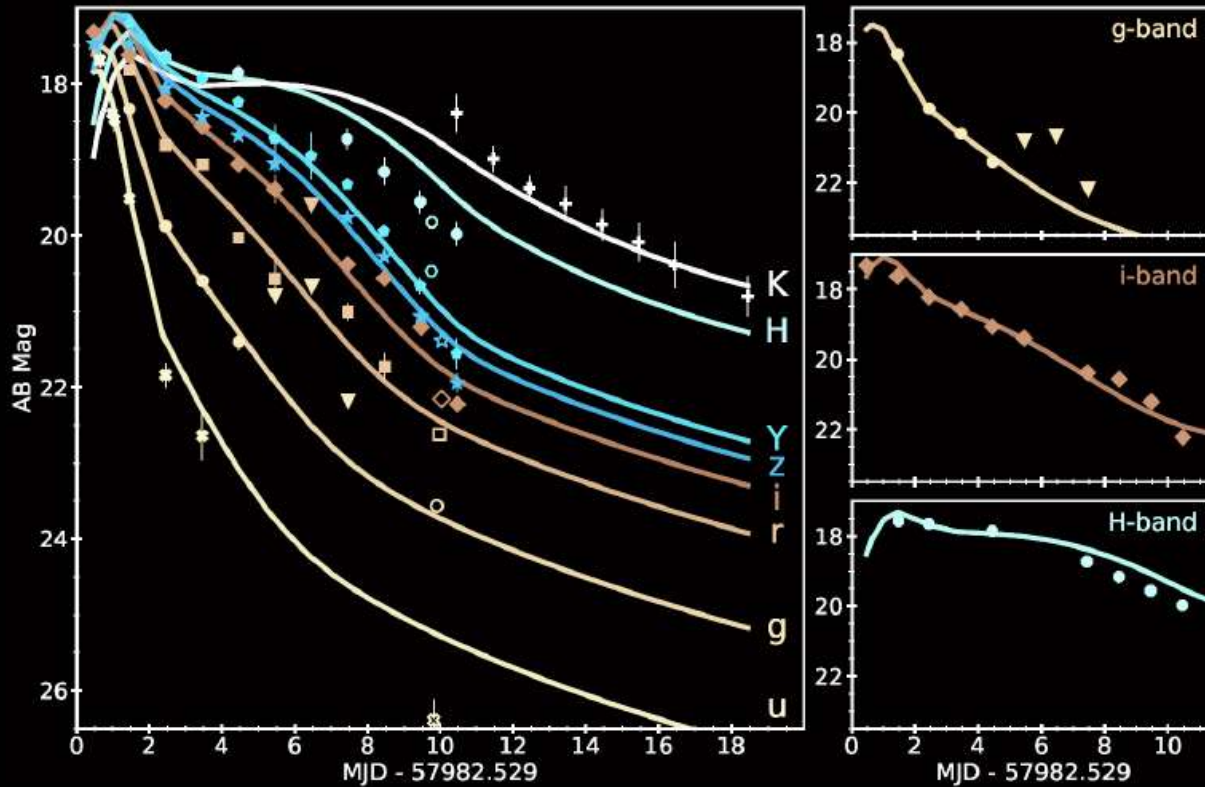
n-poor matter (Y_e large): only r-process elements with $A < 130$ are synthesized
(i.e. lanthenides are not produced)

Opacity (photon): large for lanthenide-rich ejecta (red colour)

small for lanthenide-poor ejecta (blue colour)

Shibata & Kotokezaka (2019)





UV, optical, and near-IR spectra are well fit by a two-component kilonova
 0.02 M_{sun} lanthanide-poor ejecta (blue) and 0.05 M_{sun} lanthanide rich ejecta (red)

Cowperthwaite, ..., DAB et al. ApJ **848** L17 (2017)

The following 5 slides are provided by Duncan Brown.

Distance-constrained GW observations of viewing angle are consistent with EM observations

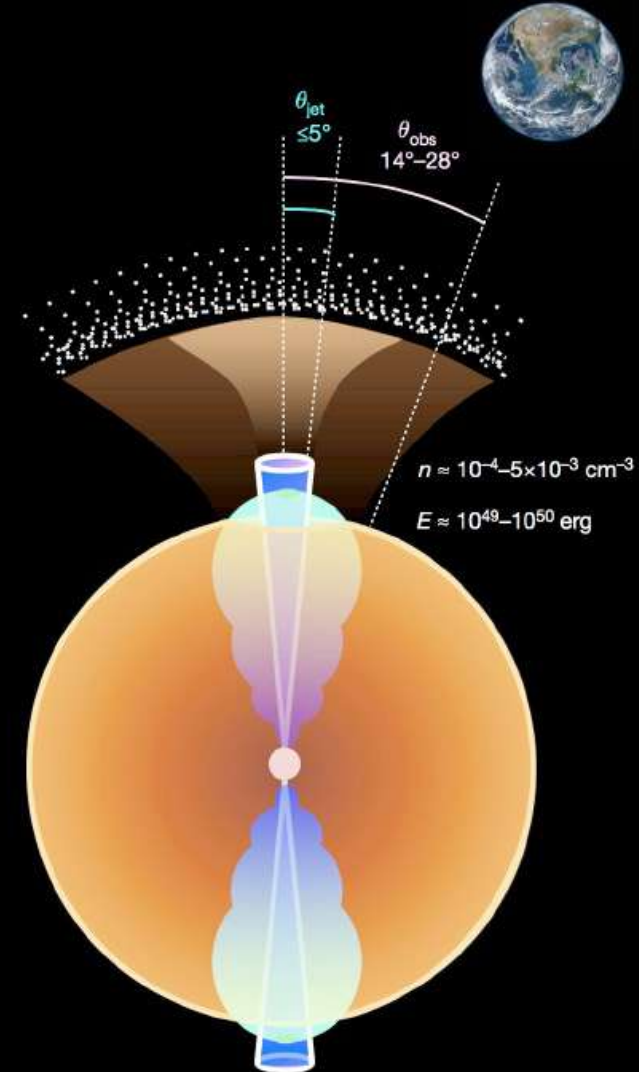
Mooley et al. report 14 - 28 deg from radio

Troja et al. report 21 - 29 deg from broad band observations

GW and EM observations support successful-jet cocoon model (structured jet)

Mooley et al. Nature **561**, 355 (2018)

Troja et al. MNRAS arXiv:1808.06617

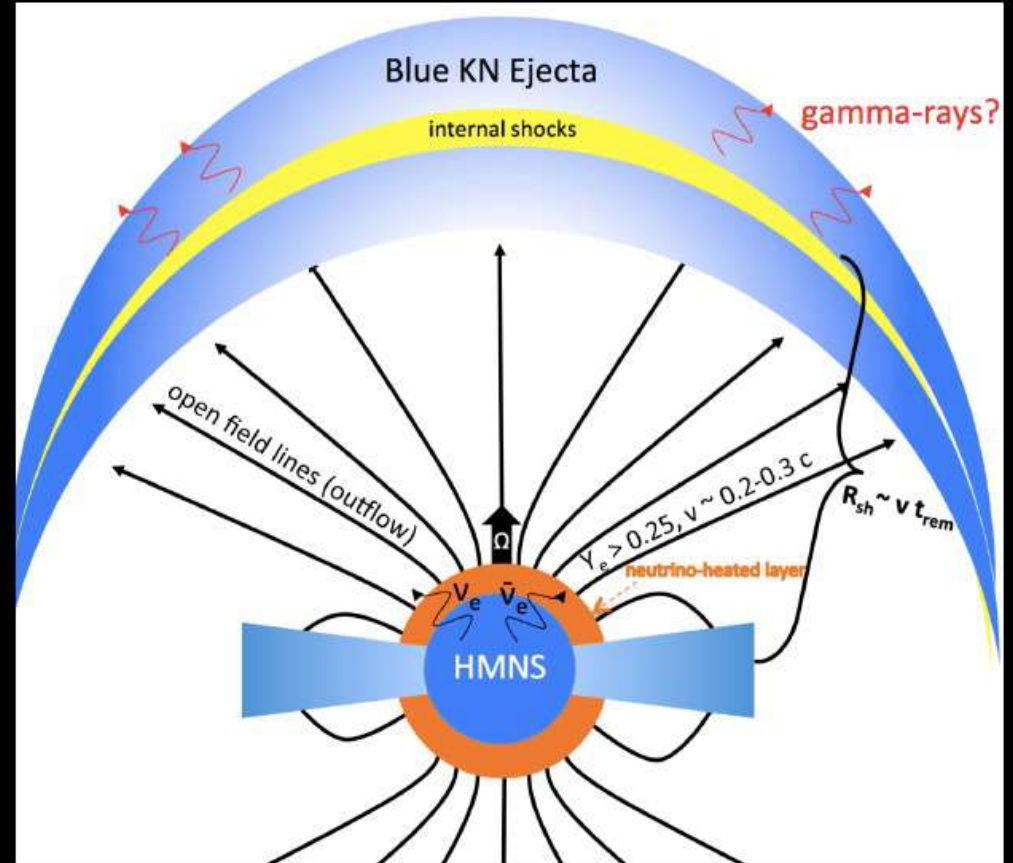


$V_{red} \sim 0.1 c$

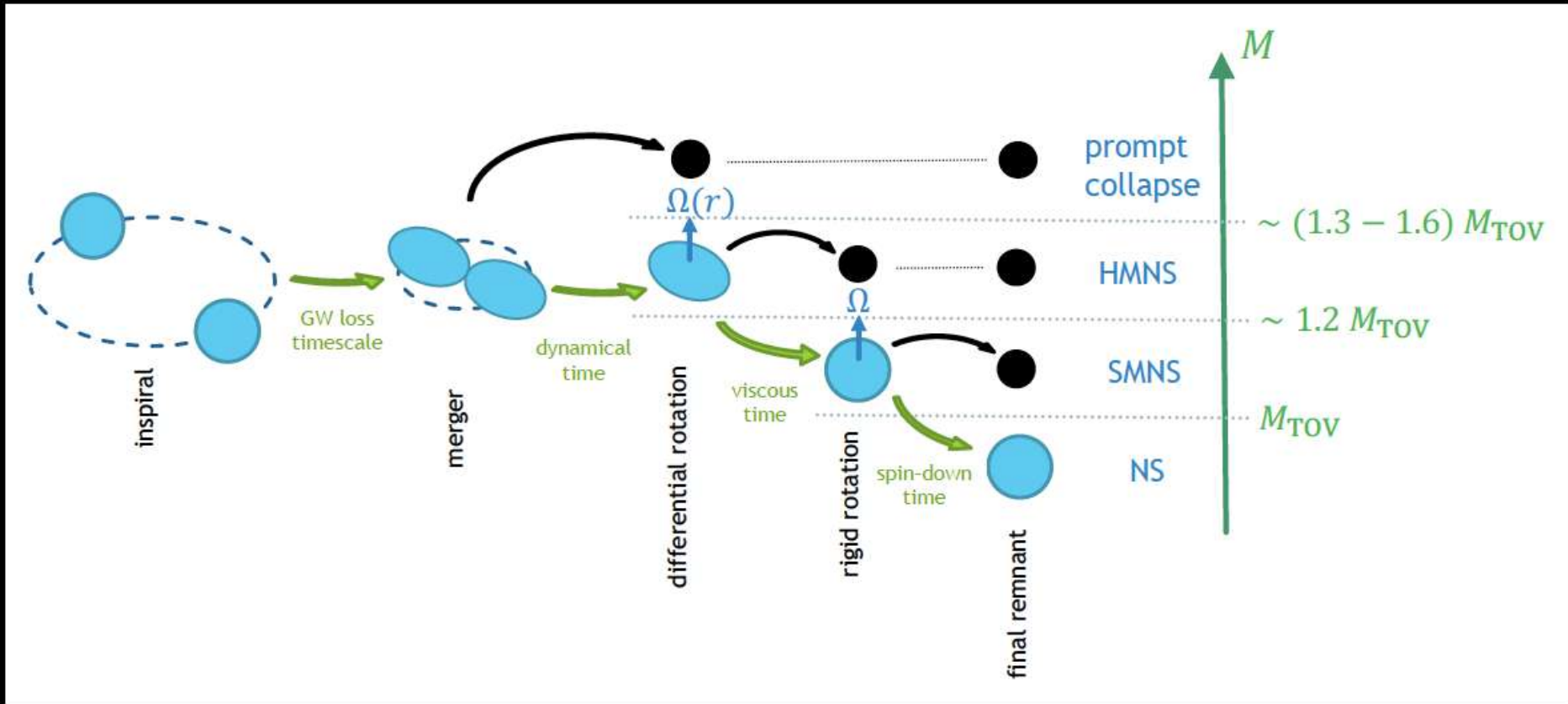
$V_{blue} \sim 0.25 c$

Kilonova light curves suggest the existence of a hyper massive neutron star prior to collapse to a black hole

EM suggests neutron star merger



Metzger, Thompson, Quataert ApJL **856** 101 (2018)



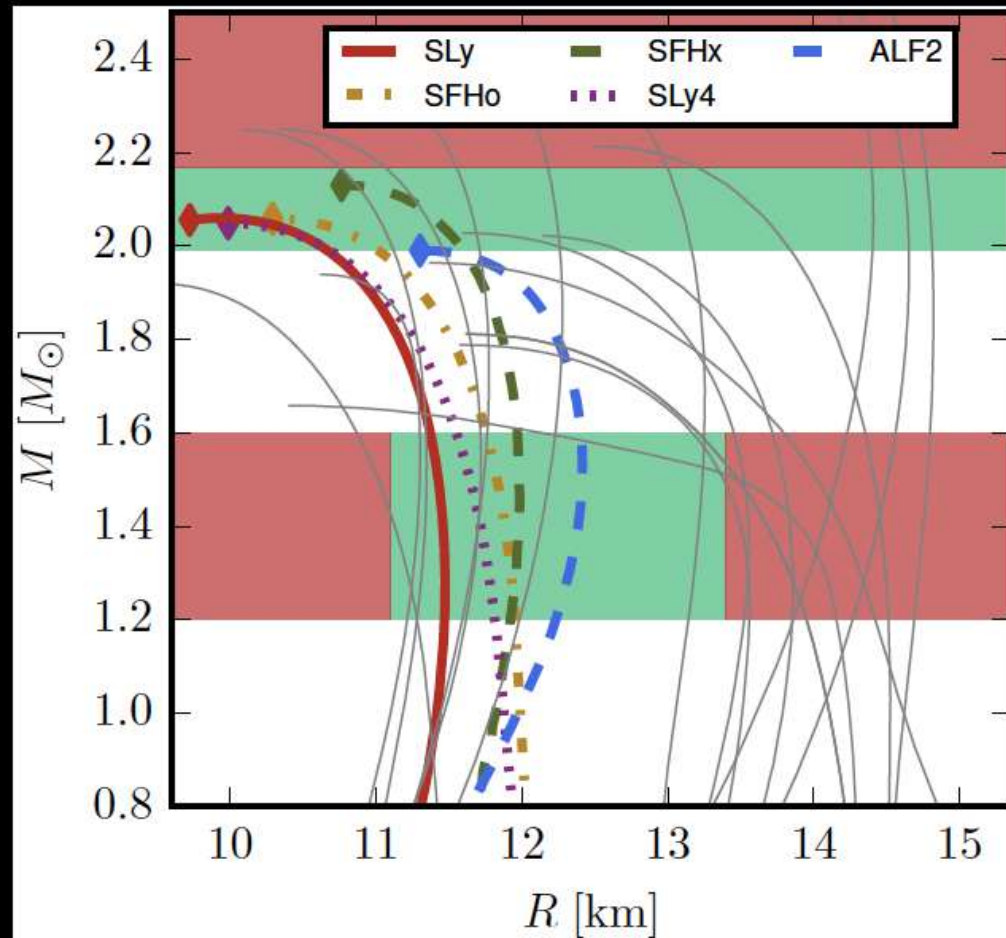
Ben Margalit

The merger remnant also places a constraint on the maximum neutron star mass

The remnant NS cannot be long lived, or there would be too much energy in the EM observantion

$$M_{\text{max}} \leq 2.17M_{\odot} \text{ (90\%)}$$

Margalit and Metzger ApJL 850 19 (2018)



Coughlin, Dietrich, Margalit, Metzger arXiv:1812:04803

The r-process: observational evidence

NS mergers versus core collapse supernovae

GW170817: first firm detection of **kilonova** (EM transient).

Ejected mass $\Delta M_{\text{eject}} \geq 1.5 \times 10^{-2} M_{\odot}$ depending on amount of energy release ending up in the observed emission (Rosswog et al. 2017)



Heavy r-process elements are also observed in atmospheres of old stars and in the Solar system. If they are explained by **core-collapse SNe** the amount of enrichment per explosion is therefore about $10^{-5} M_{\odot}$ for a Galactic SN rate of about 0.01 yr^{-1}



However, studies of a group of stars in the dwarf galaxy Reticulum II supports rare events with large ejecta (NS mergers) compared to frequent events with little ejecta (core-collapse SNe), cf. Beniamini et al. (2016).



Spin distributions of BHs and NSs

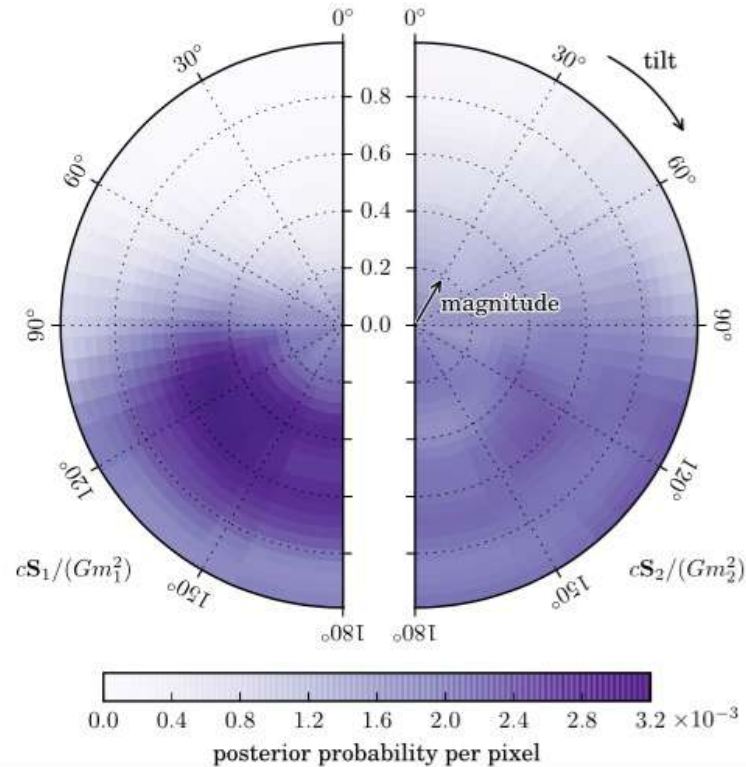
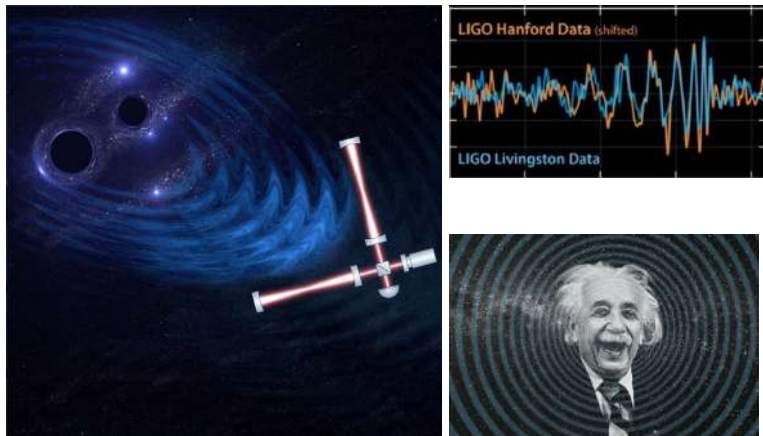
$|X_{\text{eff}}| < 0.35$ at the 90% credible level for all events!
 (degeneracy between projected spins and orbital inclination, masses)

$$\chi_{\text{eff}} \equiv \frac{1}{M} (m_1 \chi_1 + m_2 \chi_2)$$

3G ?

Provides a clue to their astrophysical origin
 e.g. Baibhav et al. (2020)

Tests of GR and other gravity theories

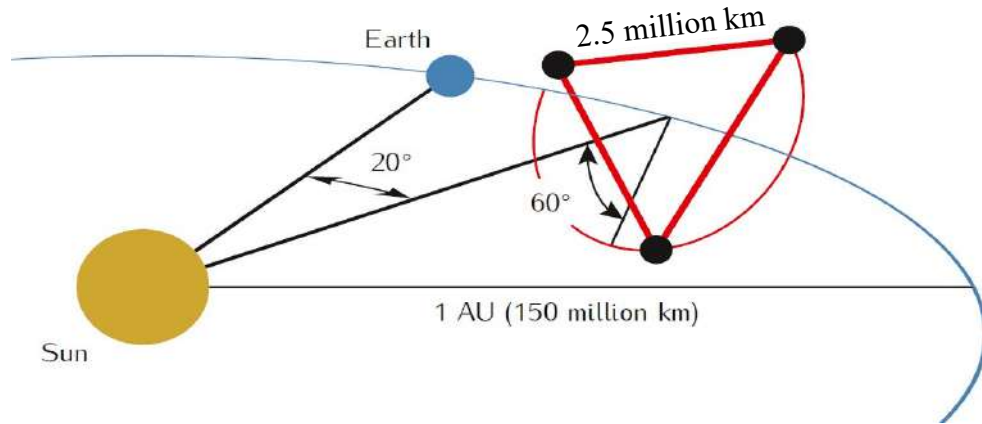
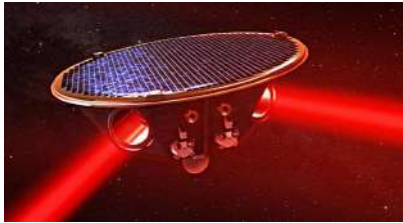


Multi-messenger astrophysics

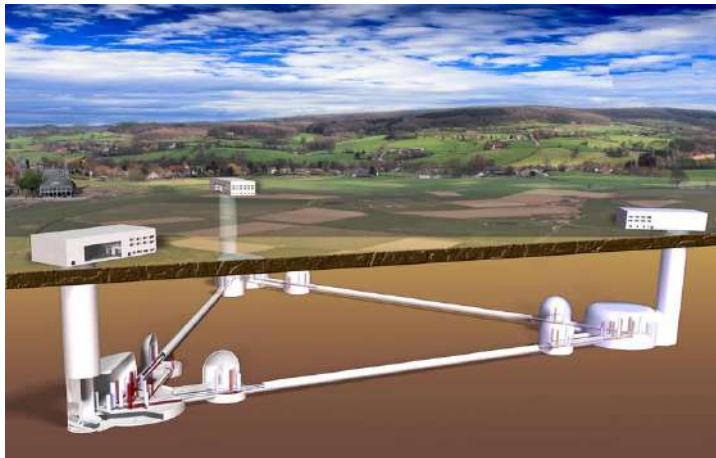
- ✓ GWs
- ✓ Optical
- ✓ X-rays
- ✓ Radio
- ✓ ...more



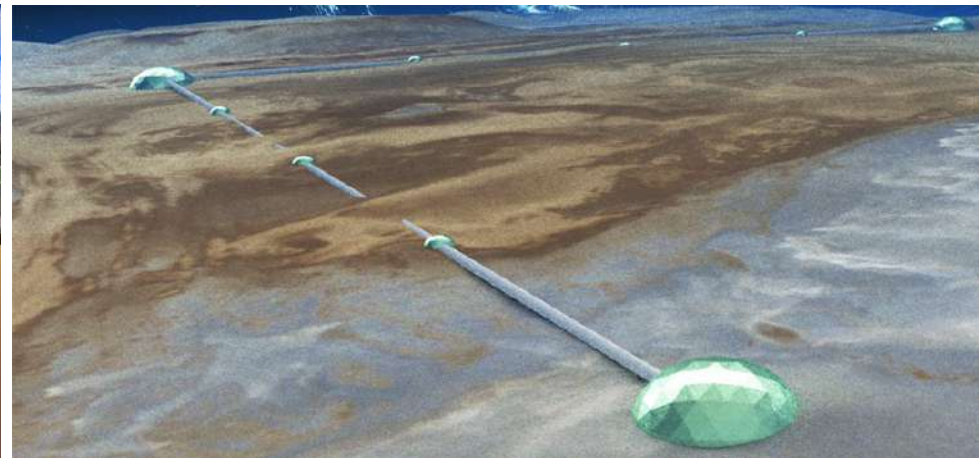
LISA
~2034



3G
~ ?

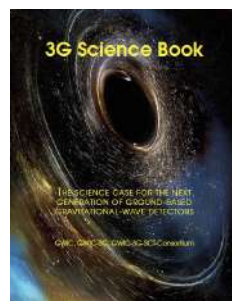


EINSTEIN TELESCOPE



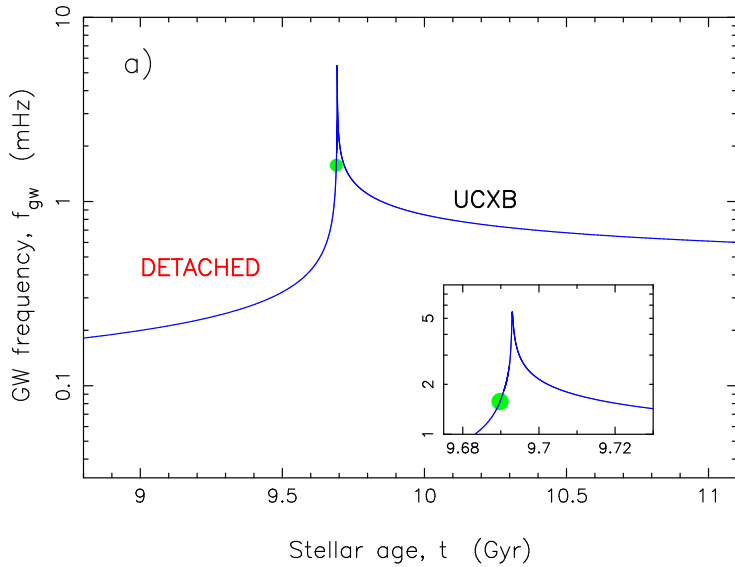
COSMIC EXPLORER

Ask for 3 detectors
(~ 1 billion € each)

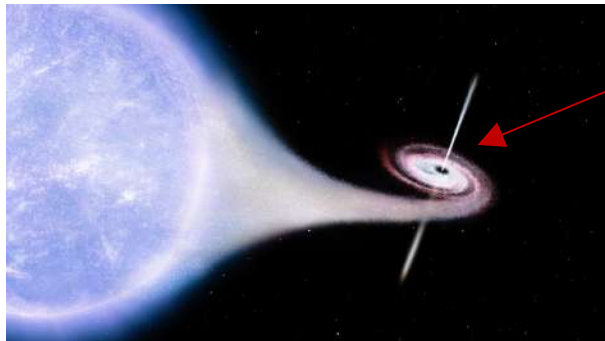
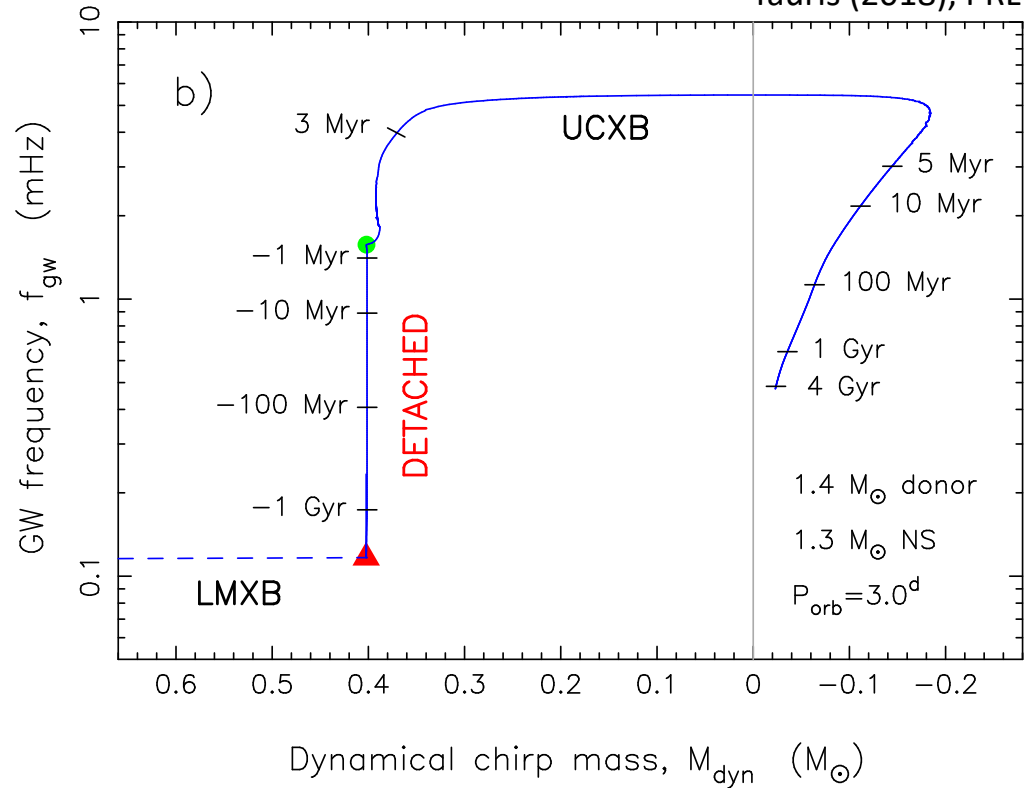


- Detect all BH-BH mergers out to $z \sim 20$
- Detect the BH seeds evolving into SMBHs
- Possibly detect primordial BHs
- Determine the NS EoS to extreme precision
- etc.

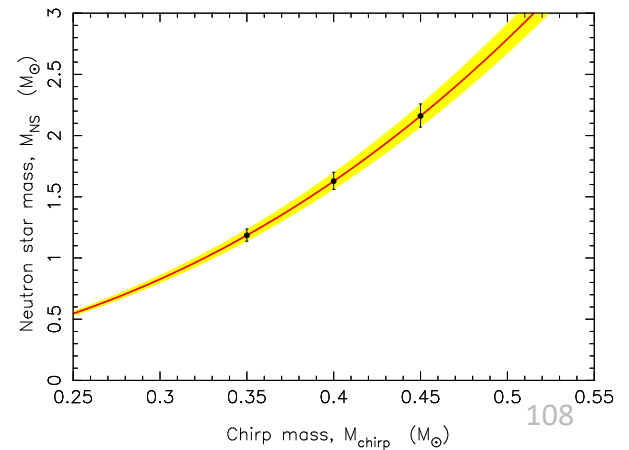
GW spectrum evolution with finite-temperature effects (specific entropy) of the WD donor

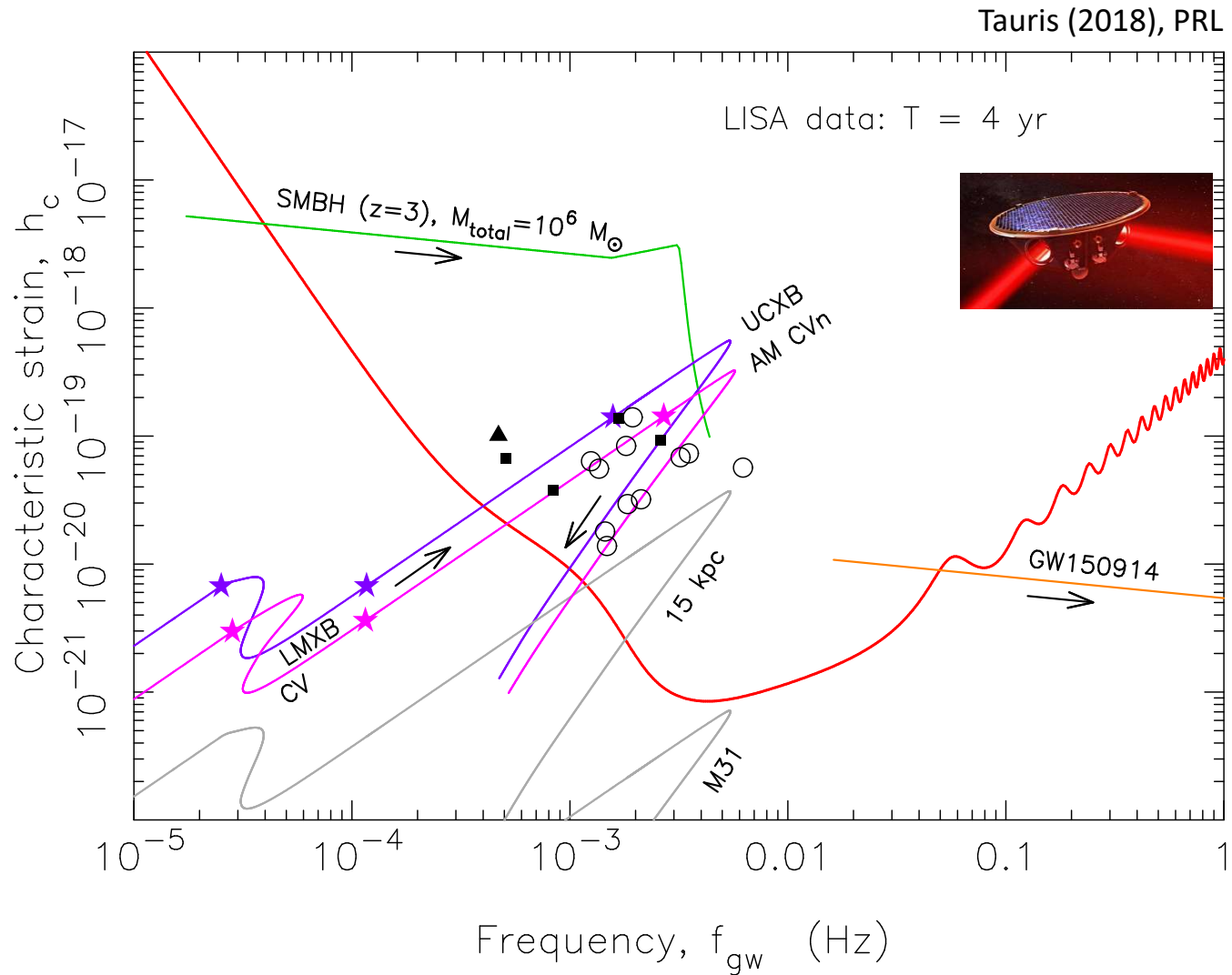


Tauris (2018), PRL



Determine NS the mass to a high accuracy via a new method

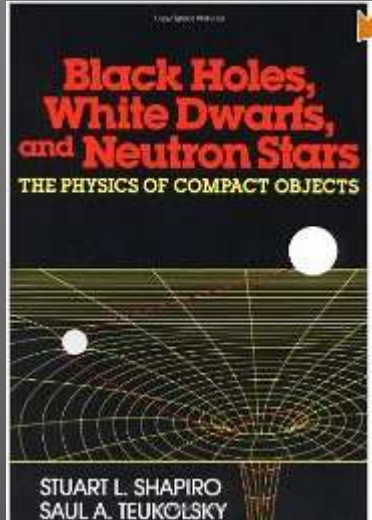




$$h_c \approx \sqrt{N_{cycles}} \sqrt{2} h_0 = \sqrt{2 f_{gw} T_{obs}} h_0 \quad h_0 = \sqrt{\frac{32}{80}} \frac{\pi^{2/3} G^{5/3} f_{gw}^{2/3} M_{chirp}^{5/3}}{c^4 d_L}$$

Physics of Compact Objects

week 10



Shapiro & Teukolsky (1983), Wiley-Interscience

Curriculum

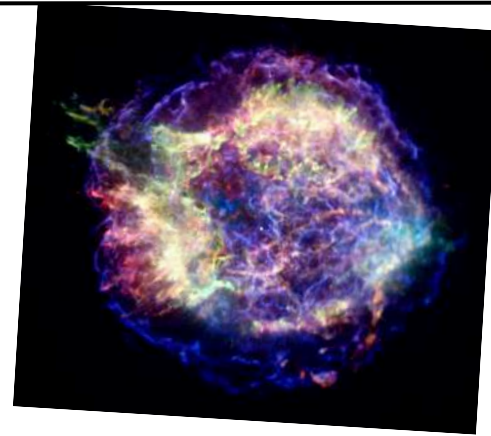
- Lecture notes
- Tauris & van den Heuvel (2023), Chapter 15 (Shapiro & Teukolsky Chapter 16) (Riles 2013; Colpi & Senasa 2017) (LIGO-Virgo-KAGRA: GWTC-3: [arXiv:2111.03634](https://arxiv.org/abs/2111.03634))

Exercises: # 7, 8

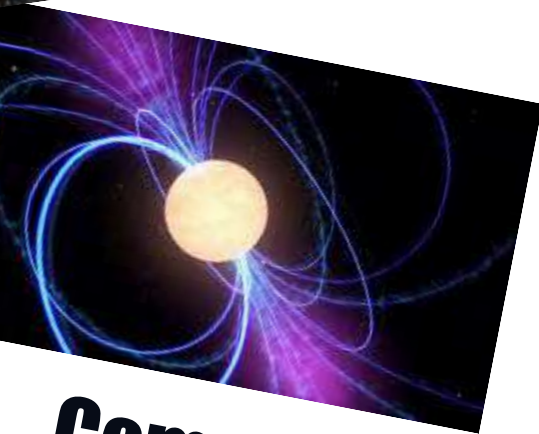
- Monday Nov. 13, 10:15-12:00
+ course evaluation



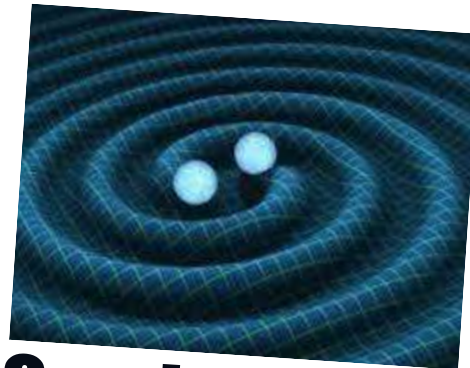
Stars



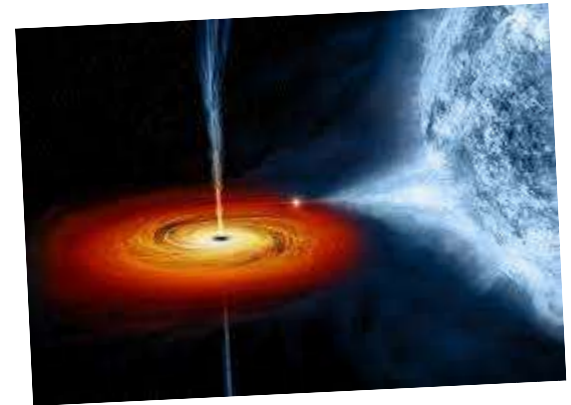
SNe



**Compact
Objects**

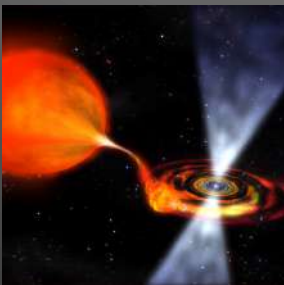
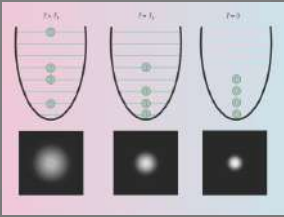


**Gravitational
Waves**

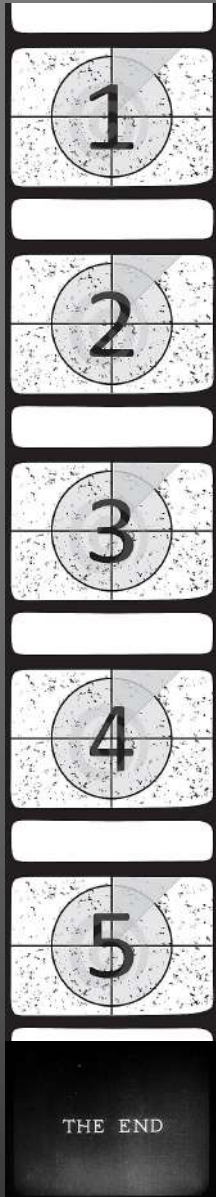
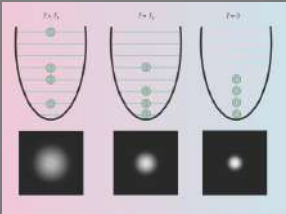


**Binary
Interactions**

Programme



- * **Introduction**
- * **Degenerate Fermi Gases**
Non-relativistic and extreme relativistic electron / (n,p,e⁻) gases
- * **White Dwarfs**
Structure, cooling models, observations
- * **Neutron Stars**
Structure and equation-of-state
- * **Radio Pulsars**
Characteristics, spin evolution, magnetars, observations
- * **Binary Evolution and Interactions**
X-ray binaries, accretion, formation of millisecond pulsars, recycling
- * **Black Holes**
Observations, characteristics and spins
- * **Gravitational Waves**
Sources and detection, kilonovae
- * **Exam**



Thanks for joining!
Remember your evaluations